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RELIABLE APPROXIMATION OF LONG
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MOLECULAR DYNAMICS

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RELIABLE APPROXIMATION OF LONG RELAXATION TIMESCALES IN MOLECULAR DYNAMICS

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Abstract. Many interesting rare events in molecular systems like ligand association, protein folding or conformational changes happen on timescales that often are not accessible by direct numerical simulation. Therefore rare event approximation approaches like interface sampling, Markov state model building or advanced reaction coordinate based free energy estimation have attracted huge attention recently. In this article we analyze the reliability of such approaches: How precise is an estimate of long relaxation timescales of molecular systems resulting from various forms of rare event approximation methods? Our results give a theoretical answer to this question by relating it with the transfer operator approach to molecular dynamics. By doing so they also allow for understanding deep connections between the different approaches.

Key words. molecular dynamics, eigenvalue problem, effective dynamics, Galerkin method, variational approach, Markov state model, reaction coordinate.

AMS subject classifications. 60J60, 60J35, 37A60

1. Introduction. The problem of accurate estimation of long relaxation timescales associated with rare events in molecular dynamics like ligand association, protein folding or conformational changes has attracted a lot of attention recently. Often these timescales are not accessible by direct numerical simulation. Therefore, different *discrete coarse graining approaches* for their approximation like Markov state model (MSM) building [30, 25] or time-lagged independent component analysis (TiCA) [26, 24] have been introduced and successfully applied to various molecular systems [22, 2]. These approaches are based on finite-dimensional Galerkin discretization [30] or variational approximation [21, 23] of the transfer operator of the molecular dynamics process [33]. In several theoretical studies the approximation error of these numerical techniques regarding the longest relaxation timescales has been analyzed resulting in error estimates in terms of the dominant eigenvalues of the transfer operator [33, 26]. In this article we first show how to get similar error estimates when replacing the transfer operator by the infinitesimal generator [32] associated with it. Furthermore, the analysis exhibits that the different approaches are deeply connected, that is, they finally lead to the identical numerical problem.

In addition to the different discrete coarse graining approaches the literature contains various alternative *reaction coordinate sampling approaches* aiming at approximation of very long relaxation processes: In these sampling approaches, one assumes that the effective dynamical behavior of the systems on long timescales can be described by a relatively low dimensional object given by some reaction coordinates. Various advanced methods such as Umbrella Sampling [34, 14], Metadynamics [16, 15], Blue-Moon sampling [3], Adaptive Biasing Force method [5], or Temperature-Accelerated Molecular Dynamics [18], as well as trajectory-based techniques like Milestoning [7], Transition Interface Sampling [20], or Forward Flux Sampling [1] may serve as some examples. These methods result in free energy barriers, transition rates or first mean passage times for the rare events of interest; they are complemented by several approaches to the effective dynamics on the reaction coordinate space [17, 8, 35] that allow for

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significantly faster simulation of these rare events [11, 13, 12] including details of the underlying molecular mechanisms. Surprisingly, our analytic tools, originally developed for discrete coarse graining approaches, can also be utilized for evaluating the approximation quality of reaction coordinate sampling approaches to the effective dynamics: We derive an explicit error estimate for the longest timescale resulting from the choice of specific reaction coordinates.

However, estimating the approximation quality is not the only way of utilizing the analytical insights presented in this article. We also demonstrate how the new techniques for simulation of the effective dynamics can be used for efficient MSM building or TiCA applications.

Mathematically, the article is based on the analysis of the dominant timescales of reversible and ergodic diffusion processes in energy landscapes. The leading eigenvalues of the transfer operator and the corresponding eigenfunctions characterize the dynamical behavior of the process on long timescales [33]. In several articles the approximation error regarding these leading eigenvalues under discretization of the transfer operator have been discussed, cf. [6, 21, 23, 26, 28, 29]. Following this work, we characterize the approximation quality for the (low-lying) eigenvalues of the infinitesimal generator. This permits us to study the connection between the effective dynamics considered in [35] and Galerkin discretization schemes for the transfer operator. Secondly, following the work [21, 23], we study the variational approach for the infinitesimal generator. In fact, we will see that this approach leads to the same generalized matrix eigenvalue problem as the one resulting from Galerkin discretization. Thirdly, numerical issues related to the estimation of the coefficient matrices by means of the effective dynamics are discussed.

The paper is organized as follows. In Section 2, we introduce the various operators associated to the reversible diffusion processes and discuss the relation between eigenvalues and relaxation timescales. Next, in Section 3, we study the Galerkin discretization of generators/-transfer operators for solving the eigenvalue problem and show that previous results can be extended to reaction coordinate subspaces. In Section 4, the variational approach to the approximation of the eigenvalue problem is considered and its relations to the Galerkin approach are worked out in detail. Then, in Section 5, we discuss numerical issues related to estimating the discretization matrices by means of simulating the effective dynamics for given reaction coordinates; the performance of this approach is studied numerically in Section 6. Finally, conclusions and some further remarks are given in Section 7.

2. Diffusion process and the associated operators. We consider a diffusion process given by the stochastic differential equation (SDE)

$$\begin{aligned} dx_s &= -\nabla V(x_s)ds + \sqrt{2\beta^{-1}}dw_s, \quad s \geq 0, \\ x_0 &= x, \end{aligned} \tag{2.1}$$

where $x_s \in \mathbb{R}^n$, parameter $\beta > 0$ is related to the inverse of system's temperature, and w_s is an n -dimensional Brownian motion. $V : \mathbb{R}^n \rightarrow \mathbb{R}$ is a potential function which is assumed to be smooth and bounded from below. The results presented in the subsequent can be extended to more general reversible diffusion processes with a state-dependent noise intensity matrix, cf. [35]. However, for the sake of simplicity of presentation we restrict our considerations to the specific case (2.1) typically studied in molecular dynamics.

The infinitesimal generator of the dynamics (2.1) is given by

$$\mathcal{L} = -\nabla V \cdot \nabla + \frac{1}{\beta} \Delta. \quad (2.2)$$

It is known that, under mild conditions on V , the solution process $(x_s)_{s \geq 0}$ of (2.1) is ergodic [19], and its unique invariant measure π is given by $\pi(dx) = \rho(x)dx$ where

$$\rho(x) = \frac{1}{Z} e^{-\beta V(x)}, \quad \text{with } Z = \int_{\mathbb{R}^n} e^{-\beta V(x)} dx. \quad (2.3)$$

We introduce the Hilbert space $H = L^2(\mathbb{R}^n, \pi)$, which is endowed with the inner product

$$\langle f, g \rangle_\pi = \int_{\mathbb{R}^n} f(x) g(x) \rho(x) dx, \quad \forall f, g \in H, \quad (2.4)$$

and the norm $\|f\|_\pi = \sqrt{\langle f, f \rangle_\pi}$, $\forall f \in H$. The domain of the operator \mathcal{L} will be denoted as $D(\mathcal{L}) \subset H$.

It is also known that the process $(x_s)_{s \geq 0}$ is a reversible process and that \mathcal{L} is a self-adjoint operator with respect to the inner product (2.4). Whenever the potential V grows to infinity fast enough at infinity, its spectrum is discrete [33]. Let $\lambda_i \in \mathbb{C}$ and $\varphi_i \in D(\mathcal{L})$ be the eigenvalues and the corresponding (normalized) eigenfunctions of $-\mathcal{L}$, that is, the solutions of the eigenproblem

$$-\mathcal{L}f = \lambda f \quad (2.5)$$

in H , or in weak form,

$$-\langle \mathcal{L}f, g \rangle_\pi = \lambda \langle f, g \rangle_\pi, \quad \forall g \in H. \quad (2.6)$$

Due to the self-adjointness of \mathcal{L} and the fact

$$\langle -\mathcal{L}f, f \rangle_\pi = \frac{1}{\beta} \int_{\mathbb{R}^n} |\nabla f(x)|^2 \rho(x) dx \geq 0, \quad \forall f \in H, \quad (2.7)$$

we can assume that $\lambda_i \in \mathbb{R}$ with

$$0 = \lambda_0 < \lambda_1 \leq \dots \leq \lambda_k \leq \dots, \quad (2.8)$$

with $\varphi_0 \equiv 1$.

Given $s \geq 0$, we define the operator $T_s : H \rightarrow H$ by

$$(T_s f)(x) = \mathbf{E}(f(x_s) \mid x_0 = x), \quad f \in H, \quad (2.9)$$

where \mathbf{E} denotes the expectation taken with respect to the paths of (2.1) under the initial condition that $x_0 = x$. It is well known that $u(s, x) = T_s f(x)$ is the solution of the Kolmogorov backward equation

$$\frac{d}{ds} u(s, \cdot) = \mathcal{L}u(s, \cdot), \quad u(0, \cdot) = f, \quad (2.10)$$

that is, the operators T_s , $s \geq 0$ form a one-parameter semigroup whose infinitesimal generator is \mathcal{L} , and thus they are self-adjoint in H as well. Because of (2.10), the formal expression $T_s = e^{s\mathcal{L}}$

is often used in the literature. Similarly to (2.8), we also know that the eigenvalues of T_s are given by

$$1 = e^{-\lambda_0 s} > e^{-\lambda_1 s} \geq \dots > 0, \quad (2.11)$$

with the same eigenfunctions φ_i , $i = 0, 1, \dots$. Because of (2.11), the relaxation timescales t_i of the dynamics (2.1) are given by [32]

$$t_i = \lambda_i^{-1}, \quad i = 1, 2, \dots$$

In the following we introduce another operator called *transfer operator*, which has been extensively considered in the literature to investigate the metastability of molecular systems and to build Markov state models (MSM) [30, 33, 2]. Fix a lag time $\tau > 0$, let $p(x, \cdot; \tau)$ be the transition density function of process (2.1) starting from $x \in \mathbb{R}^n$, i.e. $p(x, y; \tau)$ describes the probability density of starting from state x at time 0 and arriving at $y \in \mathbb{R}^n$ after time τ . For a bounded and continuous function $u \in H$, the transfer operator $\mathcal{T}_\tau : H \rightarrow H$ is defined by [27, 30, 31]

$$(\mathcal{T}_\tau u)(y) = \frac{1}{\rho(y)} \int_{\mathbb{R}^n} p(x, y; \tau) u(x) \rho(x) dx, \quad y \in \mathbb{R}^n. \quad (2.12)$$

From (2.12), it follows immediately that

$$\begin{aligned} \langle \mathcal{T}_\tau u, f \rangle_\pi &= \int_{\mathbb{R}^n} \int_{\mathbb{R}^n} p(x, y; \tau) f(y) dy u(x) \rho(x) dx \\ &= \int_{\mathbb{R}^n} \left[\mathbf{E}(f(x_\tau) | x_0 = x) \right] u(x) \rho(x) dx \\ &= \langle u, T_\tau f \rangle_\pi = \langle T_\tau u, f \rangle_\pi, \quad \forall f \in H. \end{aligned}$$

Therefore, we can conclude that $\mathcal{T}_\tau = T_\tau$, i.e. the transfer operator \mathcal{T}_τ coincides with the operator T_τ , a member within the semigroup $(T_s)_{s \geq 0}$. Denote the eigenvalues of \mathcal{T}_τ as μ_i , $i \geq 0$, such that

$$1 = \mu_0 > \mu_1 \geq \dots > 0. \quad (2.13)$$

Then from (2.11) we know that $\mu_i = e^{-\lambda_i \tau}$ and the corresponding eigenfunctions are the same as the eigenfunctions φ_i of the infinitesimal generator \mathcal{L} . This means that the dominant relaxation time scales of the dynamics (2.1) can be obtained by computing the dominant eigenvalues of \mathcal{T}_τ , cf. [32, 27]

3. Galerkin approximation of the eigenvalues of the generator. In this section, we study the Galerkin method for computing the eigenvalues of the infinitesimal generator \mathcal{L} . While Galerkin discretization of the transfer operator has been studied to some extent [33], results on the associated generator are rather sparse.

