# Eigenvalue perturbation theory of structured matrices under generic structured rank one perturbations: General results and complex matrices* 

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## Dedicated to G.W. (Pete) Stewart on the occasion of his 70th birthday


#### Abstract

We study the perturbation theory of structured matrices under structured rank one perturbations, and then focus on several classes of complex matrices. Generic Jordan structures of perturbed matrices are identified. It is shown that the perturbation behavior of the Jordan structures is substantially different from the corresponding theory for unstructured generic rank one perturbation as it has been studied in [26], [28], [29].


Key Words: structured matrices, Brunovsky form, complex Hamiltonian Jordan form, perturbation analysis, generic perturbation, rank one perturbation.

Mathematics Subject Classification: 15A63, 15A21, 15A57, 47A55, 93B10, 93B35, 93C73.

## 1 Introduction

In this paper, we consider the perturbation theory for Jordan structures associated with complex matrices in several classes of structured matrices under generic perturbations

[^0]that have rank one and are structure preserving. We also present results on the behavior of Jordan structures under rank one structured perturbations for rather general classes of structured matrices, both real and complex, that cover many particular cases and support the perturbation theory developed in this paper, and will be used in subsequent publications as well. The classes that we consider are defined as follows.

Let $\mathbb{F}$ denote either the field of complex numbers $\mathbb{C}$ or the field of real numbers $\mathbb{R}$ and let $I_{n}$ denote the $n \times n$ identity matrix. The superscript $(\cdot)^{T}$ denotes the transpose and $(\cdot)^{*}$ denotes the conjugate transpose of a matrix or vector; thus $X^{*}=X^{T}$ for $X \in \mathbb{R}^{m \times n}$.

Definition 1.1 Let $J \in \mathbb{F}^{2 n \times 2 n}$ be an invertible skew-symmetric matrix. A matrix $A \in \mathbb{F}^{2 n \times 2 n}$ is called $J$-Hamiltonian if $J A=(J A)^{T}$.

The classical and most important example in applications, see Section 1.3, are the classes obtained with the matrix

$$
J=\left[\begin{array}{cc}
0 & I_{n}  \tag{1.1}\\
-I_{n} & 0
\end{array}\right] .
$$

In this case we typically drop the prefix $J$ in the name of the matrix classes.
Other types of symmetries are introduced using an invertible symmetric matrix instead of a skew-symmetric $J$ in Definition 1.1:

Definition 1.2 Let $H \in \mathbb{F}^{n \times n}$ be an invertible symmetric matrix. A matrix $A \in \mathbb{F}^{n \times n}$ is called $H$-symmetric if $H A=(H A)^{T}$.

If $J$ is skew-symmetric invertible, and $N$ is such that $J N=-(J N)^{T}$, then $N$ is called $J$-skew-Hamiltonian. Note that the rank of any $J$-skew-Hamiltonian matrix is even, and since we are concerned only with rank one perturbations in this paper, $J$ -skew-Hamiltonian matrices will not be considered here. For a similar reason, we do not consider here matrices $N$ such that $H N=-(H N)^{T}$, where $H$ is symmetric and invertible.

In this paper we consider the complex case in the above definitions. The real case, as well as rank one perturbation analysis of $J$-symplectic matrices (defined by the equality $S^{T} J S=J$ with invertible skew symmetric matrix $J$ ) and $H$-orthogonal matrices (defined by the equality $S^{T} H S=H$ with invertible symmetric matrix $H$ ) will be studied in subsequent papers. An analogous but different perturbation theory for rank one structured perturbations can be also developed for the case when $H$ is taken to be Hermitian and the transpose is replaced by the conjugate transpose (in the complex case) in Definition 1.2. This will be addressed elsewhere as well.

### 1.1 Notation

In the following the set of positive integers is denoted by $\mathbb{N}$. $\mathcal{J}_{m}(\lambda)$ denotes an upper triangular $m \times m$ Jordan block with eigenvalue $\lambda$ and $R_{m}$ stands for the $m \times m$ matrix
with 1 on the leftbottom - topright diagonal and zeros elsewhere, i.e.,

$$
\mathcal{J}_{m}(\lambda)=\left[\begin{array}{cccc}
\lambda & 1 & & 0 \\
& \lambda & \ddots & \\
& & \ddots & 1 \\
0 & & & \lambda
\end{array}\right], \quad R_{m}=\left[\begin{array}{ccc}
0 & & 1 \\
& . & \\
1 & & 0
\end{array}\right]
$$

The $k$-th standard basis vector of length $n$ will be denoted by $e_{k, n}$ or in short $e_{k}$ if the length is clear from the context. The spectrum of a matrix $A \in \mathbb{F}^{n \times n}$, i.e., the set of eigenvalues including possibly nonreal eigenvalues of real matrices, is denoted by $\sigma(A)$. An eigenvalue $\lambda \in \sigma(A)$ is said to be simple if the corresponding algebraic multiplicity is one, i.e., $\lambda$ is a simple root of the characteristic polynomial of $A$.
$\chi(Z)=\operatorname{det}(x I-Z)$ is the characteristic polynomial of a square size matrix $Z$.
Throughout the paper we will use a fixed matrix norm $\|\cdot\|$ which denotes the spectral norm $\|\cdot\|_{2}$.

A block diagonal matrix with diagonal blocks $X_{1}, \ldots, X_{q}$ (in that order) is denoted by $X_{1} \oplus X_{2} \oplus \cdots \oplus X_{q}$.

If $v^{T}=\left[v_{1}, \ldots, v_{n}\right]^{T} \in \mathbb{C}^{n}$ then Toep $(v)$ denotes the $n \times n$ upper triangular Toeplitz matrix

$$
\operatorname{Toep}(v)=\left[\begin{array}{cccc}
v_{1} & v_{2} & \ldots & v_{n} \\
0 & v_{1} & \ddots & \vdots \\
\vdots & \ddots & \ddots & v_{2} \\
0 & \ldots & 0 & v_{1}
\end{array}\right]
$$

We also introduce the anti-diagonal matrices

$$
\Sigma_{k}=\left[\begin{array}{cccc}
0 & \cdots & 0 & (-1)^{0}  \tag{1.2}\\
\vdots & \cdot & (-1)^{1} & 0 \\
0 & . & & \vdots \\
(-1)^{k-1} & 0 & \cdots & 0
\end{array}\right]=\left[\right]=(-1)^{k-1} \Sigma_{k}^{T}
$$

i.e., $\Sigma_{k}$ is symmetric if $k$ is odd, and skew-symmetric if $k$ is even.

### 1.2 Motivation

The perturbation theory for eigenvalues of matrices is well established [31]. This is also the case if the perturbations are generic low rank matrices, see [4, 19, 26, 28, 29]. But when the perturbations are restricted to be structure preserving then surprisingly different effects may occur.

Example 1.3 Let

$$
A=\left[\begin{array}{cc}
\mathcal{J}_{3}(0) & 0 \\
0 & -\mathcal{J}_{3}(0)^{T}
\end{array}\right], \quad J=\left[\begin{array}{cc}
0 & I_{3} \\
-I_{3} & 0
\end{array}\right] .
$$

Then $A$ is $J$-Hamiltonian and has two Jordan blocks of size 3 associated with the eigenvalue 0 . The perturbation analysis under unstructured generic rank 1 perturbations, Theorem 3.1 in [26] (a particular case of which is part of Theorem 2.3 below), yields that the perturbed matrix still has one block $J_{3}(0)$, while the other block has vanished and split into three (generically different) nonzero eigenvalues.

In contrast to this (as we will show below) a generic Hamiltonian rank one perturbation will lead to a Jordan structure with a $4 \times 4$ block $\mathcal{J}_{4}(0)$ plus two (generically) nonzero simple eigenvalues. Thus the size of the largest block even increases.

This example demonstrates that the classical understanding of perturbation theory has to be changed for classes of structured matrices. The perturbation theory for structured generic low rank perturbations is dominated by two conflicting effects, the generic structured perturbation trying to destroy the most sensitive part in the Jordan structure (which is the largest Jordan block) and the structure which requires certain Jordan structures.

### 1.3 Applications

The perturbation theory that we present in this paper has several important applications in control.

Let us first discuss the problem of passivity of systems. Consider a linear timeinvariant control system

$$
\begin{align*}
\dot{x} & =A x+B u, \quad x(0)=0 \\
y & =C x+D u \tag{1.3}
\end{align*}
$$

with matrices $A \in \mathbb{F}^{n, n}, B \in \mathbb{F}^{n, m}, C \in \mathbb{F}^{p, n}, D \in \mathbb{F}^{p, m}$. Here $u$ is the input, $x$ the state, and $y$ the output. Let us assume that all eigenvalues of $A$ are in the open left half complex plane and that $D$ is square and invertible. The system is called passive, see e.g., [2], if there exists a nonnegative scalar valued function $\Theta$ such that the dissipation inequality

$$
\Theta\left(x\left(t_{1}\right)\right)-\Theta\left(x\left(t_{0}\right)\right) \leq \int_{t_{0}}^{t_{1}} u^{*} y+y^{*} u d t
$$

holds for all $t_{1} \geq t_{0}$, i.e., the system absorbes supply energy. It is well known, $[2,16]$, that one can check whether the system is passive by checking whether the Hamiltonian matrix

$$
\mathcal{H}=\left[\begin{array}{cc}
F & G  \tag{1.4}\\
H & -F^{*}
\end{array}\right]:=\left[\begin{array}{cc}
A-B R^{-1} C & -B R^{-1} B^{*} \\
-C^{*} R^{-1} C & -\left(A-B R^{-1} C\right)^{*}
\end{array}\right]
$$

has no purely imaginary eigenvalues, where $R=D+D^{*}$.
In many real world applications the system model (1.3) is only an approximation arising from a discretization of an infinite dimensional problem, a linearization of a nonlinear system, a realization or a reduced order approximate model, see, e.g., [12,
$13,16,17,27,30]$ and often in this approximation process the passivity is lost, and one tries to modify the non-passive approximate system by a small norm (typically also small rank) perturbation to a nearby passive system. Our perturbation theory will be important in understanding and computing minimal perturbations.

Another important application arises in robust control. Consider a control system

$$
\begin{align*}
& \dot{x}(t)=A x(t)+B_{1} w(t)+B_{2} u(t), \quad x\left(t_{0}\right)=x^{0}, \\
& z(t)=C_{1} x(t)+D_{11} w(t)+D_{12} u(t)  \tag{1.5}\\
& y(t)=C_{2} x(t)+D_{21} w(t)+D_{22} u(t)
\end{align*}
$$

In this system, $x$ is again the state, $u$ the input, and $w$ is an exogenous input that may include noise, linearization errors, and un-modelled dynamics. The vector $y$ contains measured outputs, while $z$ is a regulated output or an estimation error.

