



Quantile Trend Regression and Its Application to Central England Temperature

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Abstract: The identification and estimation of trends in hydroclimatic time series remains an important task in applied climate research. The statistical challenge arises from the inherent nonlinearity, complex dependence structure, heterogeneity and resulting non-standard distributions of the underlying time series. Quantile regressions are considered an important modeling technique for such analyses because of their rich interpretation and their broad insensitivity to extreme distributions. This paper provides an asymptotic justification of quantile trend regression in terms of unknown heterogeneity and dependence structure and the corresponding interpretation. An empirical application sheds light on the relevance of quantile regression modeling for analyzing monthly Central England temperature anomalies and illustrates their various heterogenous trends. Our results suggest the presence of heterogeneities across the considered seasonal cycle and an increase in the relative frequency of observing unusually high temperatures.

Keywords: temperature; trend modeling; seasonality; heterogeneity; quantile regression

JEL Classification: C02; C14; C18; C22; Q54

1. Introduction

The current phase of rapidly accelerating impacts of global warming reveals many immediate effects on direct and indirect threats to global health and clearly illustrates the role of human activities (e.g., [1,2]). Recently, the International Panel on Climate Change [3] reported unprecedented increases in global surface temperature over the last four decades since the start of geo-referenced systematic climatological records in 1850with larger increases for land- than sea-surface temperatures. Substantial increases in mean temperatures across the world and heterogeneity in temperature increases across months are noted in Vogelsang and Franses [4], who emphasize that especially winters are warming in the northern hemisphere. King et al. [1] document substantial warming in the Central England Temperature (CET) time series for 2014 based on simulated and observed datasets and provide evidence for a human-induced increase in the relative frequency of observing record high temperatures. King [5] extends their previous analysis beyond the univariate approach and provides a proxy for the local response to global warming by modeling local temperature as a nonlinear parametric function of global temperature. Rivas and Gonzalo [6] provide a definition of global warming and trending time series and present evidence for warming based on the CET and global temperature time series.

The properties of climatologic time series have been frequently studied in the literature and the related findings fuel an ongoing debate about the analysis and attribution of trends (see, e.g., [7,8], for recent summaries and discussions). When investigating the properties of climatologic time series, seasonality is typically filtered in a first step, as it may frustrate diagnostic tests and complicate inference (e.g., [9], in the context of long



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Copyright © 2022 by the authors. Licensee MDPI, Basel, Switzerland. This article is an open access article distributed under the terms and conditions of the Creative Commons Attribution (CC BY) license (https:// creativecommons.org/licenses/by/ 4.0/). range dependence and structural breaks). After filtering the seasonal annual cycle, Gao and Franzke [10] test temperature series for the presence of long range dependence. Proietti and Hillebrand [11] conduct a series of univariate (and multivariate) stationarity (and cointegration) checks against various alternatives and find evidence of nonstationarity. Their findings are in line with He et al. [12]. Among others, Fomby and Vogelsang [13], Barbosa [14], and Franzke [15] discuss previous studies suggesting that temperature series are mean stationary and exhibit long range dependence. While such dynamics do not generate trends in the long run (e.g., thousand years or longer), series such as the CET ranging over a few hundred years may exhibit stochastic trends. Whether such trends are deterministic or stochastic and due to short run or long run dependence has been investigated intensively (e.g., [4,11,16]).

Koenker and Schorfheide [17] investigate curvature in temperature trends in greater detail. They discuss an upward trend in global Hansen-Lebedeff temperatures [18] between 1880 and 1940 and after 1965, and identify a trend break in a period of decreasing temperatures between 1940 and 1965. An important challenge in the analysis of hydrologic and climatologic processes is the delicate interplay of nonlinearities, temporal dependence, bimodality, and heterogeneity (e.g., [5,17,19–23]). Rial et al. [24] provide potential explanations of the mechanisms underlying the climate system such as the superposition of shift and trends in long-term and short-term dynamics, ocean atmosphere interactions, and feedback effects from the carbon-cycle. Koenker and Schorfheide [17] suggest to broaden the perspective of analysis beyond the common practice of modeling the deviations around deterministic trends for the conditional mean temperature using ARIMA (autoregressive integrated moving average) processes. They propose the use of conditional quantiles as "serial correlation may be due to a variety of mechanisms that could plausibly effect year-to-year changes" in temperature and climatic conditions in general (e.g., [25]). As a consequence, in recent contributions to statistical modeling of temperature series, the method of quantiles introduced by Koenker and Bassett [26] plays a major role to define and identify anomalies (e.g., [27,28]).

The presence of seasonal trends has been discussed for climatologic time series by Colman [29], who argues that sea-surface temperature anomalies may drive seasonal trends in temperature series, which entails important consequences for ecologic policy (e.g., [30]). Recently, He et al. [12] emphasized that seasonal trends are difficult to model and forecast due to time-varying patterns in their components. In what follows, we avoid the perils of applying a particular approach to filter seasonality and use the assumption of cyclostationarity (e.g., [31]) in a quantile trend regression (QTR) framework. Cyclostationarity is based on the literature on periodically autocorrelated processes (e.g., [32]) and assumes that trend patterns vary cyclically over time. As a consequence of this, temperature can be understood as a stationary process fluctuating around a deterministic trend component. Thus, by splitting up the original series into *S* non-seasonal subseries, where *S* corresponds to the length of the seasonal cycle, the estimating equations in our QTR framework can then be decomposed into two parts which are allowed to vary over the temperature distribution: the conditional temperature quantile estimated by a trend polynomial and a stationary colored noise process. We provide a detailed derivation and discussion of the assumptions underlying the QTR model and show the asymptotic normality of the QTR estimator. Additionally, we analyze temperature time series which are known to be crucial factors in models explaining morbidity and mortality and which are frequently used as predictors for other entities in epidemiology, climatology, ecology and economics.

The paper is structured as follows. Section 2 introduces a sketch of a baseline mathematical model and from this derives the QTR model, its assumptions, and asymptotic properties. Section 3 provides an overview of the CET anomaly data, illustrates the relevance of quantile regression for analyzing anomalies, and summarizes and discusses the modeling results. Section 4 provides concluding remarks.