3.1. Some general results. For introducing the Galerkin method, let H_0 be a Hilbert subspace of H containing the constant function, and let \mathbf{P} denote the orthogonal projection operator from H to H_0 , which satisfies $\mathbf{P}^2 = \mathbf{P}$ and

$$\langle \mathbf{P}f, g \rangle_\pi = \langle f, g \rangle_\pi, \quad \forall f \in H, g \in H_0. \quad (3.1)$$

Then the Galerkin method aims at approximating the solution of (2.6) in the subspace H_0 . Specifically, we want to find $f \in H_0$, such that

$$-\langle \mathcal{L}f, g \rangle_\pi = \kappa \langle f, g \rangle_\pi, \quad \forall g \in H_0, \quad (3.2)$$

for some constant $\kappa \geq 0$. Using the property (3.1), we know that problem (3.2) is equivalent to the eigenvalue problem for operator $-\mathbf{P}\mathcal{L}$ on the subspace H_0 , i.e.

$$-\mathbf{P}\mathcal{L}f = \kappa f. \quad (3.3)$$

Notice that $-\mathbf{P}\mathcal{L}$ is a self-adjoint operator on H_0 . Similarly as in (2.8), let $\zeta_i \in H_0$ be the orthonormal eigenfunctions of operator $-\mathbf{P}\mathcal{L}$ corresponding to eigenvalues κ_i , where

$$0 = \kappa_0 < \kappa_1 \leq \kappa_2 \leq \dots, \quad (3.4)$$

and $\zeta_0 \equiv 1$. When H_0 is an infinite dimensional subspace, we assume $\kappa_i \rightarrow +\infty$ as $i \rightarrow +\infty$.

First we want to study the condition under which the eigenvalues of the projected generator $\mathbf{P}\mathcal{L}$ are reliable approximations of the eigenvalues of the full generator \mathcal{L} . The following approximation result was obtained in [35]:

THEOREM 3.1. *For $i \geq 0$, let φ_i and ζ_i be the normalized eigenfunctions of operators $-\mathcal{L}$ and $-\mathbf{P}\mathcal{L}$ corresponding to eigenvalues λ_i and κ_i , respectively. We have*

$$\lambda_i \leq \kappa_i \leq \lambda_i + \frac{1}{\beta} \int_{\mathbb{R}^n} |\nabla(\varphi_i - \zeta_i)(x)|^2 \rho(x) dx. \quad (3.5)$$

We include its proof for completeness:

Proof. From (3.2), we have $\kappa_i = -\langle \mathcal{L}\zeta_i, \zeta_i \rangle_\pi$. Define the subspace $E_{i+1} = \text{span}\{\zeta_0, \dots, \zeta_i\}$ for $i \geq 0$. It follows from the orthogonality of functions ζ_i that E_{i+1} is an $(i+1)$ -dimensional subspace of H . And using (3.4) it is direct to verify that

$$\kappa_i = \max_{f \in E_{i+1}, |f|_\pi=1} \langle -\mathcal{L}f, f \rangle_\pi. \quad (3.6)$$

Applying the Min-Max theorem to the eigenvalues of operator $-\mathcal{L}$, we conclude

$$\kappa_i = \max_{f \in E_{i+1}, |f|_\pi=1} \langle -\mathcal{L}f, f \rangle_\pi \geq \min_{E'_{i+1}} \max_{f \in E'_{i+1}, |f|_\pi=1} \langle -\mathcal{L}f, f \rangle_\pi = \lambda_i, \quad (3.7)$$

where E'_{i+1} goes over all $(i+1)$ -dimensional subspaces of H . For the upper bound, we can compute that

$$\begin{aligned} & \langle -\mathcal{L}(\varphi_i - \zeta_i), (\varphi_i - \zeta_i) \rangle_\pi \\ &= \langle -\mathcal{L}\varphi_i, \varphi_i \rangle_\pi + 2\langle \mathcal{L}\varphi_i, \zeta_i \rangle_\pi + \langle -\mathcal{L}\zeta_i, \zeta_i \rangle_\pi \\ &= \lambda_i - 2\lambda_i \langle \varphi_i, \zeta_i \rangle_\pi + \kappa_i \\ &= \kappa_i - \lambda_i + 2\lambda_i(1 - \langle \varphi_i, \zeta_i \rangle_\pi) \geq \kappa_i - \lambda_i, \end{aligned}$$

where we have used the fact that $\langle \varphi_i, \zeta_i \rangle_\pi \leq |\varphi_i|_\pi |\zeta_i|_\pi = 1$. The conclusion follows from (2.7). \square

Previous studies on the Galerkin approximation of the dominant eigenvalues of the transfer operator have shown that the approximation error of eigenvalues can be reliably bounded by means of the projection errors of the corresponding eigenfunctions [6, 28, 29]. Next we will

derive a similar result for the generator \mathcal{L} . To this end, we introduce the orthogonal projection \mathbf{P}^\perp from H to the complement subspace H_0^\perp of H_0 , that is, $\mathbf{P}^\perp = \mathbf{I} - \mathbf{P}$. We have

THEOREM 3.2. *Let φ be a normalized eigenfunction of the operator $-\mathcal{L}$ corresponding to the eigenvalue λ . Define constants*

$$\delta_1 = |\mathcal{L}\mathbf{P}^\perp\varphi|_\pi, \quad \delta_2 = |\mathbf{P}^\perp\varphi|_\pi, \quad (3.8)$$

and suppose that $0 < \delta_2 < 1$. Then there is an eigenvalue κ_i of the operator $-\mathbf{P}\mathcal{L}$, such that

$$|\kappa_i - \lambda| \leq \frac{\delta_1}{(1 - \delta_2^2)^{\frac{1}{2}}}. \quad (3.9)$$

Proof. Since $\delta_2 = |\mathbf{P}^\perp\varphi|_\pi = (1 - |\mathbf{P}\varphi|_\pi^2)^{\frac{1}{2}} < 1$, we have $|\mathbf{P}\varphi|_\pi > 0$. Let $\mathbf{P}\varphi = \sum_{i=0}^{+\infty} \omega_i \zeta_i$, where $\omega_i = \langle \varphi, \zeta_i \rangle_\pi$, and the summation consists of finite terms when H_0 is a finite dimensional space. For all $g \in H_0$, we can compute

$$\begin{aligned} \langle \mathbf{P}\mathcal{L}\mathbf{P}^\perp\varphi, g \rangle_\pi &= \langle \mathbf{P}\mathcal{L}(\varphi - \mathbf{P}\varphi), g \rangle_\pi \\ &= \langle \mathbf{P}\mathcal{L}\varphi, g \rangle_\pi - \langle \mathbf{P}\mathcal{L}\left(\sum_{i=0}^{+\infty} \omega_i \zeta_i\right), g \rangle_\pi \\ &= -\langle \lambda\mathbf{P}\varphi, g \rangle_\pi + \left\langle \sum_{i=0}^{+\infty} \omega_i \kappa_i \zeta_i, g \right\rangle_\pi = \left\langle \sum_{i=0}^{+\infty} \omega_i (\kappa_i - \lambda) \zeta_i, g \right\rangle_\pi. \end{aligned}$$

Therefore we have obtained $\mathbf{P}\mathcal{L}\mathbf{P}^\perp\varphi = \sum_{i=0}^{+\infty} \omega_i (\kappa_i - \lambda) \zeta_i$, and

$$|\mathbf{P}\mathcal{L}\mathbf{P}^\perp\varphi|_\pi^2 = \sum_{i=0}^{+\infty} \omega_i^2 |\kappa_i - \lambda|^2 \geq \left(\min_i |\kappa_i - \lambda|^2 \right) \sum_{i=0}^{+\infty} \omega_i^2 = \min_i |\kappa_i - \lambda|^2 |\mathbf{P}\varphi|_\pi^2.$$

And it follows that

$$\min_i |\kappa_i - \lambda| \leq \frac{|\mathbf{P}\mathcal{L}\mathbf{P}^\perp\varphi|_\pi}{|\mathbf{P}\varphi|_\pi} \leq \frac{|\mathcal{L}\mathbf{P}^\perp\varphi|_\pi}{(1 - |\mathbf{P}^\perp\varphi|_\pi^2)^{\frac{1}{2}}} = \frac{\delta_1}{(1 - \delta_2^2)^{\frac{1}{2}}}. \quad (3.10)$$

□

REMARK 1. *Notice that our error bound relies on both constants δ_1, δ_2 , while the error bound in [29] for the transfer operator only depends on one constant, the projection error δ_2 . This difference is due to the fact that the generator \mathcal{L} is an unbounded operator while the transfer operator is bounded.*

3.2. Finite dimensional subspaces. In applications, it is often assumed that H_0 is spanned by finitely many basis functions. Especially, this is the situation when constructing MSMs based on indicator functions of partition sets [29] or based on core sets [32].