The optimal $\mathcal{H}_{\infty}$ control problem is the task of designing a dynamic controller that minimizes (or at least approximately minimizes) the influence of the disturbances $w$ on the output $z$ in the $\mathcal{H}_{\infty}$-norm, see [33]. The computation of this controller is usually achieved by first solving two Hamiltonian eigenvalue problems that both are low rank perturbations (rank one in the single input case) of other Hamiltonian matrices where the perturbation matrices depend on the same parameter $\gamma$ that gives an upper bound for the $\mathcal{H}_{\infty}$-norm to be minimized. By minimizing the value of $\gamma$ under certain constraints then allows to find the optimal controller. Very often the optimal solution is obtained when an eigenvalue of the Hamiltonian matrix (as a function of $\gamma$ ) hits the imaginary axis, and thus becomes a multiple eigenvalue. The structured perturbation analysis of the eigenvalues as functions of this low rank perturbation allows the analysis and computation of the optimal controller, see [5, 25].

There are many further applications of the perturbation theory for structured matrices, such as the analysis of numerical methods for the Hamiltonian eigenvalue problem or its generalizations, see, e.g., $[1,7,6,11,24]$ or the solution of algebraic Riccati equations [18, 20]. Although in most applications the system matrices are real, in this paper we first study the complex case to lay down the basis for the structured perturbation theory. The real case will be discussed in a subsequent paper.

## 2 General results

In this section we recall and/or derive some mathematical results on structured rank one perturbations that will become important in the further analysis, in this and subsequent papers.

### 2.1 Perturbation theory under generic perturbations

We say that a set $W \subseteq \mathbb{F}^{n}$ (abbreviation for $\mathbb{F}^{n \times 1}$ ) is algebraic if there exists a finite set of polynomials $f_{1}\left(x_{1}, \ldots, x_{n}\right), \ldots, f_{k}\left(x_{1}, \ldots, x_{n}\right)$ with coefficients in $\mathbb{F}$ such that a
vector $\left[a_{1}, \ldots, a_{n}\right]^{T} \in \mathbb{F}^{n}$ belongs to $W$ if and only if

$$
f_{j}\left(a_{1}, \ldots, a_{n}\right)=0, \quad j=1,2, \ldots, k .
$$

In particular, the empty set is algebraic and $\mathbb{F}^{n}$ is algebraic. We say that a set $W \subseteq \mathbb{F}^{n}$ is generic if $W$ is not empty and the complement $\mathbb{F}^{n} \backslash W$ is contained in the union of finitely many algebraic sets. Clearly, if the set $W \subseteq \mathbb{F}^{n}$ is generic and if $S \in \mathbb{F}^{n \times n}$ is invertible then $S W$ is also generic. In the following, if we say that a set $W \subseteq \mathbb{F}^{n} \times \mathbb{F}^{n}$ is generic if $W$, canonically identified with a subset of $\mathbb{F}^{2 n}$ is generic as a subset of $\mathbb{F}^{2 n}$.

The following lemma is almost obvious, but useful.
Lemma 2.1 Let $Y\left(x_{1}, \ldots, x_{r}\right) \in \mathbb{F}^{m \times n}$ be a matrix whose entries are polynomials in the variables $x_{1}, \ldots, x_{r}$. If

$$
\operatorname{rank} Y\left(a_{1}, \ldots, a_{r}\right)=k
$$

for some $\left[a_{1}, \ldots, a_{r}\right]^{T} \in \mathbb{F}^{r}$, then the set

$$
\begin{equation*}
\left\{\left[b_{1}, \ldots, b_{r}\right]^{T} \in \mathbb{F}^{r}: \operatorname{rank} Y\left(b_{1}, \ldots, b_{r}\right) \geq k\right\} \tag{2.1}
\end{equation*}
$$

is generic.
Proof. Let

$$
f_{j}\left(x_{1}, \ldots, x_{r}\right)=\operatorname{det} Y_{j}\left(x_{1}, \ldots, x_{r}\right), \quad j=1,2, \ldots, s
$$

where $Y_{1}\left(x_{1}, \ldots, x_{r}\right), \ldots, Y_{s}\left(x_{1}, \ldots, x_{r}\right)$ are the $k \times k$ submatrices of $Y\left(x_{1}, \ldots, x_{r}\right)$. Then the complement of the set (2.1) consists of the common zeros of the polynomials $f_{1}, \ldots, f_{s}$, i.e., it is an algebraic set, and the set (2.1) is nonempty by hypothesis. This shows that (2.1) is generic.

In the following result we discuss ranks of powers of generic rank one perturbations to nilpotent matrices.

Theorem 2.2 Consider a matrix $A \in \mathbb{F}^{n \times n}$ satisfying $A^{m}=0$ for some $m \in \mathbb{N}$.
(1) If $X \in \mathbb{F}^{n \times n}$ is any rank one matrix, then

$$
\operatorname{rank}\left((A+X)^{m}\right) \leq m
$$

(2) If in addition $A^{m-1} \neq 0$, then

$$
\operatorname{rank}\left(\left(A+u v^{T}\right)^{m}\right)=m
$$

$$
\text { for a generic set of vectors }\left[\begin{array}{l}
u \\
v
\end{array}\right] \in \mathbb{F}^{2 n} \text {. }
$$

(3) If in addition $A^{m-1} \neq 0$, then for every invertible $B, C \in \mathbb{F}^{n \times n}$ we have

$$
\operatorname{rank}\left(\left(A+C u u^{T} B\right)^{m}\right)=m
$$

for a generic set of vectors $u \in \mathbb{F}^{n}$.

Proof. Multiplying out $(A+X)^{m}$ we obtain

$$
\begin{align*}
(A+X)^{m} & =A^{m}+A^{m-1} X+A^{m-2} X(A+X)+A^{m-3} X(A+X)^{2} \\
& +\cdots+A X(A+X)^{m-2}+X(A+X)^{m-1} \tag{2.2}
\end{align*}
$$

Since $A^{m}=0$ and all other summands in the right hand side of (2.2) have ranks at most one, part (1) follows.

For part (2) let us assume without loss of generality that $A$ is in Jordan canonical form, i.e.,

$$
\begin{equation*}
A=J_{k_{1}}(0) \oplus J_{k_{2}}(0) \oplus \cdots \oplus J_{k_{t}}(0) \tag{2.3}
\end{equation*}
$$

where $k_{j} \leq k_{1}=m, j=2,3, \ldots, t$.
We obviously have that

$$
\operatorname{rank}\left(\left(A+e_{m} e_{1}^{T}\right)^{m}\right)=m
$$

so by Lemma 2.1, part (2) follows (note that we cannot have $\operatorname{rank}\left(\left(A+u v^{T}\right)^{m}\right)>m$ by part (1)).

Finally, consider part (3). From

$$
A+C u u^{T} B=C\left(C^{-1} A C+u u^{T} B C\right) C^{-1}
$$

we see that without loss of generality we may assume that $C=I$. Furthermore,

$$
A+u u^{T} B=X\left(X^{-1} A X+\left(X^{-1} u\right)\left(u^{T}\left(X^{-1}\right)^{T}\right)\left(X^{T} B X\right)\right) X^{-1}
$$

and choosing the invertible matrix $X$ so that $X^{-1} A X$ is in Jordan canonical form, we may also assume without loss of generality that $A$ is given by (2.3). Denote by $\Delta(u)$, $u=\left[u_{1}, \ldots, u_{n}\right]^{T} \in \mathbb{F}^{n}$, the determinant of the $m \times m$ upper left corner of the matrix $\left(A+u u^{T} B\right)^{m}$. In view of Lemma 2.1, we only need to show that

$$
\begin{equation*}
\Delta\left(w_{1}, \ldots, w_{n}\right) \neq 0 \quad \text { for some } w_{1}, \ldots, w_{n} \in \mathbb{F} \tag{2.4}
\end{equation*}
$$

By (2.2) we have that

$$
\begin{align*}
& \left(A+u u^{T} B\right)^{m} \\
= & A^{m-1} u u^{T} B+A^{m-2} u u^{T} B\left(A+u u^{T} B\right)+A^{m-3} u u^{T} B\left(A+u u^{T} B\right)^{2}+  \tag{2.5}\\
& \cdots+A u u^{T} B\left(A+u u^{T} B\right)^{m-2}+u u^{T} B\left(A+u u^{T} B\right)^{m-1} .
\end{align*}
$$

This formula shows that $\Delta(u)$ is a polynomial in $u_{1}, \ldots, u_{n}$ of the form

$$
\Delta(u)=\Delta_{2 m}(u)+\Delta_{2 m+2}(u)+\cdots+\Delta_{2 m^{2}}(u),
$$

where $\Delta_{p}(u)$ is a homogeneous polynomial of degree $p$. Clearly, to prove (2.4), we only need to find $w_{1}, \ldots, w_{n} \in \mathbb{F}$ such that

$$
\begin{equation*}
\Delta_{2 m}\left(w_{1}, \ldots, w_{n}\right) \neq 0 \tag{2.6}
\end{equation*}
$$

(here we use the easily proved fact, that if one homogeneous component of a polynomial in several variables takes a nonzero value, then the whole polynomial takes a nonzero value).

Note that $\Delta_{2 m}(u)$ is the determinant of the upper left $m \times m$ corner of the matrix

$$
\begin{equation*}
A^{m-1} u u^{T} B+A^{m-2} u u^{T} B A+A^{m-3} u u^{T} B A^{2}+\cdots+A u u^{T} B A^{m-2}+u u^{T} B A^{m-1} \tag{2.7}
\end{equation*}
$$

(cf. formula (2.5)). Let $\left[b_{1}, \ldots, b_{n}\right]^{T}$ be the first column of $B$. Since the upper left $m \times m$ corner of $A$ is the nilpotent Jordan block $J_{m}(0)$, it follows that the upper left $m \times m$ corner of the matrix (2.7) is an upper triangular $m \times m$ matrix with

$$
u_{m}\left(u_{1} b_{1}+\ldots+u_{n} b_{n}\right)
$$

on the main diagonal. Clearly, one can choose $w_{1}, \ldots, w_{n} \in \mathbb{F}$ so that

$$
w_{m}\left(w_{1} b_{1}+\ldots+w_{n} b_{n}\right) \neq 0,
$$

(here we use the hypothesis that $B$ is invertible, and therefore at least one of $b_{1}, \ldots, b_{n}$ is nonzero), and (2.6) follows.

### 2.2 Unstructured generic rank one perturbation theory

The general perturbation analysis for generic low rank perturbations has been studied in $[19,26,28,29]$. For the case of rank one perturbations - which is of interest in this paper - we have the following result.