2. Methods

Figures 1 and 2 highlight stylized facts of the CET data from January 1850 to December 2020 analyzed in the empirical section of the paper. Figure 1 shows the CET data, Figure 2 illustrates the corresponding annual mean temperatures. A season plot for the CET data is given in Figure 3, which includes the monthly mean temperature for each month from 1850 until 1900 via the grey horizontal lines. The latter figure illustrates the presence of heterogeneity in the respective monthly levels of temperature and also in the month-specific trajectories of temperature. The figures illustrate the generic nonlinearity and nonstationarity often found in environmental and hydroclimatic processes due to seasonal cycles or trend patterns. The decomposition of time series into trend and seasonality is the starting point for our modeling framework. We thereby address the requirements identified by Proietti and Hillebrand [11] that the temperature exhibits month-specific variation, dependence and trend patterns, while a global trend affects all months in a potentially different way. We will address this problem by using an approach similar to the modeling of periodically correlated sequences (e.g., [32]).



Figure 1. CET data from 1850 until 2020.



Figure 2. Annual mean temperature in Central England from 1850 until 2020. Dashed grey line represents annual mean temperature during years 1850 until 1900.



Figure 3. CET data for each month from 1961 until 2020. From left to right, months January to December are displayed over the years. Grey lines indicate monthly mean temperature from 1850 until 1900. Different months separated by dashed vertical lines.

2.1. Baseline Model

Let $(\Omega, \mathcal{F}, \mathbb{P})$ be a complete probability space and, following Koenker and Schorfheide [17], consider temperature data $\{Y\}_{t\geq 0}$ to be an \mathcal{F} -measurable scalar random sequence. Further assume that this process has a single cyclical component due to the annual season $s = 0, \ldots, S - 1$. This assumption implies that potential nonstationarities at seasonal position s (i.e., for each of the S = 12 months in our application) can only be generated by trend components (see [12] for a discussion in autoregressive models). Our approach deviates from the concept of periodical correlation in the Hilbert space $\mathbb{L}_2(\Omega, \mathcal{F}, \mathbb{P})$ of random variables with finite second moment, as we do not require the existence of any moments of the random processes constituting the estimating equation. However, we will exploit that the quantile regression loss function ρ_{τ} defined below belongs to the class of absolutely continuous (on every interval) functions with right-hand-derivative ψ_{τ} , which is a function of bounded variation (e.g., [33], Section 2.3).

Employing a common notation for seasonal time series (e.g., [12,34]) we denote the time index *t* as t = s + Si where i = 0, 1, ... is a new (annual) time index and consider a realization $\{Y_t\}_{t=0}^n$ with sample size n + 1. We assume that the data are generated according to

$$Y_{s+Si} = h(\mathbf{x}_{s+Si,n}, \boldsymbol{\beta}_s) + U_{s+Si}, \tag{1}$$

where h(.) is an unknown function continuous in $x_{s+Si,n}$, which depends on unknown but fixed season-specific parameters $\beta_s \in \mathbb{R}^L$, the row vector $x_{t,n}$ contains known and deterministic covariates depending on n, and U_{s+Si} is an error process.

We assume that the set of control variables does not vary over the seasons. Then, following Yang and Tschernig [34], the task of variable selection and parameter estimation has to be carried out for each month and corresponding data set

$$\{x_{s+Sin}, Y_{s+Si}\}_{i=0}^{\lfloor n/S \rfloor}, \quad s = 0, \dots, S-1.$$
(2)

The Weierstrass approximation theorem states that we can approximate

$$h_S(\mathbf{x}_{s+Si}, \boldsymbol{\beta}) \approx \mathbf{x}_{s+Si} \boldsymbol{\beta}_s$$

In the semiparametric context of model (1), we assume a polynomial regression design

$$\mathbf{x}_{s+Si,n} = (1, x_{s+Si,n}, x_{s+Si,n}^2, \dots, x_{s+Si,n}^K).$$
(3)

The baseline example in our data analysis is to assume that month-specific temperature series are generated by linear deterministic trends $h(\mathbf{x}_{s+Si,n}, \boldsymbol{\beta}_s) = \mu_s + \beta_s(s+Si)/n$ and additive, stationary and weakly dependent errors U_{s+Si} .

Example 1. Our empirical analysis of the CET data investigates temperature anomalies over the last 60 years from 1961 to 2020. Thus, n = 720, S = 12, and the effective sample size is N = n/S = 60. For convenience we consider the index i = 1, ..., N. A simple special case for a month-specific trend equation is the linear trend equation

$$Y_i = x_i \beta_0 + U_i = \mu + \beta \cdot i/60 + U_i, \quad i = 1, \dots, 60.$$
(4)

Using the QTR model, we will estimate β_0 for each month s, $1 \le s \le 12$, and various quantiles τ , $1 < \tau < 1$. Least squares based diagnostic testing, estimation, and inference of temperature trend regressions (4) have been discussed among others by Fomby and Vogelsang [13], Fatichi et al. [16], Barbosa [14], and in detail by Mudelsee [23].

2.2. Quantile Trend Regression

Instead of a least squares based trend modeling of the month-specific annual average increase of temperature data over time, we employ quantile regression. The resilience of quantile regression to extreme values in the temperature distribution allows robust estimation of annual median increases but also other non-central aspects of the temperature processes. As a consequence, the QTR estimator introduced in this chapter allows to draw inferences on trend patterns and their variation over the temperature distribution. We derive the asymptotic normality of the QTR estimator under primitive assumptions. Our argumentation is similar to Pollard [35], Fitzenberger [36], and Koenker [37], but builds on the weaker assumptions and corresponding results of Oberhofer and Haupt [38,39,40].