Let H_0 be the finite dimensional space $H_0 = \text{span}\{\psi_1, \psi_2, \dots, \psi_N\}$ and consider the problem (3.2). Applying Theorem 3.1 and Theorem 3.2, we have

COROLLARY 3.3. *For Galerkin approximation of the eigenproblem using finite-dimensional ansatz spaces the following three statements are valid:*

1. Write $f = \sum_{i=1}^N \omega_i \psi_i \in H_0$ and let $X = (\omega_1, \omega_2, \dots, \omega_N)^T \in \mathbb{R}^N$. Then problem (3.2) is equivalent to the generalized matrix eigenvalue equation

$$CX = \lambda SX, \quad (3.11)$$

where C, S are $N \times N$ matrices whose entries are given by

$$C_{ll'} = \langle -\mathcal{L}\psi_l, \psi_{l'} \rangle_\pi, \quad S_{ll'} = \langle \psi_l, \psi_{l'} \rangle_\pi, \quad 1 \leq l, l' \leq N, \quad (3.12)$$

2. Let $0 = \kappa_0 \leq \kappa_1 \leq \dots \leq \kappa_k$ be the first $(k+1)$ smallest eigenvalues of problem (3.11) and

$$X_i = (X_{i1}, X_{i2}, \dots, X_{iN})^T, \quad 1 \leq i \leq k,$$

be the eigenvector corresponding to κ_i such that $X_i^T S X_i = 1$. Define $\zeta_i = \sum_{l=1}^N X_{il} \psi_l$, then we have

$$\lambda_i \leq \kappa_i \leq \lambda_i + \frac{1}{\beta} \int_{\mathbb{R}^n} |\nabla(\varphi_i - \zeta_i)(x)|^2 \rho(x) dx, \quad 1 \leq i \leq k, \quad (3.13)$$

where λ_i, φ_i are the eigenvalue and the eigenfunction of the operator $-\mathcal{L}$, respectively.

3. Let \mathbf{P} be the orthogonal projection operator from H to H_0 , and φ be an eigenfunction of the operator $-\mathcal{L}$ corresponding to eigenvalue λ . Define constants

$$\delta_1 = |\mathcal{L}\mathbf{P}^\perp \varphi|_\pi, \quad \delta_2 = |\mathbf{P}^\perp \varphi|_\pi,$$

and suppose that $\delta_2 < 1$. Then there is an eigenvalue κ_i of problem (3.11) such that

$$|\kappa_i - \lambda| \leq \frac{\delta_1}{(1 - \delta_2^2)^{\frac{1}{2}}}. \quad (3.14)$$

3.3. Infinite dimensional subspace: effective dynamics. In this subsection, we discuss Galerkin approximations based on infinitely-dimensional ansatz spaces; these cases appear when studying the effective dynamics given by a so-called reaction coordinate, cf. [35]. In order to explain the relation between Galerkin approximation and effective dynamics let us first recall some definitions and results regarding the effective dynamics. For more details, readers are referred to [17, 9] for related work.

Let $\xi : \mathbb{R}^n \rightarrow \mathbb{R}^m$ be a reaction coordinate function, $m \geq 1$. For any function $f \in H$ and $x \in \mathbb{R}^n$, we define

$$\mathbf{P}f(x) = \frac{1}{Q(z)} \int_{\mathbb{R}^n} \rho(x') f(x') \delta(\xi(x') - z) dx', \quad (3.15)$$

where $z = \xi(x) \in \mathbb{R}^m$, $\delta(\cdot)$ denotes the delta function, and $Q(z) = \int_{\mathbb{R}^n} \rho(x') \delta(\xi(x') - z) dx'$ is a normalization factor satisfying $\int_{\mathbb{R}^m} Q(z) dz = 1$. Define the probability measure ν on \mathbb{R}^m given by $\nu(dz) = Q(z) dz$ for $z \in \mathbb{R}^m$ and consider the Hilbert space $\tilde{H} = L^2(\mathbb{R}^m, \nu)$. \tilde{H} induces a (infinite dimensional) subspace of H , namely,

$$H_0 = \left\{ f \mid f \in H, f = \tilde{f} \circ \xi, \text{ for some } \tilde{f} \in \tilde{H} \right\} \subset H, \quad (3.16)$$

and (3.15) clearly implies that $\mathbf{P}f \in H_0$.

Let $\tilde{f} \in \tilde{H}$ satisfy $\mathbf{P}f = \tilde{f} \circ \xi$. Then, using (3.15), we can verify that $\mathbf{P}^2 = \mathbf{P}$ and

$$\langle f, h \rangle_\pi = \langle \mathbf{P}f, h \rangle_\pi = \langle \tilde{f}, \tilde{h} \rangle_\nu, \quad \forall h = \tilde{h} \circ \xi \in H_0. \quad (3.17)$$

Therefore the mapping $\mathbf{P} : H \rightarrow H_0$ actually is the orthogonal projection operator from H to the subspace H_0 . For $f \in H$, $z \in \mathbb{R}^m$, in the following we will also write $\mathbf{P}f(z)$ instead of $\tilde{f}(z)$, where $\tilde{f} \in \tilde{H}$ such that $\mathbf{P}f = \tilde{f} \circ \xi$. The effective dynamics for the reaction coordinate ξ is defined on \mathbb{R}^m and satisfies the SDE

$$dz_s = \tilde{b}(z_s) ds + \sqrt{2\beta^{-1}\tilde{\sigma}(z_s)} dw_s, \quad (3.18)$$

where $z_s \in \mathbb{R}^m$, w_s is a Brownian motion on \mathbb{R}^m , and the coefficients $\tilde{b} : \mathbb{R}^m \rightarrow \mathbb{R}^m$, $\tilde{\sigma} : \mathbb{R}^m \rightarrow \mathbb{R}^{m \times m}$ are given by

$$\begin{aligned} \tilde{b}_l(z) &= \mathbf{P}(\mathcal{L}\xi_l)(z) = \mathbf{P}\left(-\nabla V \cdot \nabla \xi_l + \frac{1}{\beta} \Delta \xi_l\right)(z), \\ \tilde{a}_{ll'}(z) &= (\tilde{\sigma}\tilde{\sigma}^T)_{ll'}(z) = \mathbf{P}\left(\sum_{i=1}^n \frac{\partial \xi_l}{\partial x_i} \frac{\partial \xi_{l'}}{\partial x_i}\right)(z), \end{aligned} \quad (3.19)$$

for $\forall z \in \mathbb{R}^m$, $1 \leq l, l' \leq m$. The infinitesimal generator of the process governed by (3.18) is given by

$$\tilde{\mathcal{L}} = \sum_{l=1}^m \tilde{b}_l \frac{\partial}{\partial z_l} + \frac{1}{\beta} \sum_{l, l'=1}^m \tilde{a}_{ll'} \frac{\partial^2}{\partial z_l \partial z_{l'}}, \quad (3.20)$$

which is a self-adjoint operator on space \tilde{H} with discrete spectrum under appropriate conditions on ξ . We consider the eigenvalue problem

$$-\tilde{\mathcal{L}}\tilde{f} = \tilde{\lambda}\tilde{f}, \quad \tilde{f} \in \tilde{H}, \quad (3.21)$$

and let $\tilde{\varphi}_i \in \tilde{H}$ be the orthonormal eigenfunctions of the operator $-\tilde{\mathcal{L}}$ corresponding to eigenvalues $\tilde{\lambda}_i$, where

$$0 = \tilde{\lambda}_0 < \tilde{\lambda}_1 \leq \tilde{\lambda}_2 \leq \dots. \quad (3.22)$$

Applying Theorem 3.1 and Theorem 3.2, we have the following result.

COROLLARY 3.4. *For the eigenproblem associated with the effective dynamics the following three statements are valid:*

1. For $f = \tilde{f} \circ \xi \in H_0$ where $\tilde{f} \in \tilde{H}$, we have

$$\mathbf{P}\mathcal{L}f = (\tilde{\mathcal{L}}\tilde{f}) \circ \xi. \quad (3.23)$$

2. Let φ_i and $\tilde{\varphi}_i$ be the normalized eigenfunctions of operator $-\mathcal{L}$ and $-\tilde{\mathcal{L}}$ corresponding to eigenvalues λ_i and $\tilde{\lambda}_i$, respectively. We have

$$\lambda_i \leq \tilde{\lambda}_i \leq \lambda_i + \frac{1}{\beta} \int_{\mathbb{R}^n} |\nabla(\varphi_i - \tilde{\varphi}_i \circ \xi)(x)|^2 \rho(x) dx. \quad (3.24)$$

3. Let φ be the normalized eigenfunction of operator $-\mathcal{L}$ corresponding to eigenvalue λ . Define constants

$$\delta_1 = |\mathcal{L}\mathbf{P}^\perp\varphi|_\pi, \quad \delta_2 = |\mathbf{P}^\perp\varphi|_\pi,$$

and suppose $\delta_2 < 1$. Then there is an eigenvalue $\tilde{\lambda}_i$ of problem (3.21), such that

$$|\tilde{\lambda}_i - \lambda| \leq \frac{\delta_1}{(1 - \delta_2^2)^{\frac{1}{2}}}. \quad (3.25)$$

Proof. The proof of the first assertion can be found in [35]. Using (3.23) and (3.17), we can derive

$$-\langle \mathbf{P}\mathcal{L}(\tilde{\varphi}_i \circ \xi), f \rangle_\pi = -\langle (\tilde{\mathcal{L}}\tilde{\varphi}_i) \circ \xi, f \rangle_\pi = -\langle \tilde{\mathcal{L}}\tilde{\varphi}_i, \tilde{f} \rangle_\nu = \tilde{\lambda}_i \langle \tilde{\varphi}_i \circ \xi, f \rangle_\pi, \quad \forall f = \tilde{f} \circ \xi \in H_0, \quad (3.26)$$

i.e. $\tilde{\lambda}_i$ and $\tilde{\varphi}_i \circ \xi$ are the eigenvalues and eigenfunctions of the projected operator $-\mathbf{P}\mathcal{L}$ on the subspace H_0 , respectively. Furthermore, $\langle \tilde{\varphi}_i \circ \xi, \tilde{\varphi}_i \circ \xi \rangle_\pi = \langle \tilde{\varphi}_i, \tilde{\varphi}_i \rangle_\nu = 1$, i.e. $\tilde{\varphi}_i$ is normalized. Therefore, the second assertion is implied by Theorem 3.1. The third assertion follows from Theorem 3.2 in the same way. \square

REMARK 2. As an interesting conclusion of the first assertion we see that, on the infinitesimal subspace H_0 defined in (3.16), the projected operator $-\mathbf{P}\mathcal{L}$ is essentially described by another differential operator $\tilde{\mathcal{L}}$, which is defined in the Hilbert space \tilde{H} and coincides with the infinitesimal generator of the effective dynamics on \mathbb{R}^m .