Theorem 2.3 Let $A \in \mathbb{C}^{n \times n}$ be a matrix having the pairwise distinct eigenvalues $\lambda_{1}, \ldots, \lambda_{p}$ with geometric multiplicities $g_{1}, \ldots, g_{p}$ and having the Jordan canonical form

$$
\bigoplus_{k=1}^{g_{1}} \mathcal{J}_{n_{1, k}}\left(\lambda_{1}\right) \oplus \cdots \oplus \bigoplus_{k=1}^{g_{p}} \mathcal{J}_{n_{p, k}}\left(\lambda_{p}\right),
$$

where $n_{j, 1} \geq \cdots \geq n_{j, g_{j}}$, $j=1, \ldots, p$. Consider the rank one matrix $B=u v^{T}$, with $u, v \in \mathbb{C}^{n}$. Then generically (with respect to the entries of $u$ and $v$ ) the Jordan blocks of $A+B$ with eigenvalue $\lambda_{j}$ are just the $g_{j}-1$ smallest Jordan blocks of $A$ with eigenvalue $\lambda_{j}$, and all other eigenvalues of $A+B$ are simple; if $g_{j}=1$, then generically $\lambda_{j}$ is not an eigenvalue of $A+B$.

More precisely, there is a generic set $\Omega \subseteq \mathbb{C}^{n} \times \mathbb{C}^{n}$ such that for every $(u, v) \in \Omega$, the Jordan structure of $A+u v^{T}$ is described in (a) and (b) below:
(a) the Jordan structure of $A+u v^{T}$ for the eigenvalues $\lambda_{1}, \ldots, \lambda_{p}$ is given by

$$
\bigoplus_{k=2}^{g_{1}} \mathcal{J}_{n_{1, k}}\left(\lambda_{1}\right) \oplus \cdots \oplus \bigoplus_{k=2}^{g_{p}} \mathcal{J}_{n_{p, k}}\left(\lambda_{p}\right)
$$

(b) the eigenvalues of $A+u v^{T}$ that are different from any of $\lambda_{1}, \ldots, \lambda_{p}$, are all simple.

Part (a) of Theorem 2.3 is the main theorem of [26] specialized to the case of rank one perturbations; a result similar to that of [26] has been obtained in [28]. For the proof of part (b) we need some preparations. We start with the following well known example:

Example 2.4 Let

$$
Z^{(1)}(\lambda, \alpha)=\mathcal{J}_{m}(\lambda)+\alpha e_{m} e_{1}^{T}=\left[\begin{array}{cccc}
\lambda & 1 & \ldots & 0 \\
0 & \lambda & \ddots & \vdots \\
\vdots & \ddots & \ddots & 1 \\
\alpha & \cdots & 0 & \lambda
\end{array}\right] \in \mathbb{C}^{m \times m}, \quad \lambda \in \mathbb{C}, \alpha \in \mathbb{C} \backslash\{0\}
$$

Then $\chi\left(Z^{(1)}(\lambda, \alpha)\right)=(x-\lambda)^{m}-\alpha$; in particular, $Z^{(1)}(\lambda, \alpha)$ has $m$ distinct eigenvalues.

Next, we note that by [26], it follows that there exists a generic set $\Omega^{\prime}$ of vectors $(u, v) \in \mathbb{C}^{n} \times \mathbb{C}^{n}$ for which (a) holds. Clearly, we may assume $\Omega^{\prime}$ is open. We then obtain the following lemma.

Lemma 2.5 Let $A$ be as in Theorem 2.3. Then there exists $\epsilon>0$ and an open dense (in $\left\{(u, v) \in \mathbb{C}^{n} \times \mathbb{C}^{n}:\|u\|,\|v\|<\epsilon\right\}$ ) set $\Omega^{\prime \prime} \subseteq \Omega^{\prime}$ such that for every $(u, v) \in \Omega^{\prime \prime}$, $\|u\|,\|v\|<\epsilon$, the Jordan form of $A+u v^{T}$ is as in Theorem 2.3.

Proof. Denote by $D(z, \epsilon)$ the closed disc of radius $\epsilon$ centered at $z \in \mathbb{C}$. Let $\epsilon>0$ be so small that for every $u, v \in \mathbb{C}^{n}$ with $\|u\|,\|v\|<\epsilon$, all eigenvalues of $A+B$ are within the union of the closed pairwise nonintersecting discs of radius $\epsilon$ centered at each of the points $\lambda_{1}, \ldots, \lambda_{p}$. It will be assumed from now on in that $\|u\|,\|v\|<\epsilon$.

Let $\chi\left(\lambda_{j}, u, v\right)$ for $j=1,2, \ldots, p$ be the characteristic polynomials in the independent variable $x$ for the restrictions of $A+B$ to its spectral invariant subspaces corresponding to the eigenvalues of $A+B$ within the disc $D\left(\lambda_{j}, \epsilon\right)$. Notice that the coefficients of $\chi\left(\lambda_{j}, u, v\right)$ are analytic functions of the components of $u$ and $v$. Indeed, this follows from the formula for the projection onto the spectral invariant subspace

$$
\frac{1}{2 \pi i} \int_{\Gamma}(z I-(A+B))^{-1} d z
$$

for a suitable closed simple contour $\Gamma$. (To prove the analyticity of the integral as function of $u$ and $v$, use approximation of the integral by Riemann sums, and within every summand of the Riemann sum use the formula

$$
\left(z_{0} I-(A+B)\right)^{-1}=\left(\operatorname{adj}\left(z_{0} I-(A+B)\right)\right) /\left(\operatorname{det}\left(z_{0} I-(A+B)\right)\right)
$$

where adj $Z$ stands for the algebraic adjoint of a matrix $Z$ ).

Let $q\left(\lambda_{j}, u, v\right)$ be the number of distinct eigenvalues of $A+B$ in the disc $D\left(\lambda_{j}, \epsilon\right)$ Let

$$
\left.q_{\max }\left(\lambda_{j}\right)=\max _{u, v \in \mathbb{C}^{n},}^{\|u\|,\|v\|<\epsilon}{ }^{2}\left(\lambda_{j}, u, v\right)\right\} .
$$

Next, we fix $\lambda_{j}$. Denote by $S\left(p_{1}, p_{2}\right)$ the Sylvester resultant matrix of the two polynomials $p_{1}(x), p_{2}(x)$ (see, e.g., $[3,15]$ ); note that $S\left(p_{1}, p_{2}\right)$ is a square matrix of size degree $\left(p_{1}\right)+\operatorname{degree}\left(p_{2}\right)$ and recall the well known fact (see [22] for example) that the rank deficiency of $p_{1}(x), p_{2}(x)$ coincides with the degree of the greatest common divisor of the polynomials $p_{1}(x)$ and $p_{2}(x)$. We have

$$
q\left(\lambda_{j}, u, v\right)=\operatorname{rank} S\left(\chi\left(\lambda_{j}, u, v\right), \frac{\partial \chi\left(\lambda_{j}, u, v\right)}{\partial x}\right)-\left(n_{j, 1}+\cdots+n_{j, g_{j}}\right)+1
$$

The entries of $S\left(\chi\left(\lambda_{j}, u, v\right), \frac{\partial \chi\left(\lambda_{j}, u, v\right)}{\partial x}\right)$ are scalar (independent of $\left.u, v\right)$ multiples of the coefficients of $\chi\left(\lambda_{j}, u, v\right)$, and therefore the set $Q\left(\lambda_{j}\right)$ of all vectors $(u, v) \in \mathbb{C}^{n} \times \mathbb{C}^{n}$, $\|u\|,\|v\|<\epsilon$, for which $q\left(\lambda_{j}, u, v\right)=q_{\max }\left(\lambda_{j}\right)$ is the complement of the set of common zeros of finitely many analytic functions of the components of $u$ and $v$. In particular, $Q\left(\lambda_{j}\right)$ is open and dense in

$$
\left\{(u, v) \in \mathbb{C}^{n} \times \mathbb{C}^{n}:\|u\|,\|v\|<\epsilon\right\}
$$

On the other hand, still for a fixed $\lambda_{j}$, consider

$$
u_{0}:=\frac{1}{2} \epsilon^{n}\left[\begin{array}{c}
u_{1}  \tag{2.8}\\
\vdots \\
u_{p}
\end{array}\right]
$$

where the vectors $u_{k} \in \mathbb{C}^{n_{k, 1}+\cdots+n_{k, g_{k}}}$ are such that all $u_{k}$ 's are zeros except for $u_{j}$ which has 1 in the $n_{j, 1}$ th position and zeros elsewhere. Also let

$$
v_{0}=\left[\begin{array}{c}
v_{1} \\
\vdots \\
v_{p}
\end{array}\right]
$$

partitioned conformably with (2.8), where all all $v_{k}$ 's are zeros except for $v_{j}$ which has 1 in the first position and zeros elsewhere. One checks easily (cf. Example 2.4) that in the disc $D\left(\lambda_{j}, \epsilon\right)$ the matrix $A+u_{0} v_{0}^{T}$ has:
(1) $n_{j, 1}$ simple eigenvalues different from $\lambda_{j}$; and
(2) the eigenvalue $\lambda_{j}$ with partial multiplicities $n_{j, 2}, \ldots, n_{j, g_{j}}$.

If by chance the pair $\left(u_{0}, v_{0}\right)$ is not in $\Omega^{\prime}$, then we slightly perturb $\left(u_{0}, v_{0}\right)$ to obtain a new pair $\left(u_{0}^{\prime}, v_{0}^{\prime}\right) \in \Omega^{\prime}$ such that (1) and (2) are still valid for the matrix $A+u_{0}^{\prime}\left(v_{0}^{\prime}\right)^{T}$.
(This is possible because $\Omega^{\prime}$ is generic, the property of eigenvalues being simple persists under small perturbations, and the total number of eigenvalues of $A+u v^{T}$ within $D\left(\lambda_{j}, \epsilon\right)$, counted with multiplicities, is equal to $n_{j, 1}+\cdots+n_{j, g_{j}}$, for every $(u, v) \in \mathbb{C}^{n}$, $\|u\|,\|v\|<\epsilon$.) Since $\Omega^{\prime}$ is open, clearly there exists $\delta>0$ such that (1) and (2) are valid for every $A+u v^{T}$, where $(u, v) \in \mathbb{C}^{n} \times \mathbb{C}^{n}$ and $\left\|u-u_{0}\right\|,\left\|v-v_{0}\right\|<\delta$. Since the set of all such pairs of vectors $(u, v)$ is open in $\mathbb{C}^{n} \times \mathbb{C}^{n}$, it follows from the properties of the set $Q\left(\lambda_{j}\right)$ established in the preceding paragraph that in fact we have

$$
q\left(\lambda_{j}, u, v\right)=q_{\max }\left(\lambda_{j}\right), \quad \text { for all } u, v \in \mathbb{C}^{n}, \quad\left\|u-u_{0}\right\|,\left\|v-v_{0}\right\|<\delta
$$

So for the following open set

$$
\Omega_{j}^{(1)}:=Q\left(\lambda_{j}\right) \cap \Omega^{\prime}
$$

which is dense in $\left\{(u, v) \in \mathbb{C}^{n} \times \mathbb{C}^{n}:\|u\|,\|v\|<\epsilon\right\}$, we have that the part of Jordan form of $A+u v^{T}$, where $(u, v) \in \Omega_{j}^{(1)}$, corresponding to the eigenvalues within $D\left(\lambda_{j}, \epsilon\right)$ consists of

$$
\mathcal{J}_{n_{j, 2}}\left(\lambda_{j}\right) \oplus \cdots \oplus \mathcal{J}_{n_{j, p_{j}}}\left(\lambda_{j}\right)
$$

and $n_{j, 1}$ simple eigenvalues different from $\lambda_{j}$.
Now let

$$
\Omega^{\prime \prime}=\left(\bigcap_{j=1}^{p} \Omega_{j}^{(1)}\right) \cap \Omega^{\prime}
$$

to satisfy Lemma 2.5.
Proof of Theorem 2.3. As noted above, in view of the main result of [26], we only need to prove part (b). Let $\chi(u, v)$ be the characteristic polynomial (in the independent variable $x$ ) of $A+B$. Then the number of distinct roots of $\chi(u, v)$ is given by the rank of the Sylvester resultant matrix $S\left(\chi(u, v), \frac{\partial}{\partial x} \chi(u, v)\right)$ minus $n-1$ (cf. the proof of Lemma 2.5). Therefore, the set $\Omega_{0}$ of all pairs of vectors $(u, v)$ on which the number of distinct roots of $\chi(u, v)$ is maximal, is a generic set. By Lemma 2.5, the maximal number of distinct roots of $\chi(u, v)$ is equal to

$$
n_{1,1}+\cdots+n_{p, 1}+\sum_{j=1}^{p} \min \left\{g_{j}-1,1\right\}
$$

Thus, for the generic set $U=\Omega_{0} \cap \Omega^{\prime}$ the Jordan structure of $A+u v^{T}$ is described by (a) and (b), as required.