2.2.1. Basic Structure of a Quantile Regression Model

Our goal is to study season-specific trend patterns and investigate their heterogeneity over the temperature distribution. For this purpose we formulate a QTR model for each seasonal position *s*. For the sake of simplicity we will omit the index *s* whenever the context is clear. Using annual data, we assume that the temperature processes are generated according to the QTR equation

$$Y_i = x_i \beta_0 + U_i, \quad i = 1, ..., N,$$
 (5)

where $x_i := x_{iN}$, $1 \le i \le N$, are known deterministic $(1 \times K + 1)$ design vectors, and the true parameter β_0 is an inner point of a compact subset of \mathbb{R}^{K+1} . The unknown parameter vector β is estimated by the minimizer

$$\hat{\boldsymbol{\beta}}_{N} = \arg\min_{\boldsymbol{\beta} \in \mathbb{R}^{K+1}} \sum_{i=1}^{N} \rho_{\mathcal{X}}(Y_{i} - \boldsymbol{x}_{i}\boldsymbol{\beta}), \tag{6}$$

where $0 < \tau < 1$ and the conditional quantile $Q_{\tau}(Y|x) = x\beta_0$ is the optimal predictor of Y under quantile regression loss [26]

$$\rho_{\tau}(w) := |\tau - \mathbb{1}(w \le 0)| \cdot |w|,$$
(7)

where $\mathbb{I}(.)$ is the usual indicator function. By $F_i(v)$ we denote the distribution of U_i , by $F_{i,j}(v, w)$ the joint distribution of (U_i, U_j) for $i \neq j$, and employ the usual normalization in quantile regression:

(A1)
$$F_i(0) = \mathbb{P}(U_i \le 0) = \tau, 0 < \tau < 1, \forall i.$$

Starting from these core components of a quantile regression model, we motivate further assumptions required to derive the asymptotic normality of the QTR estimator. Due to the different growth rates of the components in the polynomial design (3), a normalization has to be applied (see [41], Ch. 7.2, for a detailed discussion in the case of least squares regression). Define

$$z_{iN} := x_i D_N^{-1}, \tag{8}$$

where $D_N = \text{diag}(d_{0N}, d_{1N}, \dots, d_{KN})$ is a nonsingular $(K + 1 \times K + 1)$ diagonal matrix with positive elements d_{kN} for $0 \le k \le K$ and for large enough N. The normalization (8) has to ensure that the elements of z_{iN} are Cesàro-summable (e.g., [33]).

Example 2. For the cubic polynomial trend model $Y_i = \mu + \beta_1 x_i + \beta_2 x_i^2 + \beta_3 x_i^3 + U_i$ employed in *Franzke* [28], d_{0N} is of order 1, d_{1N} is of order N/2, d_{2N} is of order N²/3, and d_{3N} is of order N³/4.

(A2) The $(1 \times K + 1)$ vectors x_i are deterministic and known and for some real number M, $||z_{iN}|| \le M < \infty$, for $1 \le i \le N$, and all N.

For our asymptotic analysis we re-write the argument of (6) according to (5) as

$$Y_{i} - \mathbf{x}_{i}\boldsymbol{\beta} = \mathbf{x}_{i}\boldsymbol{\beta}_{0} - \mathbf{x}_{i}\boldsymbol{\beta} + U_{i}$$

$$= U_{i} - \mathbf{x}_{i}(\boldsymbol{\beta} - \boldsymbol{\beta}_{0})$$

$$= U_{i} - \frac{1}{\sqrt{N}}\mathbf{x}_{i}\boldsymbol{D}_{N}^{-1}\sqrt{N}\boldsymbol{D}_{N}(\boldsymbol{\beta} - \boldsymbol{\beta}_{0})$$
(9)

$$= U_i - \frac{1}{\sqrt{N}} z_{iN} \gamma \tag{10}$$

$$= U_i - h_i(\gamma), \qquad (11)$$

with local parameters $\gamma := \sqrt{N}D_N(\beta - \beta_0)$ and $h_i(\gamma) := N^{-1/2}z_{iN}\gamma = x_i(\beta - \beta_0)$. Following Huber [42], we can then write the QTR objective function as

$$Q_N(\boldsymbol{\gamma}) = \sum_{i=1}^N q_{iN}(\boldsymbol{\gamma}), \qquad (12)$$

with summands

$$q_{iN}(\boldsymbol{\gamma}) = \rho_{\boldsymbol{\gamma}}(\boldsymbol{U}_i - \boldsymbol{h}_i(\boldsymbol{\gamma})) - \rho_{\boldsymbol{\gamma}}(\boldsymbol{U}_i) = \begin{cases} -\vartheta h_i & \text{if } \boldsymbol{U}_i > \max(0, h_i), \\ \boldsymbol{U}_i - \vartheta h_i & \text{if } \boldsymbol{h}_i < \boldsymbol{U}_i \leq 0, \\ -\boldsymbol{U}_i + (1 - \vartheta)h_i & \text{if } 0 < \boldsymbol{U}_i \leq h_i, \\ (1 - \vartheta)h_i & \text{if } \boldsymbol{U}_i \leq \min(0, h_i). \end{cases}$$
(13)

Then, under suitable assumptions (e.g., [37], Ch. 4, for the assumption of independent but not identically distributed random variables) the limiting distribution of $\hat{\gamma}_N$ can be derived from that of $Q_N(\gamma)$. If $\hat{\gamma}_N$ is the minimand of the scalar random variable $Q_N(\gamma)$, then the QTR estimator $\hat{\beta}_N$ of β_0 is given by

$$\hat{\boldsymbol{\beta}}_N = \boldsymbol{\beta}_0 + \frac{1}{\sqrt{N}} \boldsymbol{z}_{iN} \hat{\boldsymbol{\gamma}}.$$
(14)

Remark 1. In least squares regression, asymptotic analysis relies on the twice continuously differentiable loss function and local uniform approximation of $Q_N(\gamma)$. As this approach is not feasible in quantile regression, we have to approximate $Q_N(\gamma)$ by a sufficiently smooth function $\overline{Q}_N(\gamma)$. Due to its inherent smoothing, the expectation $\mathbb{E}[Q_N(\gamma)]$ is a suitable choice, whenever a uniform LLN can be established such that

$$\operatorname{plim}_{N} \sup_{\gamma} \left| \frac{1}{N} Q_{N}(\gamma) - \mathbb{E} \left[\frac{1}{N} Q_{N}(\gamma) \right] \right| = 0, \tag{15}$$

for γ in a compact subset of \mathbb{R}^{K+1} . Further elaboration is beyond the scope of this paper (see, e.g., [40], Lemma 2C), but we will get back to (15) in Remark 2 below and then analyze the asymptotic behaviour of $\mathbb{E}[Q_N(\gamma)]$ in Lemma 2.