4. Variational approach to generator eigenproblem. In this section, we study the variational approach to approximate the eigenvalues and eigenfunctions of the operator $-\mathcal{L}$. This approach has been considered in [21, 24, 23] to study the related eigenproblem of the transfer operator. Its main idea is to approximate the dominant eigenvalues of a self-adjoint transfer operator via an appropriate form of the Rayleigh variational principle instead via Galerkin discretization [21]. Herein, we present a similar approach to the low-lying generator eigenvalues.

4.1. Variational principle. The main object of the variational approach is the following functional $\mathcal{F} : D(\mathcal{L})^{\oplus(k+1)} \rightarrow \mathbb{R}$ that acts on $k+1$ functions from $D(\mathcal{L})$.

Given arbitrary constants $\omega_i > 0$, $0 \leq i \leq k$, the functional is defined by

$$\mathcal{F}(f_0, f_1, \dots, f_k) = \sum_{i=0}^k \omega_i \langle -\mathcal{L}f_i, f_i \rangle_\pi, \quad f_i \in D(\mathcal{L}), \quad 0 \leq i \leq k. \quad (4.1)$$

Clearly, for the (normalized) leading eigenfunctions φ_i of \mathcal{L} , we have

$$\mathcal{F}(\varphi_0, \varphi_1, \dots, \varphi_k) = \sum_{i=0}^k \omega_i \lambda_i,$$

where λ_i are the corresponding eigenvalues of $-\mathcal{L}$. The main working horse of the variational principle is the following lower and upper bound:

THEOREM 4.1 (Variational principle). Let ω_i , $i = 0, 1, \dots, k$ be a decreasing sequence of positive real numbers, i.e. $\omega_0 > \omega_1 > \dots > \omega_k > 0$. For any orthonormal family of functions $f_i \in D(\mathcal{L})$, $i = 0, 1, \dots, k$, we have

$$\mathcal{F}(\varphi_0, \varphi_1, \dots, \varphi_k) \leq \mathcal{F}(f_0, f_1, \dots, f_k) \leq \mathcal{F}(\varphi_0, \varphi_1, \dots, \varphi_k) + \mathcal{F}(f_0 - \varphi_0, f_1 - \varphi_1, \dots, f_k - \varphi_k), \quad (4.2)$$

or more explicitly,

$$\sum_{i=0}^k \omega_i \lambda_i \leq \sum_{i=0}^k \omega_i \langle -\mathcal{L}f_i, f_i \rangle_\pi \leq \sum_{i=0}^k \omega_i \lambda_i + \sum_{i=0}^k \omega_i \langle -\mathcal{L}(f_i - \varphi_i), (f_i - \varphi_i) \rangle_\pi. \quad (4.3)$$

In order to prove this variational principle we need the following simple lemma:

LEMMA 4.2. *Suppose $k > 0$, and let $(\alpha_i)_{i=0,1,\dots,k}$ and $(\omega_i)_{i=0,1,\dots,k}$ be two ordered sequences of real numbers such that*

$$\alpha_0 \leq \alpha_1 \leq \dots \leq \alpha_k, \quad \omega_0 \geq \omega_1 \geq \dots \geq \omega_k.$$

Then, for any permutation (ω'_i) of the sequence (ω_i) , we have

$$\sum_{i=0}^k \alpha_i \omega'_i \geq \sum_{i=0}^k \alpha_i \omega_i. \quad (4.4)$$

Proof. The proof of Theorem 4.1 is given in two steps:

1. For the lower bound, we consider the optimization problem

$$\begin{aligned} \min_{f_i} \mathcal{F}(f_0, f_1, \dots, f_k) &= \min_{f_i} \sum_{i=0}^k \omega_i \langle -\mathcal{L}f_i, f_i \rangle_\pi, \\ \text{subject to } \langle f_i, f_j \rangle_\pi &= \delta_{ij}, \quad 0 \leq i, j \leq k. \end{aligned} \quad (4.5)$$

Next, we introduce the Lagrange multipliers λ_{ij} for $0 \leq i < j \leq k$, and consider the auxiliary functional

$$\sum_{i=0}^k \omega_i \langle -\mathcal{L}f_i, f_i \rangle_\pi - \sum_{i=0}^k \sum_{j=i}^k \lambda_{ij} (\langle f_i, f_j \rangle_\pi - \delta_{ij}). \quad (4.6)$$

Applying calculus of variation, we conclude that the minimizer of (4.5) satisfies

$$\begin{aligned} -2\omega_i \mathcal{L}f_i - \sum_{j=i}^k \lambda_{ij} f_j - \sum_{j=0}^i \lambda_{ji} f_j &= 0, \quad \forall 0 \leq i \leq k, \\ \langle f_i, f_j \rangle_\pi &= \delta_{ij}, \quad 0 \leq i, j \leq k. \end{aligned} \quad (4.7)$$

Multiplying f_j for some $i < j \leq k$ in the first equation of (4.7) and integrating, we obtain $\lambda_{ij} = -2\omega_i \langle \mathcal{L}f_i, f_j \rangle_\pi$. In the same way we could also obtain $\lambda_{ij} = -2\omega_j \langle \mathcal{L}f_j, f_i \rangle_\pi$. Using the fact that \mathcal{L} is self-adjoint and $\omega_i > \omega_j$ for $i < j$, we conclude that

$$\lambda_{ij} = \langle \mathcal{L}f_i, f_j \rangle_\pi = 0, \quad \forall 0 \leq i < j \leq k, \quad (4.8)$$

and (4.7) reduces to an eigenvalue problem

$$-\mathcal{L}f_i = \frac{\lambda_{ii}}{\omega_i} f_i, \quad 0 \leq i \leq k. \quad (4.9)$$

Therefore the minimizer of (4.5) is given by the orthonormal eigenfunctions. Applying Lemma 4.2, we can further conclude that the lower bound is obtained when $f_i = \varphi_i$, with value

$$\sum_{i=0}^k \omega_i \langle -\mathcal{L}\varphi_i, \varphi_i \rangle_\pi = \sum_{i=0}^k \omega_i \lambda_i. \quad (4.10)$$

2. For the upper bound, similarly to the proof of Theorem 3.1, direct computation gives

$$\begin{aligned}
& \sum_{i=0}^k \omega_i \langle -\mathcal{L}(f_i - \varphi_i), (f_i - \varphi_i) \rangle_\pi \\
&= \sum_{i=0}^k \omega_i \langle -\mathcal{L}f_i, f_i \rangle_\pi - \sum_{i=0}^k \omega_i \lambda_i + 2 \sum_{i=0}^k \omega_i \lambda_i (1 - \langle f_i, \varphi_i \rangle_\pi) \\
&\geq \sum_{i=0}^k \omega_i \langle -\mathcal{L}f_i, f_i \rangle_\pi - \sum_{i=0}^k \omega_i \lambda_i,
\end{aligned}$$

where we have used the fact that $-\mathcal{L}\varphi_i = \lambda_i \varphi_i$ and $\langle f_i, \varphi_i \rangle_\pi \leq |f_i|_\pi |\varphi_i|_\pi = 1$, since both f_i, φ_i are normalized functions.

□

4.2. Optimization problem. The variational principle of Theorem 4.1 allows to approximate the low-lying eigenvalues of the generator. In order to turn it into an algorithm, we again introduce N basis functions $\psi_1, \dots, \psi_N \in D(\mathcal{L})$. We want to approximate the first $k+1$ eigenvalues λ_i as well as the eigenfunctions φ_i , $0 \leq i \leq k$ by approximating the eigenfunctions by linear combinations of the basis functions. That is, we consider functions

$$f_i = \sum_{l=1}^N x_{il} \psi_l, \quad (4.11)$$

where x_{il} are real-valued coefficients to be determined, $0 \leq i \leq k$, $1 \leq l \leq N$. Inspired by Theorem 4.1 we wish to determine the coefficients x_{il} by solving the optimization problem

$$\begin{aligned}
\min_{\{x_{il}\}} \mathcal{F}(f_0, f_1, \dots, f_k) &= \min_{\{x_{il}\}} \sum_{i=0}^k \omega_i \langle -\mathcal{L}f_i, f_i \rangle_\pi, \quad f_i = \sum_{l=1}^N x_{il} \psi_l, \\
\text{subject to } \langle f_i, f_j \rangle_\pi &= \delta_{ij}, \quad 0 \leq i, j \leq k.
\end{aligned} \quad (4.12)$$

Recall the matrices C, S defined in (3.12) and define the vectors $X_i = (x_{i0}, x_{i1}, \dots, x_{iN})^T \in \mathbb{R}^N$, $0 \leq i \leq k$. With this, the optimization problem (4.12) can be reformulated as

$$\begin{aligned}
\min_{\{x_{il}\}} \sum_{i=0}^k \omega_i \left(\sum_{1 \leq l, l' \leq N} x_{il} x_{il'} C_{ll'} \right), \\
\text{subject to } \sum_{1 \leq l, l' \leq N} x_{il} S_{ll'} x_{jl'} &= \delta_{ij}, \quad 0 \leq i, j \leq k.
\end{aligned} \quad (4.13)$$

or, equivalently, in matrix form,

$$\begin{aligned}
\min_{X_0, X_1, \dots, X_k} \sum_{i=0}^k \omega_i X_i^T C X_i, \\
\text{subject to } X_i^T S X_j &= \delta_{ij}, \quad 0 \leq i, j \leq k.
\end{aligned} \quad (4.14)$$

Using a similar argument as in the proof of Theorem 4.1, we can obtain

THEOREM 4.3. *The minimum of the optimization problem (4.12) is achieved by the functions f_i as of (4.11) with the coefficients from the first $k+1$ eigenvectors X_i of the generalized*

matrix eigenvalue problem

$$CX = \lambda SX. \quad (4.15)$$

Suppose the eigenvectors X_i of (4.15) are chosen such that $X_i^T SX_j = \delta_{ij}$ and the corresponding eigenvalues are κ_i for $0 \leq i \leq k$, where $\kappa_0 \leq \kappa_1 \leq \dots \leq \kappa_k$. Then the minimum of (4.12) is

$$\sum_{i=0}^k \omega_i X_i^T CX_i = \sum_{i=0}^k \omega_i \kappa_i. \quad (4.16)$$

REMARK 3. Combining the above result with Subsection 3.2, we see that both the Galerkin method and the variational approach lead to the same generalized matrix eigenvalue problem with an identical estimate for the eigenvalue error.