We will re-prove the part (a) of Theorem 2.3 in Section 2.4, using the Brunovsky canonical form.

### 2.3 Structured canonical forms

In the following we will recall the canonical forms for $J$-Hamiltonian and $H$-symmetric matrices which is available in many sources, see, e.g., [20, 23], or [21, 32] in the framework of pairs of symmetric and skew-symmetric matrices.

Theorem 2.6 Let $H \in \mathbb{C}^{n \times n}$ be symmetric and invertible and let $A \in \mathbb{C}^{n \times n}$ be $H$ symmetric. Then there exists an invertible matrix $P \in \mathbb{C}^{n \times n}$ such that

$$
\begin{equation*}
P^{-1} A P=\mathcal{J}_{n_{1}}\left(\lambda_{1}\right) \oplus \cdots \oplus \mathcal{J}_{n_{m}}\left(\lambda_{m}\right), \quad P^{T} H P=R_{n_{1}} \oplus \cdots \oplus R_{n_{m}}, \tag{2.9}
\end{equation*}
$$

where $n_{1}, \ldots, n_{m} \in \mathbb{N}$ and $\lambda_{1}, \ldots, \lambda_{m} \in \mathbb{C}$ are not necessarily pairwise distinct. The form (2.9) is uniquely determined by the pair $(A, H)$, up to a simultaneous permutation of diagonal blocks in the right hand sides of (2.9).

Theorem 2.7 Let $J \in \mathbb{C}^{n \times n}$ be skew-symmetric and invertible (i.e., $n$ is even), and let $A \in \mathbb{C}^{n \times n}$ be J-Hamiltonian. Then there exists an invertible matrix $P \in \mathbb{C}^{n \times n}$ such that $P^{-1} A P$ and $P^{T} J P$ are block diagonal matrices

$$
\begin{equation*}
P^{-1} A P=A_{1} \oplus A_{2} \oplus A_{3}, \quad P^{T} J P=J_{1} \oplus J_{2} \oplus J_{3}, \tag{2.10}
\end{equation*}
$$

where the blocks have the following forms.

$$
\begin{equation*}
A_{1}=\mathcal{J}_{2 n_{1}}(0) \oplus \cdots \oplus \mathcal{J}_{2 n_{p}}(0), \quad J_{1}=\Sigma_{2 n_{1}} \oplus \cdots \oplus \Sigma_{2 n_{p}} \tag{i}
\end{equation*}
$$

with $n_{1}, \ldots, n_{p} \in \mathbb{N}$;

$$
\begin{gather*}
A_{2}=\left[\begin{array}{cc}
\mathcal{J}_{2 m_{1}+1}(0) & 0 \\
0 & \mathcal{J}_{2 m_{1}+1}(0)
\end{array}\right] \oplus \cdots \oplus\left[\begin{array}{cc}
\mathcal{J}_{2 m_{q}+1}(0) & 0 \\
0 & \mathcal{J}_{2 m_{q}+1}(0)
\end{array}\right],  \tag{ii}\\
J_{2}=\left[\begin{array}{cc}
0 & \Sigma_{2 m_{1}+1} \\
-\Sigma_{2 m_{1}+1} & 0
\end{array}\right] \oplus \cdots \oplus\left[\begin{array}{cc}
0 & \Sigma_{2 m_{q}+1} \\
-\Sigma_{2 m_{q}+1} & 0
\end{array}\right]
\end{gather*}
$$

with $m_{1}, \ldots, m_{q} \in \mathbb{N} \cup\{0\} ;$
(iii)

$$
A_{3}=A_{3,1} \oplus \cdots \oplus A_{3, k}, \quad J_{3}=J_{3,1} \oplus \cdots \oplus J_{3, k}
$$

where

$$
\begin{gathered}
A_{3, j}=\left[\begin{array}{cc}
\mathcal{J}_{\ell_{j, 1}}\left(\lambda_{j}\right) & 0 \\
0 & -\mathcal{J}_{\ell_{j, 1}}\left(\lambda_{j}\right)^{T}
\end{array}\right] \oplus \cdots \oplus\left[\begin{array}{cc}
\mathcal{J}_{\ell_{j, q_{j}}}\left(\lambda_{j}\right) & 0 \\
0 & -\mathcal{J}_{\ell_{j, q_{j}}}\left(\lambda_{j}\right)^{T}
\end{array}\right], \\
J_{3, j}=\left[\begin{array}{cc}
0 & I_{\ell_{j, 1}} \\
-I_{\ell_{j, 1}} & 0
\end{array}\right] \oplus \cdots \oplus\left[\begin{array}{cc}
0 & I_{\ell_{j, q_{j}}} \\
-I_{\ell_{j, q_{j}}} & 0
\end{array}\right],
\end{gathered}
$$

with $\ell_{j, 1}, \ldots, \ell_{j, q_{j}} \in \mathbb{N}$ and $\lambda_{j} \in \mathbb{C}$ with $\operatorname{Re}\left(\lambda_{j}\right)>0$ or $\operatorname{Re}\left(\lambda_{j}\right)=0$ and $\operatorname{Im}\left(\lambda_{j}\right)>0$ for $j=1, \ldots, k$. Moreover, $\lambda_{1}, \ldots, \lambda_{k}$ are pairwise distinct.

The form (2.10) is uniquely determined by the pair $(A, J)$, up to a simultaneous permutation of diagonal blocks in the right hand sides of (2.10).

### 2.4 The Brunovsky form

To analyze the effect of rank one perturbations, we will make use of the following theorem, which follows directly from the Brunovsky canonical form, [9], see also [10] or [14] for example, of general multi-input control systems $\dot{x}=A x+B u$ under transformations

$$
(A, \quad B) \quad \mapsto \quad\left(C^{-1}(A+B R) C, \quad C^{-1} B D\right)
$$

with invertible $C, D$, and arbitrary $R$ of suitable sizes.
Theorem 2.8 Let $A \in \mathbb{C}^{n \times n}$ be a matrix in Jordan canonical form

$$
\begin{equation*}
A=\mathcal{J}_{n_{1}}\left(\lambda_{1}\right) \oplus \cdots \oplus \mathcal{J}_{n_{g}}\left(\lambda_{g}\right) \oplus \mathcal{J}_{n_{g+1}}\left(\lambda_{g+1}\right) \oplus \cdots \oplus \mathcal{J}_{n_{\nu}}\left(\lambda_{\nu}\right), \tag{2.11}
\end{equation*}
$$

where $\lambda_{1}=\cdots=\lambda_{g}=: \widehat{\lambda} \in \mathbb{C}, \lambda_{g+1}, \ldots, \lambda_{\nu} \in \mathbb{C} \backslash\{\widehat{\lambda}\}, n_{1} \geq \cdots \geq n_{g}$. Moreover, let $B=u v^{T}$, where

$$
u=\left[\begin{array}{c}
u_{1} \\
\vdots \\
u_{\nu}
\end{array}\right], \quad v=\left[\begin{array}{c}
v_{1} \\
\vdots \\
v_{\nu}
\end{array}\right], \quad u_{i}, v_{i} \in \mathbb{C}^{n_{i}}, i=1, \ldots, \nu .
$$

Assume that the first component of each vector $v_{i}, i=1, \ldots, \nu$ is nonzero. Then the matrix

$$
S:=\left(\operatorname{Toep}\left(v_{1}\right) \oplus \cdots \oplus \operatorname{Toep}\left(v_{\nu}\right)\right)^{-1}
$$

exists and satisfies $S^{-1} A S=A$ and

$$
\begin{equation*}
S^{-1} B S=\left[w e_{1, n_{1}}^{T}, \ldots, w e_{1, n_{\nu}}^{T}\right] \tag{2.12}
\end{equation*}
$$

where $w=S^{-1} u$. Moreover, the matrix $S^{-1}(A+B) S$ has at least $g-1$ Jordan chains associated with $\widehat{\lambda}$ of lengths at least $n_{2}, \ldots, n_{g}$ given by

$$
\begin{array}{lll}
e_{1}-e_{n_{1}+1}, & \cdots, & e_{n_{2}}-e_{n_{1}+n_{2}} \\
e_{1}-e_{n_{1}+n_{2}+1}, & \cdots, & e_{n_{3}}-e_{n_{1}+n_{2}+n_{3}} \\
\vdots & \ddots & \vdots  \tag{2.13}\\
e_{1}-e_{n_{1}+\cdots+n_{g-1}+1}, & \cdots, & e_{n_{g}}-e_{n_{1}+\cdots+n_{g-1}+n_{g}} .
\end{array}
$$

Proof. Clearly Toep $\left(v_{i}\right)$ is invertible if the first component of $v_{i}$ is nonzero, so $S$ exists. Moreover, $S$ commutes with $A$, and $e_{1, n_{i}}^{T}\left(\operatorname{Toep}\left(v_{i}\right)\right)=v_{i}^{T}$, so we have

$$
S^{-1} B S=S^{-1} u v^{T} S=\left[w e_{1, n_{1}}^{T}, \ldots, w e_{1, n_{\nu}}^{T}\right]
$$

It is then straightforward to check that the given chains are indeed Jordan chains associated with $\widehat{\lambda}$.

We emphasize that in Theorem 2.8 there is no claim whether the Jordan chains (2.13) associated with $\widehat{\lambda}$ can be extended to a longer chain or not, nor is there a claim whether (2.13) form a full basis of the corresponding root subspace or not.