Knight [43] suggests a useful decomposition of $Q_N(\gamma)$, which allows to study its asymptotic behavior by examining that of its partial sums separately. Knight's identity is given by

$$Q_N(\gamma) = B_N(\gamma) + \zeta_N(\gamma) = \sum_{i=1}^N b_{iN}(\gamma) + \sum_{i=1}^N \zeta_{iN}(\gamma),$$
(16)

where the second summand in (16) is defined as

$$\zeta_N(\boldsymbol{\gamma}) := -\sum_{i=1}^N h_i(\boldsymbol{\gamma}) \psi_{\boldsymbol{\gamma}}(\boldsymbol{U}_i), \tag{17}$$

where

$$\psi_{\tau}(w) := \tau - \mathbb{1}(w \le 0)$$
 (18)

is the right-hand-derivative of (7), and

$$\zeta_{iN}(\boldsymbol{\gamma}) = -h_i(\boldsymbol{\gamma})\psi_{\boldsymbol{\gamma}}(\boldsymbol{U}_i) = -\frac{1}{\sqrt{N}} \boldsymbol{z}_{iN}\boldsymbol{\gamma}[\boldsymbol{\tau} - \mathbb{I}(\boldsymbol{U}_i \leq 0)]. \tag{19}$$

The first summand in (16) is defined as

$$b_{iN}(\gamma) = |h_i - U_i| [\mathbb{1}(0 < U_i \le h_i) + \mathbb{1}(h_i < U_i \le 0)] \\ = \begin{cases} 0 & \text{if } U_i > \max(0, h_i), \\ U_i - h_i & \text{if } h_i < U_i \le 0, \\ -U_i + h_i & \text{if } 0 < U_i \le h_i, \\ 0 & \text{if } U_i \le \min(0, h_i). \end{cases}$$
(20)

The definition of summands $q_{iN}(\gamma)$ implies (e.g., [44], Lemma 3.1)

$$|q_{iN}(\boldsymbol{\gamma})| \le |h_i(\boldsymbol{\gamma})|. \tag{21}$$

As a consequence, every moment of $q_{iN}(\gamma)$ exists for finite $h_i(\gamma)$. Then, due to assumption (A1), $\mathbb{E}[\psi_{\gamma}(U_i)] = 0$, and it holds that

$$\mathbb{E}[\zeta_{iN}(\gamma)] = 0, \tag{22}$$

and

$$\mathbb{E}[q_{iN}(\gamma)] = \mathbb{E}[b_{iN}(\gamma)] = \begin{cases} \int_0^{h_i} (h_i - v) f_i(v) dv, & \text{if } h_i > 0, \\ 0, & \text{if } h_i = 0, \\ \int_{h_i}^0 (v - h_i) f_i(v) dv, & \text{if } h_i < 0. \end{cases}$$
(23)

2.2.2. Modeling the QTR Dependence Structure

In the case of independence, the previous arguments suffice to describe the relevant properties of $Q_N(\gamma)$ and $\mathbb{E}[Q_N(\gamma)]$. In this case we obtain $F_{i,j}(0,0) = F_i(0)F_j(0) = \tau^2$ and $\mathbb{C}ov[\zeta_{iN}(\gamma), \zeta_{jN}(\gamma)] = 0$ for $i \neq j$. In the case of dependence, we require a suitable assumption concerning the asymptotic behavior of $\mathbb{C}ov[\zeta_{iN}(\gamma), \zeta_{jN}(\gamma)]$. From the definition of $\zeta_{iN}(\gamma)$ in (19) follows

$$\mathbb{C}\operatorname{ov}[\zeta_{iN}(\gamma),\zeta_{jN}(\gamma)] = h_i h_j \omega_{ij}, \tag{24}$$

where ω_{ij} can be interpreted as local measures of dependence, defined as

$$\begin{split} \omega_{i,j} &= \tau \mathbb{P}(U_i > 0, U_j > 0) + (1 - \tau) \mathbb{P}(U_i \le 0, U_j \le 0) \\ &- \tau (1 - \tau) [\mathbb{P}(U_i > 0, U_j \le 0) + \mathbb{P}(U_i \le 0, U_j > 0)] \\ &= \begin{cases} F_{i,j}(0, 0) - F_i(0) F_j(0) & \text{for } i \ne j, \\ F_i(0) [1 - F_i(0)] & \text{for } i = j, \end{cases} 1 \le i, j \le N. \end{split}$$
(25)

Define the $(K + 1 \times K + 1)$ matrix Ω_N with generic elements (25). Further define the $(N \times K + 1)$ design matrix X_N with rows defined in (3) and $Z_N^{\top} := X_N^{\top} D_N^{-1}$ in analogy to (8). Then

$$\operatorname{Var}[\zeta_N(\gamma)] = \frac{1}{N} \gamma^\top \mathbf{Z}_N^\top \Omega_N \mathbf{Z}_N \gamma.$$
(26)

Given (24) and (26), we can derive the limiting distribution of $\zeta_N(\gamma)$ by introducing a suitable assumption on the limiting behavior of $\omega_{i,j}$ defined in (25).

(A3) Let
$$\omega_r = \sup_i |\omega_{i,i+r}|, r = 1, 2, \dots$$
 Then, $\omega_r \to 0$ of size $-r$, where $r > 2$.