5. Numerical algorithms. In this section, we consider how the matrices C, S defined in (3.12), that is,

$$C_{ll'} = \langle -\mathcal{L}\psi_l, \psi_{l'} \rangle_\pi, \quad S_{ll'} = \langle \psi_l, \psi_{l'} \rangle_\pi, \quad 1 \leq l, l' \leq N \quad (5.1)$$

can be approximated from trajectories of the diffusion process. For the transfer operator this problem has been studied in [21, 24, 23] using trajectories of the original diffusion process given by (2.1). In contrast we herein will consider trajectories of the effective dynamics (3.18) instead of the original diffusion process.

5.1. Computing coefficient matrices using effective dynamics. Similar to the setup in Subsection 3.3, we assume that a reaction coordinate function $\xi : \mathbb{R}^n \rightarrow \mathbb{R}^m$, as well as N basis functions ψ_l , $1 \leq l \leq N$, are given. Furthermore, we suppose that the basis functions ψ_l can be written as $\psi_l = \tilde{\psi}_l \circ \xi$ for some functions $\tilde{\psi}_l \in \tilde{H}$, i.e. $\psi_l \in H_0$. In this case, it follows from the first assertion of Corollary 3.4 and the relation (3.17) that

$$\begin{aligned} S_{ll'} &= \langle \psi_l, \psi_{l'} \rangle_\pi = \langle \tilde{\psi}_l, \tilde{\psi}_{l'} \rangle_\nu, \\ C_{ll'} &= \langle -\mathcal{L}\psi_l, \psi_{l'} \rangle_\pi = \langle -\mathcal{L}\psi_l, \mathbf{P}\psi_{l'} \rangle_\pi = \langle -\mathbf{P}\mathcal{L}\psi_l, \psi_{l'} \rangle_\pi \\ &= \langle -(\tilde{\mathcal{L}}\tilde{\psi}_l) \circ \xi, \psi_{l'} \rangle_\pi = \langle -\tilde{\mathcal{L}}\tilde{\psi}_l, \tilde{\psi}_{l'} \rangle_\nu. \end{aligned} \quad (5.2)$$

These equalities are simple but interesting, because they show that the entries of coefficient matrices C, S are related to the infinitesimal generator $\tilde{\mathcal{L}}$ of the effective dynamics in (3.20). Since it is known that ν is the unique invariant measure of the effective dynamics [35], we can apply the ergodic theorem and get

$$S_{ll'} = \lim_{T \rightarrow +\infty} \frac{1}{T} \int_0^T \tilde{\psi}_l(z_s) \tilde{\psi}_{l'}(z_s) ds \approx \frac{1}{M - M_0} \sum_{i=M_0+1}^M \tilde{\psi}_l(z_{i\Delta t}) \tilde{\psi}_{l'}(z_{i\Delta t}), \quad (5.3)$$

where z_s denotes a realization of the effective dynamics (3.18), $\Delta t > 0$ is the step size, $M \in \mathbb{N}$ is a large integer, and only the parts of trajectories after time $M_0\Delta t$ are used for estimation.

For the matrix C , by using (5.2), the definition of the infinitesimal generator $\tilde{\mathcal{L}}$, as well as

the ergodic theorem, we can derive

$$\begin{aligned}
C_{ll'} &= \langle -\tilde{\mathcal{L}}\tilde{\psi}_l, \tilde{\psi}_{l'} \rangle_\nu \\
&= - \int_{\mathbb{R}^m} \lim_{s \rightarrow 0} \frac{\mathbf{E}(\tilde{\psi}_l(z_s) | z_0 = z) - \tilde{\psi}_l(z)}{s} \tilde{\psi}_{l'}(z) d\nu(z) \\
&= - \lim_{s \rightarrow 0} \int_{\mathbb{R}^m} \frac{\mathbf{E}(\tilde{\psi}_l(z_s) | z_0 = z) - \tilde{\psi}_l(z)}{s} \tilde{\psi}_{l'}(z) d\nu(z) \\
&= - \lim_{s \rightarrow 0} \mathbf{E} \left[\frac{\tilde{\psi}_l(z_s) - \tilde{\psi}_l(z_0)}{s} \tilde{\psi}_{l'}(z_0) \mid z_0 \sim \nu \right] \\
&= - \lim_{s \rightarrow 0} \lim_{T \rightarrow +\infty} \frac{1}{T} \int_0^T \frac{\tilde{\psi}_l(z_{t+s}) - \tilde{\psi}_l(z_t)}{s} \tilde{\psi}_{l'}(z_t) dt \\
&= - \lim_{s \rightarrow 0} \lim_{T \rightarrow +\infty} \frac{1}{T} \int_0^T \frac{\tilde{\psi}_{l'}(z_{t+s}) - \tilde{\psi}_{l'}(z_t)}{s} \tilde{\psi}_l(z_t) dt.
\end{aligned} \tag{5.4}$$

In the above, \mathbf{E} denotes the mathematical expectation with respect to the effective dynamics z_s , and the last equality follows from the symmetry of matrix C .

To compute $C_{ll'}$ numerically, we further introduce a parameter $\tau \ll 1$, and approximate (5.4) by

$$\begin{aligned}
C_{ll'} &\approx - \frac{1}{2(M - M_0)} \left[\sum_{i=M_0+1}^M \frac{\tilde{\psi}_l(z_{i\Delta t + \tau}) - \tilde{\psi}_l(z_{i\Delta t})}{\tau} \tilde{\psi}_{l'}(z_{i\Delta t}) \right. \\
&\quad \left. + \sum_{i=M_0+1}^M \frac{\tilde{\psi}_{l'}(z_{i\Delta t + \tau}) - \tilde{\psi}_{l'}(z_{i\Delta t})}{\tau} \tilde{\psi}_l(z_{i\Delta t}) \right] \\
&= - \frac{1}{2(M - M_0)} \sum_{i=M_0+1}^M \frac{\tilde{\psi}_l(z_{i\Delta t + \tau})\tilde{\psi}_{l'}(z_{i\Delta t}) + \tilde{\psi}_l(z_{i\Delta t})\tilde{\psi}_{l'}(z_{i\Delta t + \tau}) - 2\tilde{\psi}_l(z_{i\Delta t})\tilde{\psi}_{l'}(z_{i\Delta t})}{\tau}.
\end{aligned} \tag{5.5}$$

Formulas (5.3) and (5.5) can be used to estimate the coefficient matrices C, S , provided that we can obtain a long trajectory of the effective dynamics (3.18).

REMARK 4. *From the discussions in Section 2 we know that the eigenvalues of the transfer operator \mathcal{T}_τ and that of the operator $-\mathcal{L}$ satisfy the relation $\mu_i = e^{-\lambda_i \tau}$, $i \geq 0$. When the lag time τ is small, the approximation $\mu_i \approx 1 - \lambda_i \tau$ holds for the leading eigenvalues since λ_i is small. In fact, estimating the matrix C using the last expression in (5.5), we will have $C = \frac{S - \bar{C}}{\tau}$, where the matrix \bar{C} is given by*

$$\bar{C}_{ll'} = \frac{1}{2(M - M_0)} \sum_{i=M_0+1}^M \left[\tilde{\psi}_l(z_{i\Delta t + \tau})\tilde{\psi}_{l'}(z_{i\Delta t}) + \tilde{\psi}_l(z_{i\Delta t})\tilde{\psi}_{l'}(z_{i\Delta t + \tau}) \right]. \tag{5.6}$$

It is easy to observe that the eigenvalue estimates resulting from problem (4.15) are related to those of the problem $\bar{C}X = \mu SX$ by $\mu = 1 - \lambda\tau$. Note that (5.6) is very similar to the estimator derived in [26] except for the fact that here we use trajectories of the effective dynamics instead of the original dynamics. To summarize, when the lag time τ is small, the above discussion implies that after solving the problem (4.15) we can approximate the leading eigenvalues of the transfer operator by $\mu_i = 1 - \lambda_i \tau$.

5.2. Algorithms for simulating the effective dynamics. In order to utilize the above results we have to be able to efficiently compute (long) realizations of the effective dynamics (3.18). In this subsection, we discuss two numerical algorithms for realizing this.

First algorithm. The first one is based on the following formula for the coefficients \tilde{b}, \tilde{a} given in (3.19)

$$\begin{aligned}\tilde{b}_l(z) &= \lim_{s \rightarrow 0^+} \mathbf{E} \left(\frac{\xi_l(x_s) - z_l}{s} \mid x_0 \sim \mu_z \right), \quad 1 \leq l \leq m, \\ \tilde{a}_{ll'}(z) &= \frac{\beta}{2} \lim_{s \rightarrow 0^+} \mathbf{E} \left(\frac{(\xi_l(x_s) - z_l)(\xi_{l'}(x_s) - z_{l'})}{s} \mid x_0 \sim \mu_z \right), \quad 1 \leq l, l' \leq m.\end{aligned}\tag{5.7}$$

where x_s is a realization of the original diffusive dynamics (2.1) and μ_z is restriction of invariant measure to the submanifold $\xi^{-1}(z) = \{x \in \mathbb{R}^n \mid \xi(x) = z\}$. We refer the readers to [35] for more details.

In order to utilize this for simulation, fix two parameters $0 < \Delta s \ll \Delta t$ and proceed as follows:

Algorithm 1.

1. At step $k \geq 0$, starting from $x_0 \sim \mu_z$, generate N trajectories $x_{\Delta s}^{(i)}$ of length Δs of the (unconstrained) full dynamics x_s by discretizing (2.1). Compute the coefficients \tilde{b}, \tilde{a} by

$$\begin{aligned}\tilde{b}_l &= \frac{1}{N} \sum_{i=1}^N \frac{\xi_l(x_{\Delta s}^{(i)}) - z_{k\Delta t, l}}{\Delta s}, \\ \tilde{a}_{ll'} &= \frac{\beta}{2} \left[\frac{1}{N} \sum_{i=1}^N \frac{(\xi_l(x_{\Delta s}^{(i)}) - z_{k\Delta t, l})(\xi_{l'}(x_{\Delta s}^{(i)}) - z_{k\Delta t, l'})}{\Delta s} - \tilde{b}_l \tilde{b}_{l'} \Delta s \right],\end{aligned}\tag{5.8}$$

where $1 \leq l, l' \leq m$.

2. Compute $\tilde{\sigma}$ from $\tilde{a} = \tilde{\sigma} \tilde{\sigma}^T$ by matrix decomposition. Update $z_{(k+1)\Delta t}$ by

$$z_{(k+1)\Delta t, l} = z_{k\Delta t, l} + \tilde{b}_l \Delta t + \sqrt{\frac{2\Delta t}{\beta}} \sum_{i=1}^m \tilde{\sigma}_{li} \eta_i^{(k)}, \quad 1 \leq l \leq m,\tag{5.9}$$

where $\eta_i^{(k)}$ are independent standard Gaussian variables, $1 \leq i \leq m$.