Example 2.9 If $\widehat{\lambda}=0, \nu=g=3, n_{1}=4, n_{2}=3, n_{3}=2$, then the Brunovsky form of $A+B$ and the corresponding Jordan chains associated with $\widehat{\lambda}=0$ of length 3 and 2 are given by

$$
\left[\begin{array}{cccc|ccc|cc}
w_{1} & 1 & 0 & 0 & w_{1} & 0 & 0 & w_{1} & 0 \\
w_{2} & 0 & 1 & 0 & w_{2} & 0 & 0 & w_{2} & 0 \\
w_{3} & 0 & 0 & 1 & w_{3} & 0 & 0 & w_{3} & 0 \\
w_{4} & 0 & 0 & 0 & w_{4} & 0 & 0 & w_{4} & 0 \\
\hline w_{5} & 0 & 0 & 0 & w_{5} & 1 & 0 & w_{5} & 0 \\
w_{6} & 0 & 0 & 0 & w_{6} & 0 & 1 & w_{6} & 0 \\
w_{7} & 0 & 0 & 0 & w_{7} & 0 & 0 & w_{7} & 0 \\
\hline w_{8} & 0 & 0 & 0 & w_{8} & 0 & 0 & w_{8} & 1 \\
w_{9} & 0 & 0 & 0 & w_{9} & 0 & 0 & w_{9} & 0
\end{array}\right], \quad\left[\begin{array}{c}
1 \\
0 \\
0 \\
0 \\
-1 \\
0 \\
0 \\
0 \\
0
\end{array}\right],\left[\begin{array}{c}
0 \\
1 \\
0 \\
0 \\
0 \\
-1 \\
0 \\
0 \\
0
\end{array}\right],\left[\begin{array}{c}
0 \\
0 \\
1 \\
0 \\
0 \\
0 \\
-1 \\
0 \\
0
\end{array}\right], \quad\left[\begin{array}{c}
1 \\
0 \\
0 \\
0 \\
0 \\
0 \\
0 \\
-1 \\
0
\end{array}\right],\left[\begin{array}{c}
0 \\
1 \\
0 \\
0 \\
0 \\
0 \\
0 \\
0 \\
-1
\end{array}\right] .
$$

In the following, we want to apply Theorem 2.8 to the canonical forms in Section 2.3 which are close to but not quite in Jordan canonical form. Therefore, we will introduce the so called partial Brunovsky form with respect to a particular eigenvalue $\widehat{\lambda}$. With this form, the characteristic polynomial associated with the eigenvalue $\hat{\lambda}$ can be conveniently characterized.

Theorem 2.10 (Partial Brunovsky form) Let $A \in \mathbb{C}^{n \times n}$ be a matrix in the form

$$
\begin{equation*}
A=\mathcal{J}_{n_{1}}(\widehat{\lambda}) \oplus \cdots \oplus \mathcal{J}_{n_{g}}(\widehat{\lambda}) \oplus \widetilde{A} \tag{2.14}
\end{equation*}
$$

where $\widehat{\lambda} \in \mathbb{C}, \sigma(\widetilde{A}) \subseteq \mathbb{C} \backslash\{\widehat{\lambda}\}, n_{1} \geq \cdots \geq n_{g}$. Assume that $\widetilde{A}$ has the Jordan canonical form

$$
\mathcal{J}_{n_{g+1}}\left(\lambda_{g+1}\right) \oplus \cdots \oplus \mathcal{J}_{n_{\nu}}\left(\lambda_{\nu}\right)
$$

where $\lambda_{g+1}, \ldots, \lambda_{\nu}$ are not necessarily pairwise distinct. Moreover, let $B=u v^{T}$, where $u, v=\left[v_{1}^{T}, \ldots, v_{g}^{T}, \widetilde{v}^{T}\right]^{T} \in \mathbb{C}^{n}, v_{i} \in \mathbb{C}^{n_{i}}, i=1, \ldots, g$. Assume that the first component of each vector $v_{i}, i=1, \ldots, g$ is nonzero. Then the following statements hold:
(1) The matrix $S:=\left(\operatorname{Toep}\left(v_{1}\right) \oplus \cdots \oplus \operatorname{Toep}\left(v_{g}\right)\right)^{-1} \oplus I \in \mathbb{C}^{n \times n}$ exists and satisfies

$$
\begin{equation*}
S^{-1} A S=A \quad \text { and } \quad S^{-1} B S=\left[w e_{1, n_{1}}^{T}, \ldots, w e_{1, n_{g}}^{T}, z^{T}\right] \tag{2.15}
\end{equation*}
$$

where $w=S^{-1} u$ and for some appropriate vector $z$.
(2) The matrix $S^{-1}(A+B) S$ has at least $g-1$ Jordan chains associated with $\hat{\lambda}$ of lengths at least $n_{2}, \ldots, n_{g}$ given by (2.13).
(3) Partition $w=S^{-1} u$ as

$$
w=\left[\begin{array}{c}
w_{1} \\
\vdots \\
w_{g} \\
\widetilde{w}
\end{array}\right], \quad w_{i}=\left[\begin{array}{c}
w_{i, 1} \\
\vdots \\
w_{i, n_{i}}
\end{array}\right] \in \mathbb{C}^{n_{i}}, \quad i=1, \ldots, g
$$

let $a=n_{1}+\cdots+n_{g}$ denote the algebraic multiplicity of the eigenvalue $\hat{\lambda}$ of $A$, and let $\lambda_{1}, \ldots, \lambda_{n-a}$ denote the $n-a$ eigenvalues of $\widetilde{A}$ (counted with multiplicities). Set $\mu_{i}=\lambda_{i}-\widehat{\lambda}, i=1, \ldots, n-a$.
Then the characteristic polynomial $p_{\hat{\lambda}}(\lambda)$ of $A+B-\hat{\lambda} I$ has the form

$$
\begin{equation*}
p_{\hat{\lambda}}(\lambda)=\lambda^{a} q(\lambda)+\left(\prod_{i=g+1}^{\nu}\left(\mu_{i}-\lambda\right)^{n_{i}}\right) \cdot\left((-\lambda)^{a}+(-1)^{a-1} \sum_{i=1}^{g} \sum_{k=1}^{n_{i}} w_{i, k} \lambda^{a-k}\right) \tag{2.16}
\end{equation*}
$$

for some polynomial $q(\lambda)$.
Proof. Part (1) and (2) follow exactly as in Theorem 2.8. Concerning part (3), we may assume without loss of generality that $A+B$ is in Brunovsky form. Indeed, all that is needed is another similarity transformation with a matrix of the form $I_{a} \oplus \widetilde{S}^{-1} \in \mathbb{C}^{n \times n}$ which leaves the vectors $w_{1}, \ldots, w_{g}$ invariant in $\left(I_{a} \oplus \widetilde{S}^{-1}\right) w$.

With $A+B$ also $A+B-\widehat{\lambda} I_{n}$ is in Brunovsky form and the list of its diagonal elements is given by $\left(0, \ldots, 0, \mu_{g+1}, \ldots, \mu_{\nu}\right)$. Let $M$ denote the matrix that is obtained from $(A+B-\widehat{\lambda} I)-\lambda I$ by subtracting the first column from the columns $n_{1}+1$, $n_{1}+n_{2}+1, \ldots, n_{1}+n_{2}+\cdots+n_{\nu-1}+1$. Note that the column $n_{1}+\cdots+n_{i}+1$ then becomes zero except for $\lambda$ in the first entry, for $-\lambda$ in the $\left(n_{1}+\cdots+n_{i}+1\right)$-st entry if $i=1,2, \ldots, g-1$, and for $\mu_{i}-\lambda$ in the $\left(n_{1}+\cdots+n_{i}+1\right)$-st entry if $i=g, \ldots, \nu-1$. Then clearly $p_{\hat{\lambda}}(\lambda)=\operatorname{det} M$. If $\nu>g$, then partition $M$ as

$$
M=\left[\begin{array}{cccccc}
w_{1,1}-\lambda & z^{T} & \lambda & 0 & \cdots & 0 \\
\widetilde{w} & T & 0 & 0 & \cdots & 0 \\
w_{\nu, 1} & 0 & \mu_{\nu}-\lambda & 1 & & 0 \\
w_{\nu, 2} & 0 & 0 & \mu_{\nu}-\lambda & \ddots & 0 \\
\vdots & \vdots & \vdots & \ddots & \ddots & 1 \\
w_{\nu, n_{\nu}} & 0 & 0 & \cdots & 0 & \mu_{\nu}-\lambda
\end{array}\right]
$$

where $T \in \mathbb{C}^{\left(n-n_{\nu}-1\right) \times\left(n-n_{\nu}-1\right)}, n=n_{1}+\ldots+n_{\nu}$, is an upper triangular matrix whose first $a-1$ diagonal elements are equal to $-\lambda$. Thus, applying Laplace expansion successively, we obtain that

$$
\operatorname{det} M=\lambda^{a} \widetilde{q}(\lambda)+\left(\mu_{\nu}-\lambda\right)^{n_{\nu}} \operatorname{det}\left[\begin{array}{cc}
w_{1,1}-\lambda & z^{T}  \tag{2.17}\\
\widetilde{w} & T
\end{array}\right]
$$

for some polynomial $\widetilde{q}(\lambda)$. Indeed, for $n_{\nu}=1$ this is obvious and for $n_{\nu}>1$ we obtain

$$
\begin{aligned}
\operatorname{det} M & =(-1)^{n-1} w_{\nu, n_{\nu}} \cdot \underbrace{1 \cdot \ldots \cdot 1}_{n_{\nu} \text { times }} \cdot(-1)^{n-n_{\nu}-1} \lambda \operatorname{det} T+\left(\mu_{\nu}-\lambda\right) \operatorname{det} M_{n-1} \\
& =\lambda^{a} \breve{q}(\lambda)+\left(\mu_{\nu}-\lambda\right) \operatorname{det} M_{n-1},
\end{aligned}
$$

where $\breve{q}(\lambda)$ is some polynomial and $M_{n-1}$ is the principal $(n-1) \times(n-1)$ submatrix of $M$. Note that $M_{n-1}$ has the same structure as $M$ just with $n_{\nu}$ replaced with $n_{\nu-1}$. The claim then follows by induction. By further induction, we then obtain from (2.17) that

$$
\begin{equation*}
\operatorname{det} M=\lambda^{a} q(\lambda)+\left(\prod_{i=g+1}^{\nu}\left(\mu_{i}-\lambda\right)^{n_{i}}\right) \operatorname{det} M_{a} \tag{2.18}
\end{equation*}
$$

where $M_{a}$ is the principal $a \times a$ submatrix of $M$. It remains to compute the determinant of $M_{a}$. To this end, partition $M_{a}$ as

$$
M_{a}=\left[\begin{array}{cccccc}
w_{1,1}-\lambda & \widehat{z}^{T} & \lambda & 0 & \ldots & 0 \\
\widehat{w} & \widehat{T} & 0 & 0 & \ldots & 0 \\
w_{g, 1} & 0 & -\lambda & 1 & & 0 \\
w_{g, 2} & 0 & 0 & -\lambda & \ddots & 0 \\
\vdots & \vdots & \vdots & \ddots & \ddots & 1 \\
w_{g, n_{g}} & 0 & 0 & \ldots & 0 & -\lambda
\end{array}\right]
$$