As has been pointed out by Oberhofer and Haupt [39], only the properties of the Bernoulli process { $\psi_{\tau}(U_i)$ } and the local behavior of the distribution functions and densities in the neighborhood of v = 0 and (v, w) = (0, 0), respectively, are essential for dependence in quantile regression. Assumption (A3) implies that { $\psi_{\tau}(U_i)$ } = $\tau - \mathbb{I}(U_i \le 0)$) is strong mixing of size -r, where r > 2. Let $\sigma(\psi_i)$ denote the sigma-algebra generated by ψ_i . Then, $\sup_j |\mathbb{P}(\mathcal{B}_i \cap \mathcal{B}_{i+r}) - \mathbb{P}(\mathcal{B}_i)\mathbb{P}(\mathcal{B}_{i+r})| \to 0$, as r approaches infinity, where $\mathcal{B}_i \in \sigma(\psi_i)$ and $\mathcal{B}_{i+r} \in \sigma(\psi_{i+r})$. For mathematical convenience it is typically assumed that the error process { U_i } is strong mixing (e.g., [40]), as there exist numerous examples in time series analysis on mixing processes and a rich mathematical theory (e.g., [33,45–47]).

Lemma 1. Under assumptions (A1), (A2), and (A3), $\sum_{i=1}^{N} \zeta_{iN}(\gamma)$ converges in distribution to $c^{\top}\gamma$, where *c* is normally distributed with mean zero and covariance

$$\lim_{N\to\infty}\frac{1}{N}\boldsymbol{\gamma}^{\top}\boldsymbol{Z}_{N}^{\top}\boldsymbol{\Omega}_{N}\boldsymbol{Z}_{N}\boldsymbol{\gamma}_{N}$$

Proof. Assumption (A3) implies that $\{U_i\}$ is strong mixing of size -r, where r > 2, with mixing coefficient ω_r . Since $\{\zeta_{iN}(\gamma)\}$ is near-epoch dependent (L_1 -approximable) on $\{\psi_{\tau}(U_i)\}$, we can apply Theorem 10.2 in Pötscher and Prucha [47]. Then, the assertion follows from (26) and upon application of the Cramér-Wold device [48]. \Box

Instead of relying on the existence of moments as in least squares regression, in quantile regression we require that F_i has a density f_i (with respect to Lebesgue measure), and that f_i is informative at the conditional quantile we want to estimate.

(A4) For some $\epsilon_f > 0$, the density $f_i(v)$ of $F_i(v)$ exists for every i and $|v| \le \epsilon_f$, is uniformly continuous in i at v = 0, and $\sup_{1 \le i \le N} f_i(0) = o(\sqrt{N})$.

Define the $(N \times N)$ diagonal matrix Φ_N with elements $\varphi_i = f_i(0), 1 \le i \le N$. Then, in our main result we show that

$$\left(N^{-1}\mathbf{Z}_{N}^{\mathsf{T}}\mathbf{\Phi}_{N}\mathbf{Z}_{N}\right)^{-1}\left(N^{-1}\mathbf{Z}_{N}^{\mathsf{T}}\mathbf{\Omega}_{N}\mathbf{Z}_{N}\right)\left(N^{-1}\mathbf{Z}_{N}^{\mathsf{T}}\mathbf{\Phi}_{N}\mathbf{Z}_{N}\right)^{-1}$$
(27)

converges to the asymptotic covariance matrix of the QTR estimator as $N \rightarrow \infty$. Expression (27) reveals that dependence and heterogeneity can be treated separately in quantile regressions under non-iid settings:

$$\frac{F_{i,j}(0,0) - \tau^2}{f_i(0)f_j(0)}$$
(28)

corresponds to covariances and

$$\frac{\tau(1-\tau)}{f_i^2(0)}\tag{29}$$

to variances. This discussion reveals the need for further assumptions on the dependence structure guaranteeing the existence of (27) as $N \rightarrow \infty$.

(A5) For some $\epsilon_f > 0$, the density $f_{i,j}(v, w)$ of $F_{i,j}(v, w)$ exists for $|v| \le \epsilon_f$, $|w| \le \epsilon_f$, and $i \ne j$, is continuous at (v, w) = (0, 0) uniformly in *i* and *j*, and $\lim_{N\to\infty} N^{-1} \sum_{|r|=1}^{N} \alpha_0(r|u) = 0$, where $\alpha_0(r|u) = \sup |f_{i,i+r}(0,0) - f_i(0)f_{i+r}(0)|$, and the supremum is taken over *i*, $i + r \in \mathbb{N}$.

- (A6) $N^{-1} \mathbf{Z}_N^{\top} \Omega_N \mathbf{Z}_N$ converges for $N \to \infty$ to a $(K + 1 \times K + 1)$ matrix Σ .
- $(A7)N^{-1}Z_N^{\top}\Phi_N Z_N$ converges for $N \to \infty$ to a non-singular $(K + 1 \times K + 1)$ matrix V.

If the errors are independent, $f_{i,j}(0,0) = f_i(0)f_j(0)$ and the sum in assumption (A5) is equal to zero for all N. Then assumptions (A4) and (A5) are implied by the existence of f(v)in the neighborhood of v = 0, and the continuity of f(v) at v = 0. Assumption (A5) can be interpreted as an infinitesimal weak dependence condition similar to the "dependence index sequence" of Castellana and Leadbetter [49]. Assumptions (A6) and (A7) ensure the existence of the covariance matrix (27) in the limit. While the dependence structure is contained in assumption (A6), the heterogeneity is captured by assumption (A7). Note that a too strong dependence hinders convergence in (A6). Obviously, if the limit of the matrix $Z_N^T \Omega_N Z_N$ is singular, then the limiting distribution of $\sqrt{ND_N}(\hat{\beta} - \beta_0)$ is singular, too.

2.2.3. Main Result

The proof of our main result rests on Lemma 1 and two additional Lemmas. Lemma 1 derived the limiting distribution of $Q_N(\gamma) - \mathbb{E}[Q_N(\gamma)]$. Lemma 2 establishes that $\overline{Q}_N(\gamma) = \mathbb{E}[Q_N(\gamma)] = \mathbb{E}[B_N(\gamma)]$ (see Remark 1) converges uniformly for γ in a compact subset of \mathbb{R}^{K+1} to $\gamma^\top V \gamma$. Lemma 3 establishes that $\mathbb{V}ar[B_N(\gamma)]$ converges uniformly to 0 for γ in a compact subset of \mathbb{R}^{K+1} .