In the above, $z_{k\Delta t, l}$ denotes the l th components of $z_{k\Delta t} \in \mathbb{R}^m$. The initial states x_0 are sampled from the probability measure μ_z ; this can be achieved by using the numerical schemes proposed in [3, 4], which simulate the original dynamics (2.1) and then project the state onto the submanifold $\xi^{-1}(z)$.

Second algorithm. The second algorithm is inspired by the TAMM method as of [18]. In the following we provide a slightly different argument which motivates the method. The main idea is to consider the extended dynamics

$$\begin{aligned}dx_{s,i} &= -\frac{\partial V}{\partial x_i}(x_s) ds - \kappa \sum_{j=1}^m (\xi_j(x_s) - z_{s,j}) \frac{\partial \xi_j}{\partial x_i}(x_s) ds + \sqrt{2\beta^{-1}} dw_{s,i}, \quad 1 \leq i \leq n, \\ dz_{s,l} &= \kappa \sum_{k=1}^m (\xi_k(x_s) - z_{s,k}) \sum_{j=1}^n \frac{\partial \xi_k}{\partial x_j}(x_s) \frac{\partial \xi_l}{\partial x_j}(x_s) ds + \sqrt{2\beta^{-1}} \sum_{j=1}^n \frac{\partial \xi_l}{\partial x_j}(x_s) d\bar{w}_{s,j}, \quad 1 \leq l \leq m,\end{aligned}\tag{5.10}$$

where κ is a large constant, w_s, \bar{w}_s are independent Brownian motions on \mathbb{R}^n , and $x_{s,i}$ denotes the i th component of state x_s (similar notations for z_s, w_s, \bar{w}_s). Note that the invariant measure of dynamics (5.10) has a probability density

$$\rho_\kappa(x, z) \propto e^{-\beta(V(x) + \frac{\kappa}{2}|\xi(x) - z|^2)}, \quad (x, z) \in \mathbb{R}^{n+m}, \quad (5.11)$$

with respect to the Lebesgue measure on the extended space \mathbb{R}^{n+m} . If we choose $(x, z) \rightarrow z$ as the reaction coordinate function and derive the effective dynamics of (5.10) following [17, 35], we can obtain

$$dz_s = \tilde{b}^{(\kappa)}(z_s) ds + \sqrt{2\beta^{-1}\tilde{\sigma}^{(\kappa)}(z_s)} dw_s \quad (5.12)$$

where w_s is a Brownian motion on \mathbb{R}^m , and

$$\begin{aligned} \tilde{b}_l^{(\kappa)}(z) &= \kappa \int_{\mathbb{R}^n} \sum_{k=1}^m (\xi_k(x) - z_k) \sum_{i=1}^n \frac{\partial \xi_k}{\partial x_i}(x) \frac{\partial \xi_l}{\partial x_i}(x) \rho_\kappa(x, z) dx = \int_{\mathbb{R}^n} \mathcal{L} \xi_l(x) \rho_\kappa(x, z) dx \\ \tilde{a}_{ll'}^{(\kappa)} &= (\tilde{\sigma}^{(\kappa)}(\tilde{\sigma}^{(\kappa)})^T)_{ll'}(z) = \int_{\mathbb{R}^n} \sum_{i=1}^n \frac{\partial \xi_l}{\partial x_i}(x) \frac{\partial \xi_{l'}}{\partial x_i}(x) \rho_\kappa(x, z) dx, \end{aligned} \quad (5.13)$$

for $z \in \mathbb{R}^m$, $1 \leq l, l' \leq m$. Note that in (5.13), \mathcal{L} is the generator given in (2.2) and integration by parts has been used to derive the second equality for $\tilde{b}^{(\kappa)}$. It is not difficult to show that $\tilde{b}^{(\kappa)} \rightarrow \tilde{b}$ and $\tilde{a}^{(\kappa)} \rightarrow \tilde{a}$, when $\kappa \rightarrow +\infty$. Therefore (5.12) is an approximation of the effective dynamics (3.18) when $\kappa \gg 1$. For numerical simulations, we can express (5.13) as time averages

$$\begin{aligned} \tilde{b}_l^{(\kappa)}(z) &= \lim_{T \rightarrow \infty} \frac{\kappa}{T} \int_0^T \sum_{k=1}^m (\xi_k(x_s) - z_k) \sum_{i=1}^n \frac{\partial \xi_l}{\partial x_i}(x_s) \frac{\partial \xi_k}{\partial x_i}(x_s) ds, \quad 1 \leq l \leq m, \\ \tilde{a}_{ll'}^{(\kappa)}(z) &= \lim_{T \rightarrow \infty} \frac{1}{T} \int_0^T \sum_{i=1}^n \frac{\partial \xi_l}{\partial x_i}(x_s) \frac{\partial \xi_{l'}}{\partial x_i}(x_s) ds, \quad 1 \leq l, l' \leq m, \end{aligned} \quad (5.14)$$

where x_s satisfies SDE (5.10) with fixed $z_s = z$, i.e.,

$$dx_{s,i} = -\frac{\partial V}{\partial x_i}(x_s) ds - \kappa \sum_{j=1}^m (\xi_j(x_s) - z_j) \frac{\partial \xi_j}{\partial x_i}(x_s) ds + \sqrt{2\beta^{-1}} dw_{s,i}, \quad 1 \leq i \leq n, \quad (5.15)$$

The main steps of the algorithm can be summarized as follows:

Algorithm 2.

1. Denote $z = z_{k\Delta t}$ at step $k \geq 0$. Simulate dynamics (5.15) for M steps with time step size Δs . Compute the coefficients

$$\begin{aligned} \tilde{b}_l &= \frac{\kappa}{M - M_0} \sum_{j=M_0+1}^M \sum_{l'=1}^m (\xi_{l'}(x_{j\Delta s}) - z_{l'}) \sum_{i=1}^n \frac{\partial \xi_l}{\partial x_i}(x_{j\Delta s}) \frac{\partial \xi_{l'}}{\partial x_i}(x_{j\Delta s}), \quad 1 \leq l \leq m, \\ \tilde{a}_{ll'} &= \frac{1}{M - M_0} \sum_{j=M_0+1}^M \sum_{i=1}^n \frac{\partial \xi_l}{\partial x_i}(x_{j\Delta s}) \frac{\partial \xi_{l'}}{\partial x_i}(x_{j\Delta s}), \quad 1 \leq l, l' \leq m. \end{aligned} \quad (5.16)$$

2. Compute $\tilde{\sigma}$ from $\tilde{a} = \tilde{\sigma}\tilde{\sigma}^T$ by matrix decomposition. Update state $z_{(k+1)\Delta t}$ according to

$$z_{(k+1)\Delta t, l} = z_{k\Delta t, l} + \tilde{b}_l \Delta t + \sqrt{\frac{2\Delta t}{\beta}} \sum_{i=1}^m \tilde{\sigma}_{li} \eta_i^{(k)}, \quad 1 \leq l \leq m, \quad (5.17)$$

where $\eta_i^{(k)}$ are independent standard Gaussian variables, $1 \leq i \leq m$.

6. Illustrative example. In order to illustrate the analysis and the performance of the numerical methods presented in the previous sections, we study a simple two dimensional dynamics

$$\begin{aligned} dx_{s,1} &= -\frac{\partial V(x_s)}{\partial x_1} ds + \sqrt{2\beta^{-1}} dw_{s,1}, \\ dx_{s,2} &= -\frac{\partial V(x_s)}{\partial x_2} ds + \sqrt{2\beta^{-1}} dw_{s,2}, \end{aligned} \quad (6.1)$$

where $\beta > 0$, $x_s = (x_{s,1}, x_{s,2}) \in \mathbb{R}^2$ and $w_{s,1}, w_{s,2}$ are two independent one-dimensional Brownian motions.

The potential V in dynamics (6.1) is defined as

$$V(x) = V_1(\theta) + \frac{1}{\epsilon} V_2(r, \theta), \quad (6.2)$$

where $\epsilon > 0$,

$$V_1(\theta) = \begin{cases} [1 - \frac{9}{\pi^2} (\theta - \frac{\pi}{3})^2]^2 & \theta > \frac{\pi}{3}, \\ \frac{3}{5} - \frac{2}{5} \cos 3\theta & -\frac{\pi}{3} \leq \theta \leq \frac{\pi}{3}, \\ [1 - \frac{9}{\pi^2} (\theta + \frac{\pi}{3})^2]^2 & \theta < -\frac{\pi}{3}, \end{cases} \quad \text{and} \quad V_2(r, \theta) = \left(r^2 - 1 - \frac{1}{1 + 4r\theta^2} \right)^2,$$

and (r, θ) is the polar coordinate of state $x = (x_1, x_2)$ satisfying

$$\begin{aligned} x_1 &= r \cos \theta, & x_2 &= r \sin \theta, \\ \theta &\in [-\pi, \pi], & r &\geq 0. \end{aligned} \quad (6.3)$$

Under polar coordinate, it is easy to see that the potential V contains three local minima at $\theta = 0, \pm \frac{2\pi}{3}$ where the radius is determined by the relation $r^2 = 1 + \frac{1}{1 + 4r\theta^2}$. Furthermore, when parameter ϵ is small, one can expect that the dynamics (6.1) will be mainly confined in the neighbourhood of the curve defined by the relation $r^2 = 1 + \frac{1}{1 + 4r\theta^2}$, where the potential is relatively flat. Profiles of potential V_1 and V are displayed in Figure 6.1.

The main purpose of this numerical experiment is to demonstrate that the leading eigenvalues of operator $-\mathcal{L}$ corresponding to dynamics (6.1) can be approximated with the help of its effective dynamics, provided that the reaction coordinate function as well as the basis functions are chosen appropriately.

We choose parameters $\beta = 4.0$ and $\epsilon = 0.05$ in the following numerical experiment. In fact, for this two-dimensional problem, it is possible to directly solve the eigenvalue problem (2.5) by discretizing the operator \mathcal{L} . First of all, we note that the generator can be written as $\mathcal{L} = \frac{e^{\beta V}}{\beta} \nabla (e^{-\beta V} \nabla)$. Defining the operator \mathcal{D} such that $\mathcal{D}f = e^{-\frac{\beta}{2} V} f$ for function f , it is straightforward to see that the operator $-\mathcal{L}_{\mathcal{D}} = -\mathcal{D}\mathcal{L}\mathcal{D}^{-1}$ has the same eigenvalues λ_i as $-\mathcal{L}$ and the corresponding eigenfunctions are given by $\varphi_i^{\mathcal{D}} = \mathcal{D}\varphi_i = e^{-\frac{\beta}{2} V} \varphi_i$, where φ_i are the eigenfunctions of $-\mathcal{L}$. Furthermore, $\mathcal{L}_{\mathcal{D}}$ is a self-adjoint operator under the standard L^2 inner product. Instead of considering $-\mathcal{L}$, we will work with $-\mathcal{L}_{\mathcal{D}}$ and solve the eigenvalue problem $-\mathcal{L}_{\mathcal{D}}f = \lambda f$ because the discretized matrix will be symmetric and the corresponding eigenfunctions $\varphi_i^{\mathcal{D}}$ decay rapidly.