Applying the cofactor expansion of $\operatorname{det} M_{a}$ by the first column, and using $\operatorname{det} \widehat{T}=$ $(-\lambda)^{a-n_{g}-1}$, we obtain that

$$
\begin{aligned}
& \operatorname{det} M_{a} \\
= & (-1)^{a-1} w_{g, n_{g}}(-1)^{a-n_{g}-1} \lambda \operatorname{det} \widehat{T}+(-\lambda)(-1)^{a-2} w_{g, n_{g}-1}(-1)^{a-n_{g}-1} \lambda \operatorname{det} \widehat{T} \\
& +\cdots+(-\lambda)^{n_{g}-1}(-1)^{a-n_{g}} w_{g, 1}(-1)^{a-n_{g}-1} \lambda \operatorname{det} \widehat{T}+(-\lambda)^{n_{g}} \operatorname{det}\left[\begin{array}{cc}
w_{1,1}-\lambda & \widehat{z}^{T} \\
\widehat{w} & \widehat{T}
\end{array}\right] \\
= & (-1)^{a-1}\left(\sum_{k=1}^{n_{g}} w_{g, k} \lambda^{a-k}\right)+(-\lambda)^{n_{g}} \operatorname{det}\left[\begin{array}{cc}
w_{1,1}-\lambda & \widehat{z}^{T} \\
\widehat{w} & \widehat{T}
\end{array}\right] .
\end{aligned}
$$

By induction, we finally obtain

$$
\operatorname{det} M_{a}=(-\lambda)^{a}+(-1)^{a-1} \sum_{i=1}^{g} \sum_{k=1}^{n_{i}} w_{i, k} \lambda^{a-k}
$$

where the extra term $(-\lambda)^{a}$ appears due to the fact that the first entry of the first column of $M_{a}$ is not $w_{1,1}$, but $w_{1,1}-\lambda$.

In the next section, we will need explicit formulas for some of the coefficients of the characteristic polynomial of $A+B$ with $A, B$ as in Theorem 2.10. We establish those in the following corollary which follows immediately from Theorem 2.10. For the ease of future reference, we group together Jordan blocks of the same size in the Jordan canonical form of $A$ and repeat part of Theorem 2.10.

Corollary 2.11 Let

$$
\begin{equation*}
A=\left(\bigoplus_{j=1}^{\ell_{1}} \mathcal{J}_{n_{1}}(\widehat{\lambda})\right) \oplus \cdots \oplus\left(\bigoplus_{j=1}^{\ell_{m}} \mathcal{J}_{n_{m}}(\widehat{\lambda})\right) \oplus \widetilde{A} \in \mathbb{C}^{n \times n} \tag{2.19}
\end{equation*}
$$

where $n_{1}>\cdots>n_{m}$ and $\sigma(\widetilde{A}) \subseteq \mathbb{C} \backslash\{\widehat{\lambda}\}$. Moreover, let $a=\ell_{1} n_{1}+\cdots+\ell_{m} n_{m}$ denote the algebraic multiplicity of $\widehat{\lambda}$ and let $B=u v^{T}$, where $u \in \mathbb{C}^{n}$,

$$
v=\left[\begin{array}{c}
v^{(1)} \\
\vdots \\
v^{(m)} \\
\widetilde{v}
\end{array}\right], \quad v^{(i)}=\left[\begin{array}{c}
v^{(i, 1)} \\
\vdots \\
v^{\left(i, \ell_{i}\right)}
\end{array}\right], \quad v^{(i, j)} \in \mathbb{C}^{n_{i}}, \quad j=1, \ldots, \ell_{i}, \quad i=1, \ldots, m .
$$

Assume that the first component of each vector $v^{(i, j)}, j=1, \ldots, \ell_{i}, i=1, \ldots, m$ is nonzero. Then the following statements hold:
(1) The matrix $S:=\left(\bigoplus_{j=1}^{\ell_{1}} \operatorname{Toep}\left(v^{(1, j}\right) \oplus \cdots \oplus \bigoplus_{j=1}^{\ell_{m}} \operatorname{Toep}\left(v^{(m, j)}\right)^{-1} \oplus I_{n-a}\right.$ exists and satisfies

$$
S^{-1} A S=A, \quad S^{-1} B S=w[\underbrace{e_{1, n_{1}}^{T}, \ldots, e_{1, n_{1}}^{T}}_{\ell_{1} \text { times }}, \ldots, \underbrace{e_{1, n_{m}}^{T}, \ldots, e_{1, n_{m}}^{T}}_{\ell_{m} \text { times }}, z^{T}]
$$

where $w=S^{-1} u$ and for some appropriate vector $z \in \mathbb{C}^{n-a}$.
(2) The matrix $S^{-1}(A+B) S$ has at least $\ell_{1}+\cdots+\ell_{m}-1$ Jordan chains associated with $\widehat{\lambda}$ given as follows:
a) $\ell_{1}-1$ Jordan chains of length at least $n_{1}$ :

$$
\begin{array}{lll}
e_{1}-e_{n_{1}+1}, & \ldots, & e_{n_{1}}-e_{2 n_{1}} \\
\vdots & \ddots & \vdots  \tag{2.20}\\
e_{1}-e_{\left(\ell_{1}-1\right) n_{1}+1}, & \ldots, & e_{n_{1}}-e_{\ell_{1} n_{1}}
\end{array}
$$

b) $\ell_{i}$ Jordan chains of length at least $n_{i}$ for $i=2, \ldots, m$ :

$$
\begin{array}{lll}
e_{1}-e_{\ell_{1} n_{1}+\cdots+\ell_{i-1} n_{i-1}+1}, & \cdots, & e_{n_{i}}-e_{\ell_{1} n_{1}+\cdots+\ell_{i-1} n_{i-1}+n_{i}} \\
e_{1}-e_{\ell_{1} n_{1}+\cdots+\ell_{i-1} n_{i-1}+n_{i}+1}, & \cdots, & e_{n_{i}}-e_{\ell_{1} n_{1}+\cdots+\ell_{i-1} n_{i-1}+2 n_{i}} \\
\vdots & \ddots & \vdots  \tag{2.21}\\
e_{1}-e_{\ell_{1} n_{1}+\cdots+\ell_{i-1} n_{i-1}+\left(\ell_{i}-1\right) n_{i}+1}, & \cdots, & e_{n_{i}}-e_{\ell_{1} n_{1}+\cdots+\ell_{i-1} n_{i-1}+\ell_{i} n_{i}}
\end{array}
$$

(3) Partition $w=S^{-1} u$ as

$$
w=\left[\begin{array}{c}
w^{(1)} \\
\vdots \\
w^{(m)} \\
\widetilde{w}
\end{array}\right], \quad w^{(i)}=\left[\begin{array}{c}
w^{(i, 1)} \\
\vdots \\
w^{\left(i, \ell_{i}\right)}
\end{array}\right], \quad w^{(i, j)}=\left[\begin{array}{c}
w_{1}^{(i, j)} \\
\vdots \\
w_{n_{i}}^{(i, j)}
\end{array}\right] \in \mathbb{C}^{n_{i}}
$$

and let $\lambda_{1}, \ldots, \lambda_{q}$ be the pairwise distinct eigenvalues of $A$ different from $\hat{\lambda}$ having the algebraic multiplicities $r_{1}, \ldots, r_{q}$, respectively. Set $\mu_{i}=\lambda_{i}-\widehat{\lambda}, i=1,2, \ldots, q$. Then the characteristic polynomial $p_{\hat{\lambda}}$ of $A+B-\widehat{\lambda} I$ has the form

$$
p_{\hat{\lambda}}(\lambda)=c_{n} \lambda^{n}+\cdots+c_{a-n_{1}+1} \lambda^{a-n_{1}+1}+c_{a-n_{1}} \lambda^{a-n_{1}}
$$

where

$$
\begin{equation*}
c_{a-n_{1}}=(-1)^{a}\left(\prod_{i=1}^{q} \mu_{i}^{r_{i}}\right)\left(\sum_{j=1}^{\ell_{1}} w_{n_{1}}^{(1, j)}\right) \tag{2.22}
\end{equation*}
$$

and

$$
\begin{equation*}
c_{a-n_{1}+1}=-\left(\sum_{\nu=1}^{q} r_{\nu} \mu_{\nu}^{r_{\nu}-1} \prod_{\substack{i=1 \\ i \neq \nu}}^{q} \mu_{i}^{r_{i}}\right)\left(\sum_{j=1}^{\ell_{1}} w_{n_{1}}^{(1, j)}\right)+\left(\prod_{i=1}^{q} \mu_{i}^{r_{i}}\right)\left(\sum_{j=1}^{\ell_{1}} w_{n_{1}-1}^{(1, j)}\right) \tag{2.23}
\end{equation*}
$$

if $n_{1}-1>n_{2}$ or, if $n_{1}-1=n_{2}$, then

$$
\begin{align*}
c_{a-n_{1}+1} & =-\left(\sum_{\nu=1}^{q} r_{\nu} \mu_{\nu}^{r_{\nu}-1} \prod_{\substack{i=1 \\
i \neq \nu}}^{q} \mu_{i}^{r_{i}}\right)\left(\sum_{j=1}^{\ell_{1}} w_{n_{1}}^{(1, j)}\right) \\
& +\left(\prod_{i=1}^{q} \mu_{i}^{r_{i}}\right)\left(\sum_{j=1}^{\ell_{1}} w_{n_{1}-1}^{(1, j)}+\sum_{j=1}^{\ell_{2}} w_{n_{2}}^{(2, j)}\right) . \tag{2.24}
\end{align*}
$$

Proof. The parts (1) and (2) follow immediately from Theorem 2.10. Moreover, the characteristic polynomial of $A+B$ is given by

$$
p_{0}(\lambda)=(-\lambda)^{a} q(\lambda)+\left(\prod_{i=1}^{q}\left(\mu_{i}-\lambda\right)^{r_{i}}\right) \cdot\left((-\lambda)^{a}+(-1)^{a-1} \sum_{i=1}^{m} \sum_{j=1}^{\ell_{i}} \sum_{k=1}^{n_{i}} w_{k}^{(i, j)} \lambda^{a-k}\right),
$$

where $q(\lambda)$ is some polynomial. The lowest possible power of $\lambda$ associated with a nonzero coefficient in $p_{0}(\lambda)$ is clearly $a-n_{1}$ and the corresponding coefficient $c_{a-n_{1}}$ is

$$
c_{a-n_{1}}=(-1)^{a}\left(\prod_{i=1}^{q} \mu_{i}^{r_{i}}\right)\left(\sum_{j=1}^{\ell_{1}} w_{n_{1}}^{(1, j)}\right)
$$

while the coefficient $c_{a-n_{1}+1}$ of $\lambda^{a-n_{1}+1}$ in $p_{0}(\lambda)$ is as in (2.23) or (2.24) depending on whether $n_{1}-1>n_{2}$ or $n_{1}-1=n_{2}$.