Remark 2. In our derivations we assume $h_i \leq \epsilon_f$ and apply assumptions (A4) and (A5). Hence we implicitly assume the consistency of $\hat{\gamma}$. The ULLN (15) required for this can be derived from (i) a weak LLN for weakly dependent sequences using the L₁-approximability of $\{\zeta_{iN}(\gamma)\}$ by $\{\psi_{\tau}(U_i)\}$ (see [47], Theorem 6.2, and the proof of Lemma 1) and (ii) stochastic equicontinuity of the $q_{iN}(\gamma)$ (e.g., [33], Section 21.3). Then, weak consistency of $\hat{\gamma}$ in the QTR model can be established using the arguments in (Oberhofer and Haupt [40], Lemma 2C, Theorem 1) or (Fitzenberger [36], Theorem 2.2).

Lemma 2. Under assumptions (A2), (A4), and (A7), $\mathbb{E}[B_N(\gamma)]$ converges for $N \to \infty$ to $0.5\gamma^{\top}V\gamma$. The convergence is uniform for γ in a compact subset of \mathbb{R}^{K+1} .

Proof. For $h_i > 0$, from (20) and assumption (A4) follows from the mean value theorem

$$\frac{1}{2}h_i^2\inf_{0\leq v\leq h_t}f_i(v)\leq \mathbb{E}[b_{iN}(\gamma)]\leq \frac{1}{2}h_i^2\sup_{0\leq v\leq h_t}f_i(v).$$

Apply an analogous argumentation for $h_i < 0$. From assumption (A2) follows $\sum_{i=1}^{N} h_i^2 = O(1)$. Finally, due to assumption (A4),

$$\lim_{N\to\infty} \mathbb{E}\left[\sum_{i=1}^N b_{iN}(\gamma)\right] = \lim_{N\to\infty} \frac{1}{2} \sum_{i=1}^N h_i^2 f_i(0) = \frac{1}{2} \gamma^\top \left[\lim_{N\to\infty} \frac{1}{N} \mathbf{Z}_N^\top \mathbf{\Phi}_N \mathbf{Z}_N\right] \gamma,$$

and the assertion follows from assumption (A7) (e.g., [37], Theorem 4.1).

Lemma 3. Under assumptions (A2), (A4), and (A5), $\lim_{N\to\infty} Var[B_N(\gamma)] = 0$. The convergence is uniform for γ in a compact subset of \mathbb{R}^{K+1} .

Proof. For $i \neq j$, $h_i > 0$, and $h_j > 0$,

$$\mathbb{E}[b_{iN}(\gamma)b_{jN}(\gamma)] - \mathbb{E}[b_{iN}(\gamma)b_{jN}(\gamma)] = \int_{\substack{0 \le v \le h_i \\ 0 \le v \le h_i \ 0 \le w \le h_j}} \int_{\substack{(h_i - v)(h_j - w)[f_{i,j}(v, w) - f_i(v)f_j(w)]dvdw}} \le \frac{1}{4}h_i^2h_j^2 \sup_{\substack{0 \le v \le h_i \\ 0 \le w \le h_j}} |f_{i,j}(v, w) - f_i(v)f_j(w)|.$$

An analogous argumentation applies for $h_i > 0$, $h_i < 0$ and all other cases. Then,

$$\sum_{\substack{i,j=1\\i\neq j}}^{N} \mathbb{C}ov[(b_{iN}(\gamma), b_{jN}(\gamma))] \le \frac{1}{4} \sum_{\substack{i,j=1\\i\neq j}}^{N} h_i^2 h_j^2 \sup_{\substack{|v|\le |h_i|\\|w|\le |h_j|}} |f_{i,j}(v,w) - f_i(v)f_j(w)|.$$
(30)

The expression on the right hand side is bounded from above by

$$\frac{1}{4} \frac{1}{N} \sum_{\substack{|r|=0}}^{N-1} \sup_{\substack{|v|\leq |h_i|\\|w|\leq |h_{i+r}|}} |f_{i,i+r}(v,w) - f_i(v)f_{i+r}(w)| \frac{1}{N} \sum_{j=1}^{N-r} (z_{jN}\gamma)^2 (z_{j+rN}\gamma)^2.$$

Thus, according to assumptions (A5) and (A2), the left hand side of (30) tends to zero as $N \rightarrow \infty$. Analogously, for i = j,

$$\sum_{i=1}^{N} \mathbb{V}\mathrm{ar}[b_{iN}(\gamma)] \le \sum_{i=1}^{N} \left[\frac{1}{3} |h_i|^3 \sup_{|v| \le |h_i|} f_i(v) \right] \le \frac{1}{N} \sum_{i=1}^{N} |z_{jN}\gamma|^3 \frac{1}{\sqrt{N}} \sup_{1 \le j \le N} f_j(0).$$

Then, the assertion follows from assumptions (A2) and (A4).

Theorem 1. Under assumptions (A1)–(A7), the minimizing value $\hat{\gamma}_N = \sqrt{N}D_N(\hat{\beta}_N - \beta_0)$ of $Q_N(\gamma)$ converges in distribution to a normal distribution with mean zero and covariance matrix $V^{-1}\Sigma V^{-1}$.