Taking into account the profile of potential V in Figure 6.1(b), we truncate the whole space \mathbb{R}^2 into a finite domain $[-2, 2] \times [-2, 2]$, which is then discretized using a 500×500 uniform

mesh, leading to the cell resolution $\Delta x_1 = \Delta x_2 = \frac{4}{500} = 0.008$. For $1 \leq i, j \leq 500$, let $f_{i,j}, V_{i,j}$ denote the values of the functions f, V evaluated at state $(-2.0 + (i - \frac{1}{2})\Delta x_1, -2.0 + (j - \frac{1}{2})\Delta x_2)$, respectively. Other notations such as $V_{i \pm \frac{1}{2}, j}$ are defined in a similar way. Approximating $-\mathcal{L}_{\mathcal{D}}f = -\frac{1}{\beta}e^{\frac{\beta}{2}V}\nabla(e^{-\beta V}\nabla(e^{\frac{\beta}{2}V}f))$ by centered finite difference scheme, we obtain

$$\begin{aligned}
-(\mathcal{L}_{\mathcal{D}}f)_{i,j} \approx & \frac{e^{\frac{\beta}{2}V_{i,j}}}{\beta} \left[\frac{e^{-\beta V_{i-\frac{1}{2},j}} e^{\frac{\beta}{2}V_{i,j}} f_{i,j} - e^{\frac{\beta}{2}V_{i-1,j}} f_{i-1,j}}{\Delta x_1} - \frac{e^{-\beta V_{i+\frac{1}{2},j}} e^{\frac{\beta}{2}V_{i+1,j}} f_{i+1,j} - e^{\frac{\beta}{2}V_{i,j}} f_{i,j}}{\Delta x_1} \right. \\
& \left. + \frac{e^{-\beta V_{i,j-\frac{1}{2}}} e^{\frac{\beta}{2}V_{i,j}} f_{i,j} - e^{\frac{\beta}{2}V_{i,j-1}} f_{i,j-1}}{\Delta x_2} - \frac{e^{-\beta V_{i,j+\frac{1}{2}}} e^{\frac{\beta}{2}V_{i,j+1}} f_{i,j+1} - e^{\frac{\beta}{2}V_{i,j}} f_{i,j}}{\Delta x_2} \right], \tag{6.4}
\end{aligned}$$

for $1 < i, j < 500$. For boundary cells, Neumann condition is applied when the neighboring cells are lying outside of the truncated domain. From (6.4), it can be observed that the resulting discretization matrix is both symmetric and sparse. Solving the eigenvalues of this matrix (of order 250000) using Krylov-Schur method through numerical package SLEPc [10], we obtain the first 4 eigenvalues

$$\lambda_0 = 0.000, \quad \lambda_1 = 0.010, \quad \lambda_2 = 0.044, \quad \lambda_3 = 1.458, \tag{6.5}$$

with relative residual errors smaller than 1.1×10^{-6} . The corresponding eigenvectors are shown in Figure 6.2.

With the above reference result at hand, we continue to study the approximation quality of the effective dynamics with respect to the leading eigenvalues. For this purpose, we choose the reaction coordinate function as $\xi(x) = \theta(x) \in [-\pi, \pi]$, i.e., our reaction coordinate is the angle of the polar coordinate representation. Direct calculation shows that the coefficients $\tilde{b}, \tilde{\sigma}$ in (3.19) reduces to

$$\tilde{b}(z) = \mathbf{P}(-\nabla V \cdot \nabla \theta)(z), \quad \tilde{a}(z) = (\tilde{\sigma} \tilde{\sigma}^T)(z) = \mathbf{P}\left(\frac{1}{r^2}\right)(z), \quad z \in [-\pi, \pi]. \tag{6.6}$$

Discretizing the interval $[-\pi, \pi]$ into 1000 subintervals and applying the projection scheme proposed in [4] for each fixed $z = -\pi + \frac{2\pi j}{1000}$, $0 \leq j \leq 1000$, we can compute the coefficients of the effective dynamics; the resulting profiles are shown in Figure 6.3(a) and Figure 6.3(b). After these preparations, we can generate trajectories of the effective dynamics by simulating the SDE (3.18) using standard time stepping schemes. As shown in Figure 6.3(c), the effective dynamics spends long time around values $-\frac{2\pi}{3}, 0$ and $\frac{2\pi}{3}$, which is accordance with the behavior of dynamics (6.1) as well as with the profile of potential V in Figure 6.1(b). Since the effective dynamics is one-dimensional, we can also discretize its infinitesimal generator $\tilde{\mathcal{L}}$ in (3.20) and compute the eigenvalues of $-\tilde{\mathcal{L}}$ which gives

$$\tilde{\lambda}_0 = 0.000, \quad \tilde{\lambda}_1 = 0.012, \quad \tilde{\lambda}_2 = 0.044, \quad \tilde{\lambda}_3 = 2.068.$$

Comparing to (6.5), we conclude that the eigenvalues $\lambda_0, \lambda_1, \lambda_2$ of the original dynamics (6.1) are quite well approximated by those of the effective dynamics.

As the final step of our experiment, we test the trajectory-based method proposed in Subsection 5.1. First of all, we define basis functions $\tilde{\psi}_1(z) \equiv 1.0$ and $\tilde{\psi}_i(z) = \exp(-\frac{(z-c_i)^2}{2\gamma_i^2})$, $2 \leq i \leq 7$, where

$$c_i = \left\{ -\frac{2\pi}{3}, -\frac{2\pi}{3}, 0, 0, \frac{2\pi}{3}, \frac{2\pi}{3} \right\}, \quad r_i = \left\{ 0.4, 0.7, 0.4, 0.7, 0.4, 0.7 \right\}.$$

That is, we have located two Gaussian-like basis functions with different radius (0.4 and 0.7) at each of the three local minima $\theta = 0, \pm \frac{2\pi}{3}$. The matrices S and C are then estimated according to (5.3) and (5.5) by generating 4 long trajectories of the effective dynamics with time step size $\Delta t = 5 \times 10^{-4}$, and parameters $\tau = 20\Delta t$, $M_0 = 1000$, $M = 2 \times 10^7$ are used for each trajectories. Solving the generalized matrix eigenvalue problem $CX = \lambda SX$, we obtain the leading eigenvalues

$$\tilde{\lambda}_0 = 0.000, \quad \tilde{\lambda}_1 = 0.013, \quad \tilde{\lambda}_2 = 0.045, \quad \tilde{\lambda}_3 = 3.776.$$

As before, we conclude that the eigenvalues $\lambda_0, \lambda_1, \lambda_2$ of the original dynamics are relatively well approximated.

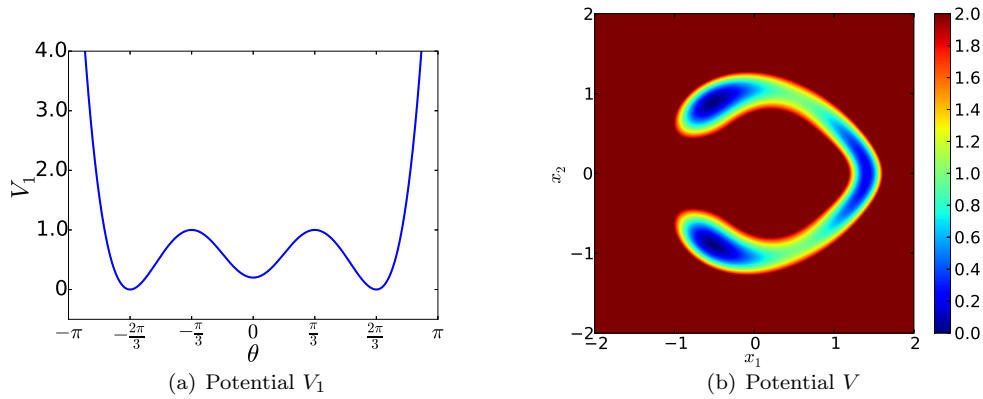


Figure 6.1: (a) Function V_1 as a function of angle θ . (b) Potential V defined in (6.2) with parameter $\epsilon = 0.05$.

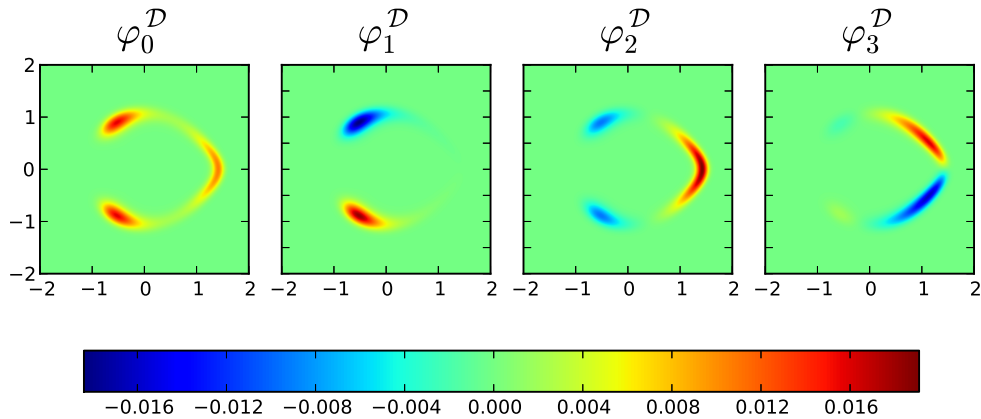


Figure 6.2: Eigenfunctions $\varphi_i^{\mathcal{D}}$ of operator $-\mathcal{L}_{\mathcal{D}}$ corresponding to the first 4 eigenvalues in (6.5).