### 2.5 Extension of Jordan chains

In this section, we discuss the extension of some Jordan chains of a matrix in Brunovsky form. This will be needed to prove our main results. However, it is not always possible to extend a given set of Jordan chains to a set of Jordan chains that forms a basis as we will illustrate in the following example.

Example 2.12 Consider the rank one perturbation of $A=\mathcal{J}_{2}(0) \oplus \mathcal{J}_{2}(0) \oplus \mathcal{J}_{1}(0)$ given by

$$
\widetilde{A}=\left[\begin{array}{rrrrr}
a & 1 & a & 0 & a \\
b & 0 & b & 0 & b \\
c & 0 & c & 1 & c \\
-b & 0 & -b & 0 & -b \\
d & 0 & d & 0 & d
\end{array}\right], \quad b, d \neq 0
$$

which is obviously in Brunovsky form. By Theorem 2.8 we know that $\widetilde{A}$ has at least two Jordan chains associated with 0 of lengths at least 2 and 1 , given by

$$
\begin{equation*}
e_{1}-e_{3}, e_{2}-e_{4}, \quad \text { and } \quad e_{1}-e_{5}, \tag{2.25}
\end{equation*}
$$

respectively. Let us check whether the first chain can be extended to a Jordan chain of length three. For this, we would have to show that $e_{2}-e_{4}$ is in the range of $\widetilde{A}$. However, the linear system

$$
\left[\begin{array}{r}
0 \\
1 \\
0 \\
-1 \\
0
\end{array}\right]=\left[\begin{array}{rrrrr}
a & 1 & a & 0 & a \\
b & 0 & b & 0 & b \\
c & 0 & c & 1 & c \\
-b & 0 & -b & 0 & -b \\
d & 0 & d & 0 & d
\end{array}\right]\left[\begin{array}{c}
x_{1} \\
x_{2} \\
x_{3} \\
x_{4} \\
x_{5}
\end{array}\right]=\left[\begin{array}{c}
x_{2}+a\left(x_{1}+x_{3}+x_{5}\right) \\
b\left(x_{1}+x_{3}+x_{5}\right) \\
x_{4}+c\left(x_{1}+x_{3}+x_{5}\right) \\
-b\left(x_{1}+x_{3}+x_{5}\right) \\
d\left(x_{1}+x_{3}+x_{5}\right)
\end{array}\right]
$$

with unknowns $x_{1}, \ldots, x_{5}$ does not have a solution, because $d \neq 0$, so the chain $e_{1}-$ $e_{3}, e_{2}-e_{4}$ cannot be extended to a Jordan chains of length 3. Nevertheless, it can be shown that $\widetilde{A}$ does have a Jordan chain of length at least 3 associated with the eigenvalue zero. To this end, consider the vectors

$$
\left[\begin{array}{r}
1 \\
0 \\
-1 \\
0 \\
0
\end{array}\right],\left[\begin{array}{r}
\alpha \\
1 \\
0 \\
-1 \\
-\alpha
\end{array}\right]
$$

that form a Jordan chain of $\widetilde{A}$ associated with 0 of length 2 . We now show that this chain can be extended for a particular choice of $\alpha$. Indeed, for $\alpha=-d / b$ the linear
system

$$
\left[\begin{array}{r}
\alpha \\
1 \\
0 \\
-1 \\
-\alpha
\end{array}\right]=\left[\begin{array}{rrrrr}
a & 1 & a & 0 & a \\
b & 0 & b & 0 & b \\
c & 0 & c & 1 & c \\
-b & 0 & -b & 0 & -b \\
d & 0 & d & 0 & d
\end{array}\right]\left[\begin{array}{c}
x_{1} \\
x_{2} \\
x_{3} \\
x_{4} \\
x_{5}
\end{array}\right]=\left[\begin{array}{c}
x_{2}+a\left(x_{1}+x_{3}+x_{5}\right) \\
b\left(x_{1}+x_{3}+x_{5}\right) \\
x_{4}+c\left(x_{1}+x_{3}+x_{5}\right) \\
-b\left(x_{1}+x_{3}+x_{5}\right) \\
d\left(x_{1}+x_{3}+x_{5}\right)
\end{array}\right]
$$

has $x_{1}=1 / b, x_{2}=-d / b-a / b, x_{3}=0, x_{4}=-c / b, x_{5}=0$ as a solution.
Note that the Jordan chain that could be extended in Example 2.12 can be considered as a "linear combination" of the two Jordan chains in (2.25). We will need similar constructions later in this paper and therefore, we introduce the following "sum" of Jordan chains.

Definition 2.13 Let $A \in \mathbb{C}^{n \times n}$ and let $X=\left(x_{1}, \ldots, x_{p}\right)$ and $Y=\left(y_{1}, \ldots, y_{q}\right)$ be two Jordan chains of $A$ associated with the same eigenvalue $\widehat{\lambda}$ of (possibly different) lengths $p$ and $q$. Then the sum $X+Y$ of $X$ and $Y$ is defined to be the chain $Z=\left(z_{1}, \ldots, z_{\max (p, q)}\right)$, where

$$
z_{j}=\left\{\begin{array}{ll}
x_{j} & \text { if } p \geq q \\
y_{j} & \text { if } p<q
\end{array}, \quad j=1, \ldots,|p-q|\right.
$$

and

$$
z_{j}=\left\{\begin{array}{ll}
x_{j}+y_{j-p+q} & \text { if } p \geq q \\
y_{j}+x_{j-p+q} & \text { if } p<q
\end{array}, \quad j=|p-q|+1, \ldots, \max (p, q)\right. \text {. }
$$

To illustrate this construction, consider e.g., $X=\left(x_{1}, x_{2}, x_{3}, x_{4}\right)$ and $Y=\left(y_{1}, y_{2}\right)$ then $X+Y=\left(x_{1}, x_{2}, x_{3}+y_{1}, x_{4}+y_{2}\right)$.

It is straightforward to check that the sum $Z=X+Y$ of two Jordan chains associated with an eigenvalue $\widehat{\lambda}$ is again a Jordan chain associated with $\widehat{\lambda}$ of the given matrix $A$, but it should be noted that this sum is not commutative.

With these preliminary results, we have now set the stage to derive the desired perturbation theorems for structured matrices under generic rank one perturbations in the following sections.

## 3 Generic structured rank one perturbations for general classes of matrices with symmetries

In this section we state and prove general theorems concerning generic structured rank one perturbations. Although we focus on symmetry structures with respect to bilinear forms in this paper, the theorems cover a much wider class of structured matrices including matrices that are structured with respect to sesquilinear forms. To this end in the next two theorems, we will use the notation ${ }^{\star}$ to denote either the transpose ${ }^{T}$ or the conjugate transpose *.

Theorem 3.1 Let $A \in \mathbb{F}^{n \times n}$ and let $T, H \in \mathbb{F}^{n \times n}$ be invertible such that

$$
\begin{align*}
T^{-1} A T & =\left(\bigoplus_{j=1}^{\ell_{1}} \mathcal{J}_{n_{1}}(\widehat{\lambda})\right) \oplus\left(\bigoplus_{j=1}^{\ell_{2}} \mathcal{J}_{n_{2}}(\widehat{\lambda})\right) \oplus \cdots \oplus\left(\bigoplus_{j=1}^{\ell_{m}} \mathcal{J}_{n_{m}}(\widehat{\lambda})\right) \oplus \widetilde{A}  \tag{3.1}\\
T^{\star} H T & =\left(\bigoplus_{j=1}^{\ell_{1}} H^{(1, j)}\right) \oplus H^{(2)} \oplus \cdots \oplus H^{(m)} \oplus \widetilde{H} \tag{3.2}
\end{align*}
$$

where $\widehat{\lambda} \in \mathbb{F}$ and the decompositions (3.1) and (3.2) have the following properties:
(1) $n_{1}>n_{2}>\cdots>n_{m}$;
(2) $H^{(j)} \in \mathbb{F}^{\ell_{j} n_{j} \times \ell_{j} n_{j}}, j=2, \ldots, m$ and the matrices

$$
H^{(1, j)}=\left[\begin{array}{cccc}
0 & \ldots & 0 & h_{1, n_{1}}^{(1, j)} \\
\vdots & . \cdot & h_{2, n_{1}-1}^{(1, j)} & h_{2, n_{1}}^{(1, j)} \\
0 & . \cdot & . \cdot & \vdots \\
h_{n_{1}, 1}^{(1, j)} & h_{n_{1}, 2}^{(1, j)} & \ldots & h_{n_{1}, n_{1}}^{(1, j)}
\end{array}\right], \quad j=1,2, \ldots, \ell_{1} ;
$$

are anti-triangular (necessarily invertible);
(3) $\widetilde{H}, \widetilde{A} \in \mathbb{F}^{(n-a) \times(n-a)}$, where $a=\sum_{j=1}^{m} \ell_{j} n_{j}$ and $\sigma(\widetilde{A}) \subseteq \mathbb{C} \backslash\{\widehat{\lambda}\}$.

If $B \in \mathbb{F}^{n \times n}$ is a rank one matrix of the form $B=u u^{\star} H$, then generically (with respect to the components of $u$ if $\star=T$, and with respect to the real and imaginary parts of the components of $u$ if $\star=*) A+B$ has the Jordan canonical form

$$
\left(\bigoplus_{j=1}^{\ell_{1}-1} \mathcal{J}_{n_{1}}(\widehat{\lambda})\right) \oplus\left(\bigoplus_{j=1}^{\ell_{2}} \mathcal{J}_{n_{2}}(\widehat{\lambda})\right) \oplus \cdots \oplus\left(\bigoplus_{j=1}^{\ell_{m}} \mathcal{J}_{n_{m}}(\widehat{\lambda})\right) \oplus \widetilde{\mathcal{J}}
$$

where $\widetilde{\mathcal{J}}$ contains all the Jordan blocks of $A+B$ associated with eigenvalues different from $\widehat{\lambda}$.