Proof. According to Lemmas 1–3 and assumption (A6), $Q_N(\gamma)$ converges for $N \to \infty$ in distribution to

$$Q(\boldsymbol{\gamma}) = \frac{1}{2} \boldsymbol{\gamma}^{\top} \boldsymbol{V} \boldsymbol{\gamma} + \boldsymbol{c}^{\top} \boldsymbol{\gamma}, \qquad (31)$$

with minimizing value $\hat{\gamma} = -V^{-1}c$. The limiting value $Q(\gamma)$ can be interpreted as the limit of a second order Taylor expansion of $\mathbb{E}[Q_N(\gamma)]$. The convergence in distribution of the minimizing value $\hat{\gamma}_N$ to $\hat{\gamma}$ requires the convexity of $Q_N(\gamma)$ [35] and the uniform convergence results established in our Lemmas for γ in a compact subset of \mathbb{R}^{K+1} . \Box

3. Data and Results

This section provides descriptives for the CET data from 1850 until 2020 used for fitting the QTR models and summarizes the modeling results. All temperatures are measured in Celsius degrees (°C). The CET time series is recorded by a network of monitoring stations (see [1], Figure 1a for a visualization of the locations) and designed to represent the climate of the English Midlands. The time series decomposition of the CET time series and the description of its key characteristics has received considerable attention in the literature (see, e.g., [1,4,11,12,15,29,50,51] for contributions over the past two decades).

All computations and visualizations were obtained with the statistical software R version 4.1.1 [52] using the addon packages quantreg [53], forecast [54], ggplot2 [55], and gridExtra [56]. Data and code (Supplementary Materials) is available from an online repository [57].

3.1. Descriptives for CET Anomalies

In the following analysis of the CET data, we focus on monthly temperature anomalies Y_s^* (hereafter anomalies). To construct the anomalies, we first compute mean monthly temperatures for each month over the years 1850 until 1900 \bar{Y}_s . This time period is generally considered to approximate pre-industrial levels and facilitates comparability of our results for Central England with other regions due to the availability of gridded surface temperature records (see, e.g., Chapter 3, p. 20 in [3], and the references cited therein). The anomalies Y_s^* are obtained by subtracting \bar{Y}_s for each month of the years 1961 until 2020 shown in Figure 3. Table 1 displays \bar{Y}_s for the years 1961 until 2020 (second row) and the anomalies Y_s^* . Note that all displayed anomalies are positive.

Table 1. Mean monthly temperatures (°C) for years 1850 until 1900 \bar{Y}_s (first row) and years 1961 until 2020 \tilde{Y}_s (second row). The third row displays corresponding anomalies Y_s^* for years 1961 until 2020.

	Years	Jan.	Feb.	Mar.	Apr.	May	Jun.	Jul.	Aug.	Sep.	Oct.	Nov.	Dec.
\bar{Y}_s	1850-1900	3.48	4.13	5.11	8.07	10.79	14.28	15.87	15.48	13.20	9.34	5.84	3.99
\tilde{Y}_s	1961-2020	4.24	4.33	6.21	8.42	11.53	14.42	16.42	16.14	13.88	10.77	6.98	4.81
Y_s^*	1961–2020	0.76	0.20	1.09	0.35	0.74	0.15	0.55	0.66	0.69	1.43	1.14	0.82

Figure 4 shows the time series plot of the anomalies for the years 1961 until 2020. Similar to King et al. [21], we include threshold lines for $0 \degree C$ (green), 1.5 $\degree C$ (orange), and $2 \degree C$ (red) in all visualizations of the anomalies (for detailed reasoning and brief summaries, see [58–60]).



Figure 4. Anomalies for years 1961 until 2020 (grey line). Horizontal lines indicate anomalies of 0 °C (green), 1.5 °C (orange), and 2 °C (red).

Figure 5 visualizes the anomalies for the different months (left display) and the empirical density of the anomalies (right display). Each line in the left display visualizes one time series of anomalies for a particular month from Figure 4. The anomalies show an upward trend and most of the anomaly time series are located above the green line after the year 2000. The empirical density of the anomalies shows a location shift when comparing the years 1961 until 2020 (green area) to the years 1850 until 1900 (grey area), with a mean anomaly during the years 1961 until 2020 of 0.71. Note that the empirical



density that an anomaly exceeds 2 °C (marked by the red line) triples for the years 1961 until 2020 compared to the years 1850 until 1900.

Figure 5. Anomalies for years 1961 until 2020 (**left** display), where each line represents anomalies for one month. **Right** display shows empirical densities of anomalies for years 1850 until 1900 (grey) and 1961 until 2020 (green). Horizontal lines indicate anomalies of 0 °C (green), 1.5 °C (orange), and 2 °C (red).

To obtain a more complete picture of potential heterogeneities across the different months, we investigate the months separately. Figure 6 shows boxplots for the anomalies for each month over the years 1850 until 1900 (grey) and the years 1961 until 2020 (green). The different boxplots indicate substantial heterogeneities across the months with respect to the dispersion of the distribution of the anomalies. The dispersion is higher for the winter months compared to the summer months. Additionally the boxplots illustrate that location shifts in the anomalies are more pronounced for the autumn and winter months compared to the other months. Figure 7 shows the empirical densities for the anomalies from January 1850 until December 1900 (darkgrey) and from January 1961 until December 2020 (green) separately for each month. Starting from December in the top left corner to November in the bottom right corner, the rows of the figure represent the winter (first), spring (second), summer (third), and fall (fourth) months. The figure indicates the presence of a location and scale shift in the conditional distribution of temperature anomalies, which varies across the different months and is more pronounced in fall and winter.



Figure 6. Boxplots for anomalies for years 1850 until 1900 (grey) and 1961 until 2020 (green) for each month. Horizontal lines indicate anomalies of 0 $^{\circ}$ C (green), 1.5 $^{\circ}$ C (orange), and 2 $^{\circ}$ C (red).



Figure 7. Empirical densities of anomalies plotted for each month from January 1850 until December 1900 (grey) and from January 1961 until December 2020 (green). Vertical lines indicate anomalies of $0 \,^{\circ}$ C (green), 1.5 $\,^{\circ}$ C (orange), and 2 $\,^{\circ}$ C (red).

3.2. Relevance of Quantile Regression for Analyzing Anomalies

In this subsection, we highlight properties of the anomalies that indicate that quantile regression analysis is more appropriate in this case compared to least squares analysis. Figure 8 displays QQ-plots of the sample quantiles of the anomalies (ordinate) for each month against the theoretical quantiles of the standard normal distribution (abscissa) and reveals that the empirical anomaly distributions appear "almost normal" only for months

May and November, while there are a number of deviations from normality in all other months. In particular, we observe evidence of skewness (left: January, February, June; right: April, July, August) and outliers (December, January, February, March), which may heavily influence the accuracy of least squares based statistics. Additional evidence of fat tails (September, October, December) may trouble the existence of moments.