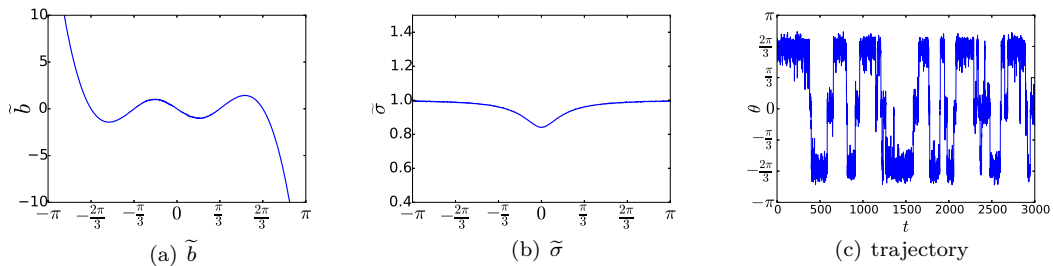


Figure 6.3: (a)(b) Coefficients \tilde{b} and $\tilde{\sigma}$ as given in (6.6). For each $z = -\pi + \frac{2\pi j}{1000}$, $0 \leq j \leq 1000$, the coefficients $\tilde{b}(z)$, $\tilde{\sigma}(z)$ are estimated by generating a trajectory of the constrained version of dynamics (6.1) using the projection scheme proposed in [4] with the time step size 2×10^{-5} , and 3×10^6 steps are simulated. (c) A typical sample trajectory of the effective dynamics for dynamics (6.1) with reaction coordinate function $\xi(x) = \theta(x)$.

7. Conclusions. In this work we have studied the approximation of eigenvalues and eigenfunctions of the infinitesimal generator associated with the longest relaxation processes of diffusive processes in energy landscapes. Following the previous studies on transfer operators, we consider the Galerkin discretization method, the variational approach and the effective dynamics given by a low-dimensional reaction coordinate for solving the eigenvalue problem in application to the generator. It turns out that (1) there are rather similar results for the approximation error of the three methods, and that (2) the first two methods lead to the same generalized matrix eigenvalue problem while the third can be used for efficient estimation of the associated coefficient matrices.

While we have assumed that the dynamics is diffusive, we emphasize that the analysis in the current work can be directly applied to more general reversible processes, see [35] for details. For non-reversible dynamics, as, for example, for Langevin dynamics, it is not immediately clear how the results in the current work can be applied. However, the approach in [33], Section 5.3, shows that the extended reversibility of Langevin dynamics may well allow for a generalization of our results. Furthermore, applications of the algorithms presented to more complicated systems are not considered in this work; they are topics of ongoing research.

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REFERENCES

- [1] N. B. BECKER, R. J. ALLEN, AND P. R. TEN WOLDE, *Non-stationary forward flux sampling*, J. Chem. Phys., 136 (2012), p. 174118.
- [2] G. R. BOWMAN, V. S. PANDE, AND F. NOÉ, eds., *An Introduction to Markov State Models and Their Application to Long Timescale Molecular Simulation*, vol. 797 of Advances in Experimental Medicine and Biology, Springer, 2014.
- [3] G. CICCOTTI, R. KAPRAL, AND E. VANDEN-EIJNDEN, *Blue moon sampling, vectorial reaction coordinates, and unbiased constrained dynamics*, ChemPhysChem, 6 (2005), pp. 1809–1814.
- [4] G. CICCOTTI, T. LELIÈVRE, AND E. VANDEN-EIJNDEN, *Projection of diffusions on submanifolds: Application to mean force computation*, Commun. Pur. Appl. Math., 61 (2008), pp. 371–408.

- [5] E. DARVE, D. RODRIGUEZ-GÓMEZ, AND A. POHORILLE, *Adaptive biasing force method for scalar and vector free energy calculations*, J. Chem. Phys., 128 (2008), p. 144120.
- [6] N. DJURDJEVIC, M. SARICH, AND CH. SCHÜTTE, *Estimating the eigenvalue error of Markov state models*, Multiscale Model. Simul., 10 (2012), pp. 61 – 81.
- [7] A. K. FARADJIAN AND R. ELBER, *Computing time scales from reaction coordinates by milestoning*, J. Chem. Phys., 120 (2004), pp. 10880–10889.
- [8] G. FROYLAND, G. A. GOTTWALD, AND A. HAMMERLINDL, *A computational method to extract macroscopic variables and their dynamics in multiscale systems*, SIAM J. Appl. Dyn. Syst., 13 (2014), pp. 1816–1846.
- [9] I. GYÖNGY, *Mimicking the one-dimensional marginal distributions of processes having an Ito differential*, Probab. Th. Rel. Fields, 71 (1986), pp. 501–516.
- [10] V. HERNANDEZ, J. E. ROMAN, AND V. VIDAL, *SLEPC: A scalable and flexible toolkit for the solution of eigenvalue problems*, ACM Trans. Math. Software, 31 (2005), pp. 351–362.
- [11] I. G. KEVREKIDIS, C. W. GEAR, AND G. HUMMER, *Equation-free: The computer-aided analysis of complex multiscale systems*, AIChE J., 50 (2004), pp. 1346–1355.
- [12] I. G. KEVREKIDIS, C. W. GEAR, J. M. HYMAN, P. G. KEVREKIDIS, O. RUNBORG, AND C. THEODOROPOULOS, *Equation-free, coarse-grained multiscale computation: Enabling microscopical simulators to perform system-level analysis*, Commun. Math. Sci., 1 (2003), pp. 715–762.
- [13] I. G. KEVREKIDIS AND G. SAMAIEY, *Equation-free multiscale computation: Algorithms and applications*, Annu. Rev. Phys. Chem., 60 (2009), pp. 321–344.
- [14] S. KUMAR, J. M. ROSENBERG, D. BOUZIDA, R. H. SWENDSEN, AND P. A. KOLLMAN, *THE weighted histogram analysis method for free-energy calculations on biomolecules. I. The method*, J. Comput. Chem., 13 (1992), pp. 1011–1021.
- [15] A. LAIO AND F. L. GERVASIO, *Metadynamics: a method to simulate rare events and reconstruct the free energy in biophysics, chemistry and material science*, Rep. Prog. Phys., 71 (2008), p. 126601.
- [16] A. LAIO AND M. PARRINELLO, *Escaping free-energy minima*, Proc. Natl. Acad. Sci. USA, 99 (2002), pp. 12562–12566.
- [17] F. LEGOLL AND T. LELIÈVRE, *Effective dynamics using conditional expectations*, Nonlinearity, 23 (2010), pp. 2131–2163.
- [18] L. MARAGLIANO AND E. VANDEN-ELJNDEN, *A temperature accelerated method for sampling free energy and determining reaction pathways in rare events simulations*, Chem. Phys. Lett., 426 (2006), pp. 168 – 175.
- [19] J. C. MATTINGLY, A. M. STUART, AND D. J. HIGHAM, *Ergodicity for SDEs and approximations: locally Lipschitz vector fields and degenerate noise*, Stoch. Proc. Appl., 101 (2002), pp. 185 – 232.
- [20] D. MORONI, T. VAN ERP, AND P. BOLHUIS, *Investigating rare events by transition interface sampling*, Physica A, 340 (2004), pp. 395–401.
- [21] F. NOÉ AND F. NÜSKE, *A variational approach to modeling slow processes in stochastic dynamical systems*, Multiscale Model. Simul., 11 (2013), pp. 635–655.
- [22] F. NOÉ, CH. SCHÜTTE, E. VANDEN-ELJNDEN, L. REICH, AND T. R. WEIKL, *Constructing the full ensemble of folding pathways from short off-equilibrium simulations*, Proc. Natl. Acad. Sci. USA, 106 (2009), pp. 19011–19016.
- [23] F. NÜSKE, B. G. KELLER, G. PÉREZ-HERNÁNDEZ, A. MEY, AND F. NOÉ, *Variational approach to molecular kinetics*, J. Chem. Theory Comput., 10 (2014), pp. 1739–1752.
- [24] F. NÜSKE, R. SCHNEIDER, F. VITALINI, AND F. NOÉ, *Variational tensor approach for approximating the rare-event kinetics of macromolecular systems*, J. Chem. Phys., 144 (2016).
- [25] V. S. PANDE, K. BEAUCHAMP, AND G. R. BOWMAN, *Everything you wanted to know about Markov state models but were afraid to ask*, Methods, 52 (2010), pp. 99–105.
- [26] G. PÉREZ-HERNÁNDEZ, F. PAUL, T. GIORGINO, G. DE FABRITHS, AND F. NOÉ, *Identification of slow molecular order parameters for Markov model construction*, J. Chem. Phys., 139 (2013), p. 015102.
- [27] J.-H. PRINZ, H. WU, M. SARICH, B. KELLER, M. SENNE, M. HELD, J. D. CHODERA, CH. SCHÜTTE, AND F. NOÉ, *Markov models of molecular kinetics: Generation and validation*, J. Chem. Phys., 134 (2011), p. 174105.
- [28] M. SARICH, F. NOÉ, AND CH. SCHÜTTE, *On the approximation quality of Markov state models*, Multiscale Model. Simul., 8 (2010), pp. 1154–1177.
- [29] M. SARICH AND CH. SCHÜTTE, *Approximating selected non-dominant timescales by Markov state models*, Comm. Math. Sci., 10 (2012), pp. 1001 – 1013.

- [30] CH. SCHÜTTE, A. FISCHER, W. HUISINGA, AND P. DEUFLHARD, *A direct approach to conformational dynamics based on hybrid Monte Carlo*, J. Comput. Phys., 151 (1999), pp. 146–168.
- [31] CH. SCHÜTTE, W. HUISINGA, AND P. DEUFLHARD, *Transfer operator approach to conformational dynamics in biomolecular systems*, in Ergodic Theory, Analysis, and Efficient Simulation of Dynamical Systems, B. Fiedler, ed., Springer Berlin Heidelberg, 2001, pp. 191–223.
- [32] CH. SCHÜTTE, F. NOÉ, J. LU, M. SARICH, AND E. VANDEN-EIJNDEN, *Markov state models based on milestone*, J. Chem. Phys., 134 (2011).
- [33] CH. SCHÜTTE AND M. SARICH, *Metastability and Markov State Models in Molecular Dynamics: Modeling, Analysis, Algorithmic Approaches*, Courant Lecture Notes, American Mathematical Society/Courant Institute of Mathematical Science, 2014.
- [34] G. M. TORRIE AND J. P. VALLEAU, *Nonphysical sampling distributions in Monte Carlo free-energy estimation: Umbrella sampling*, J. Comput. Phys., 23 (1977), pp. 187–199.
- [35] W. ZHANG, C. HARTMANN, AND CH. SCHUTTE, *Effective dynamics along given reaction coordinates, and reaction rate theory*, Faraday Discuss., 195 (2016), pp. 365–394.