Figure 8. Sample quantiles of anomalies for each month from January 1961 until December 2020 (ordinate) and theoretical quantiles of standard normal distribution (abscissa). Dashed grey line connects points constituted of respective 0.25%- and 0.75%-quantiles.

Figure 9 displays estimated mean versus median trends of anomalies for each month starting from December (top left) to November (bottom right). The trend regressions were fitted by employing a polynomial of degree three in a quantile regression with $\tau = 0.5$ (solid green line) and a least squares regression (dashed green line). The figure highlights that there are considerable differences between the conditional mean and the conditional median for all months except December. In particular, skewness and extreme values lead to a much more wiggly behavior of the mean trend compared to the median trend.



Figure 9. Estimated trend polynomial of degree three by quantile regression ($\tau = 0.5$; solid green line) and least squares regression (dashed green line) for anomalies for each month from January 1961 until December 2020 (grey lines).

3.3. Results of QTR Estimation

In the sense of our mathematical model in Equation (1), we analyze all monthly CET anomaly time series separately (e.g., [4,34,61]) by fitting a QTR model as defined in Equation (5) for time series of monthly temperature anomalies Y_s^* to capture month-specific patterns of trend and heterogeneity. Similar to Franzke [15], we use a low-order polynomial for trend modeling and fix the order of the polynomial regression design given in Equation (3) to K = 3:

$$\mathbf{x}_{s+Si,n} = (1, x_{s+Si,n}, x_{s+Si,n}^2, x_{s+Si,n}^3).$$

Figure 10 illustrates the QTR models fitted to the anomalies. Each display represents the QTR model fit to a particular month from December (top left) to November (bottom right). In our further analysis of the CET data, we focus on hot and cold extremes. Following Franzke [28], we define the $\tau = 0.95$ quantile as threshold for hot extremes and the $\tau = 0.05$ quantile as threshold for cold extremes. Visualizations based on the exceedance of such thresholds can be found, among others, in Mudelsee [62] and Dissanayake et al. [63]. Alternatively, Glick [64] considers the maximum and minimum as thresholds for hot and cold extremes. The area between the estimated quantile trend curves for hot ($\tau = 0.95$) and cold ($\tau = 0.05$) extremes is marked in green. The estimates indicate that there are substantial heterogeneities with respect to two dimensions: Across the quantiles (when contrasting the lines that bound the green area above and below for each display) and across the monthly seasonal cycle (when contrasting the green areas over the different displays). Upward trends in the anomalies are visible for almost all months, as larger fractions of the areas move above the zero line over time. The divergence between the hot and cold quantiles is larger for winter months compared to summer months. Overall, our results are in accordance with earlier literature, which provided evidence for global warming and trending in temperature time series and noted heterogeneity across the different months (e.g., [4,6,28]).





Figure 10. Anomalies for each month from January 1961 until December 2020. Green shaded area colors region between QTR models for hot ($\tau = 0.95$ quantile) and cold ($\tau = 0.05$ quantile) extremes. Horizontal lines indicate anomalies of 0 °C (green), 1.5 °C (orange), and 2 °C (red).

4. Concluding Remarks

We propose a model for season-specific trends motivated by cyclostationarity of temperature anomalies. From this we derive a quantile trend regression (QTR) model by introducing a set of assumptions reflecting that environmental processes such as temperature anomalies are neither homoscedastic nor independent. The key assumption about the trend is that temperature anomalies can be decomposed into a conditional temperature anomaly quantile modeled as a polynomial trend and a stationary weakly dependent remainder process. We provide asymptotic results for the QTR model and apply it to seasonal temperature anomalies of the Central England temperature (CET).

Our empirical results suggest a location and scale shift in the distribution of CET anomalies during the years 1961 until 2020 compared to 1850 until 1900 and an increase of the relative frequency of observing hot extremes (i.e., anomalies, which exceed the 95% quantile of temperature anomalies during 1850 until 1900). Additionally, our modeling results indicate the presence of substantial heterogeneities of trend patterns over the seasonal cycle for the years from 1961 to 2020, where warming is most pronounced for autumn and winter months. Overall, our findings are in line with earlier work, which documented upward trends for CET temperatures [1,6,17] and heterogeneity in trend across winter and summer months Vogelsang and Franses [4].

The QTR approach imposes a weak set of assumptions and is particularly suited for analyzing trends in temperature anomalies as highlighted by the visualizations provided in the previous section. First, robust alternatives to least squares based modeling are necessary due to the presence of outliers (in trend and remainder component), which may impact the fitted trend regression polynomials substantially. Second, accounting for heterogeneity across the different months is vital in the analysis of trends in hot and cold extremes. Finally, on a theoretical level, the non-normal behaviour of the anomaly data suggests that the assumptions underlying quantile regression are more suitable compared to the moment-based assumptions of least squares methods. There are of course limitations of the proposed QTR modeling of anomalies. Polynomials may tend to overfit in certain areas of the domain spanned by the design vectors. Local overfitting could be mitigated by imposing suitable restrictions such as shape-constraints or specifying particular distributions (e.g., [22,65]). This, however, requires a priori information on the nature of the restrictions and may seriously flaw the modeling results if the imposed restrictions are wrong. Future work could investigate suitable restrictions for temperature time series based on different time horizons and/or regions. Alternatively, QTR modeling could be extended to other processes which exhibit similar structures to the temperature process. An example are air pollutant processes, which result from the complex interplay of natural and anthropogenic influences such as diurnal cycles due to traffic or cycles according to the annual seasons. This renders the usual filtering procedures from time series analysis infeasible and QTR could be a valuable modeling approach.

Supplementary Materials: The code to reproduce the results is available online at https://github.com/markusfritsch/quantWarming.

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