Proof. Without loss of generality, let $A, H$ be in the forms (3.1) and (3.2) already. In view of Theorem 2.10 it is sufficient to show that the algebraic multiplicity of the eigenvalue $\hat{\lambda}$ of $A+B$ is $a-n_{1}$ generically. Let

$$
u=\left[\begin{array}{c}
u^{(1)} \\
\vdots \\
u^{(m)} \\
\widetilde{u}
\end{array}\right], \quad u^{(i)}=\left[\begin{array}{c}
u^{(i, 1)} \\
\vdots \\
u^{\left(i, \ell_{i}\right)}
\end{array}\right], \quad u^{(i, j)}=\left[\begin{array}{c}
u_{1}^{(i, j)} \\
\vdots \\
u_{n_{i}}^{(i, j)}
\end{array}\right] \in \mathbb{F}^{n_{i}}, \quad \widetilde{u} \in \mathbb{F}^{n-a}
$$

and

$$
v=H^{\star} u=\left[\begin{array}{c}
v^{(1)} \\
\vdots \\
v^{(m)} \\
\widetilde{v}
\end{array}\right], \quad v^{(i)}=\left[\begin{array}{c}
v^{(i, 1)} \\
\vdots \\
v^{\left(i, \ell_{i}\right)}
\end{array}\right], \quad v^{(i, j)}=\left[\begin{array}{c}
v_{1}^{(i, j)} \\
\vdots \\
v_{n_{i}}^{(i, j)}
\end{array}\right] \in \mathbb{F}^{n_{i}}, \quad \widetilde{v} \in \mathbb{F}^{n-a} .
$$

Generically (in the sense of the theorem), we have $v_{1}^{(i, j)} \neq 0$, because $H$ is invertible. In particular, we have

$$
v_{1}^{(1, j)}= \begin{cases}h_{n_{1}, 1}^{(1, j)} u_{n_{1}}^{(1, j)} & \text { if } \star=T,  \tag{3.3}\\ h_{n_{1}, 1}^{(1, j)} u_{n_{1}}^{(1, j)} & \text { if } \star=* .\end{cases}
$$

So by Theorem 2.10 we can compute $S^{-1}=\widehat{S} \oplus I_{n-a}$, where

$$
\begin{aligned}
& \widehat{S}=\bigoplus_{j=1}^{\ell_{1}} \operatorname{Toep}\left(v^{(1, j)}\right) \oplus \cdots \oplus \bigoplus_{j=1}^{\ell_{m}} \operatorname{Toep}\left(v^{(m, j)}\right) \in \mathbb{F}^{a \times a}, \quad \text { if } \star=T, \\
& \widehat{S}=\bigoplus_{j=1}^{\ell_{1}} \operatorname{Toep}\left(\overline{v^{(1, j)}}\right) \oplus \cdots \oplus \bigoplus_{j=1}^{\ell_{m}} \operatorname{Toep}\left(\overline{v^{(m, j)}}\right) \in \mathbb{C}^{a \times a}, \quad \text { if } \star=* .
\end{aligned}
$$

Thus, we obtain that

$$
S^{-1}(A+B) S=S^{-1}\left(A+u v^{\star}\right) S
$$

is in partial Brunovsky form (2.14) with respect to $\widehat{\lambda}$ and

$$
w:=S^{-1} u=\left[\begin{array}{c}
w^{(1)} \\
\vdots \\
w^{(m)} \\
\widetilde{w}
\end{array}\right], \quad w^{(i)}=\left[\begin{array}{c}
w^{(i, 1)} \\
\vdots \\
w^{\left(i, \ell_{i}\right)}
\end{array}\right], \quad w^{(i, j)}=\left[\begin{array}{c}
w_{1}^{(i, j)} \\
\vdots \\
w_{n_{i}}^{(i, j)}
\end{array}\right] \in \mathbb{F}^{n_{i}}, \quad \widetilde{w} \in \mathbb{F}^{n-a},
$$

where

$$
w_{n_{1}}^{(1, j)}=\left(v_{1}^{(1, j)}\right)^{\star} u_{n_{1}}^{(1, j)}= \begin{cases}h_{n_{1}, 1}^{(1, j)}\left(u_{n_{1}}^{(1, j)}\right)^{2} & \text { if } \star=T, \\ h_{n_{1}, 1}^{(1, j)}\left|u_{n_{1}}^{(1, j)}\right|^{2} & \text { if } \star=* .\end{cases}
$$

By Corollary 2.11, and taking into account formula (2.12), the characteristic polynomial of $A+B-\widehat{\lambda} I$ is given by

$$
p_{\hat{\lambda}}(\lambda)=\sum_{i=a-n_{1}}^{n} c_{i} \lambda^{i}
$$

where

$$
c_{a-n_{1}}=M \sum_{j=1}^{\ell_{1}} w_{n_{1}}^{(1, j)}= \begin{cases}\sum_{j=1}^{\ell_{1}} h_{n_{1}, 1}^{(1, j)}\left(u_{n_{1}}^{(1, j)}\right)^{2} & \text { if } \star=T, \\ \sum_{j=1}^{\ell_{1}} h_{n_{1}, 1}^{(1, j)} \mid\left(\left.u_{n_{1}}^{(1, j)}\right|^{2}\right. & \text { if } \star=* ;\end{cases}
$$

here $M \neq 0$ is a constant independent of $B$. Clearly, $c_{a-n_{1}}$ is generically (in the sense stated in the theorem) nonzero and hence the algebraic multiplicity of the eigenvalue $\widehat{\lambda}$ of $A+B$ is $a-n_{1}$. Together with Corollary 2.11, we obtain that the only possible Jordan canonical form for $A+B$ is given by

$$
\bigoplus_{j=1}^{\ell_{1}-1} \mathcal{J}_{n_{1}}(\widehat{\lambda}) \oplus \bigoplus_{j=1}^{\ell_{2}} \mathcal{J}_{n_{2}}(\widehat{\lambda}) \oplus \cdots \oplus \bigoplus_{j=1}^{\ell_{m}} \mathcal{J}_{n_{m}}(\widehat{\lambda}) \oplus \widetilde{\mathcal{J}}
$$

where $\widetilde{\mathcal{J}}$ contains all the Jordan blocks associated with eigenvalues of $A+B$ different from $\widehat{\lambda}$.

Theorem 3.2 Let $A \in \mathbb{F}^{n \times n}$ and let $T, H \in \mathbb{F}^{n \times n}$ be invertible matrices such that

$$
T^{-1} A T=\widehat{A} \oplus \breve{A} \oplus \widetilde{A}, \quad T^{\star} H T=\left[\begin{array}{cc}
0 & I_{a}  \tag{3.4}\\
\widehat{H} & 0
\end{array}\right] \oplus \widetilde{H},
$$

where the decomposition (3.4) has the following properties:
(a)

$$
\widehat{A}=\left(\bigoplus_{j=1}^{\ell_{1}} \mathcal{J}_{n_{1}}(\widehat{\lambda})\right) \oplus \cdots \oplus\left(\bigoplus_{j=1}^{\ell_{m}} \mathcal{J}_{n_{m}}(\widehat{\lambda})\right)
$$

where $n_{1}>n_{2}>\cdots>n_{m}$ and $\hat{\lambda} \in \mathbb{F}$;
(b) $a=\sum_{j=1}^{m} \ell_{j} n_{j}$ and $\widehat{H}, \breve{A} \in \mathbb{F}^{a \times a}, \widetilde{H} \in \mathbb{F}^{(n-2 a) \times(n-2 a)}$;
(c) $\sigma(\breve{A}), \sigma(\widetilde{A}) \subseteq \mathbb{C} \backslash\{\widehat{\lambda}\}$.

If $B \in \mathbb{F}^{n \times n}$ is a rank one perturbation of the form $B=u u^{\star} H, u \in \mathbb{F}^{n}$, then generically (with respect to the components of $u$ if $\star=T$, and with respect to the real and imaginary parts of the components of $u$ if $\star=*) A+B$ has the Jordan canonical form

$$
\left(\bigoplus_{j=1}^{\ell_{1}-1} \mathcal{J}_{n_{1}}(\widehat{\lambda})\right) \oplus\left(\bigoplus_{j=1}^{\ell_{2}} \mathcal{J}_{n_{2}}(\widehat{\lambda})\right) \oplus \cdots \oplus\left(\bigoplus_{j=1}^{\ell_{m}} \mathcal{J}_{n_{m}}(\widehat{\lambda})\right) \oplus \widetilde{\mathcal{J}}
$$

where $\mathcal{J}$ contains all the Jordan blocks of $A+B$ associated with eigenvalues different from $\widehat{\lambda}$.

Note that $\widehat{H}, \widetilde{H}$ are necessarily invertible.
Proof. As in the proof of Theorem 3.1, we may assume that $A$ and $H$ are in the forms (3.4). Partition

$$
u=\left[\begin{array}{c}
\widehat{u} \\
\breve{u} \\
\widetilde{u}
\end{array}\right], \quad \widehat{u}=\left[\begin{array}{c}
\widehat{u}^{(1)} \\
\vdots \\
\widehat{u}^{(m)}
\end{array}\right], \quad \widehat{u}^{(i)}=\left[\begin{array}{c}
\widehat{u}^{(i, 1)} \\
\vdots \\
\widehat{u}^{\left(i, \ell_{i}\right)}
\end{array}\right], \quad \widehat{u}^{(i, j)}=\left[\begin{array}{c}
\widehat{u}_{1}^{(i, j)} \\
\vdots \\
\widehat{u}_{n_{i}}^{(i, j)}
\end{array}\right] \in \mathbb{F}^{n_{i}}
$$

and

$$
\breve{u}=\left[\begin{array}{c}
\breve{u}^{(1)} \\
\vdots \\
\breve{u}^{(m)}
\end{array}\right], \quad \breve{u}^{(i)}=\left[\begin{array}{c}
\breve{u}^{(i, 1)} \\
\vdots \\
\breve{u}^{\left(i, \ell_{i}\right)}
\end{array}\right], \quad \breve{u}^{(i, j)}=\left[\begin{array}{c}
\breve{u}_{1}^{(i, j)} \\
\vdots \\
\breve{u}_{n_{i}}^{(i, j)}
\end{array}\right] \in \mathbb{F}^{n_{i}} .
$$

Observe that the vector $v=H^{\star} u$ has the form

$$
v=\left[\begin{array}{c}
\breve{u} \\
\widehat{H} \widehat{u} \\
\widetilde{H} \widetilde{u}
\end{array}\right] .
$$

Generically (in the sense of the theorem), we can now form the matrix $S=\widehat{S} \oplus I_{n-a}$, where

$$
\begin{aligned}
& \widehat{S}^{-1}=\left(\bigoplus_{j=1}^{\ell_{1}} \operatorname{Toep}\left(\breve{u}^{(1, j)}\right)\right) \oplus \cdots \oplus\left(\bigoplus_{j=1}^{\ell_{m}} \operatorname{Toep}\left(\breve{u}^{(m, j)}\right)\right) \quad \text { if } \star=T, \\
& \widehat{S}^{-1}=\left(\bigoplus_{j=1}^{\ell_{1}} \operatorname{Toep}\left(\overline{\breve{u}^{(1, j)}}\right)\right) \oplus \cdots \oplus\left(\bigoplus_{j=1}^{\ell_{m}} \operatorname{Toep}\left(\overline{\left.\breve{u}^{(m, j)}\right)}\right) \quad \text { if } \star=* .\right.
\end{aligned}
$$

Then $S^{-1}(A+B) S=S^{-1}\left(A+u v^{\star}\right) S$ is in partial Brunovsky form (2.14) as in Theorem 2.10. Next, consider the vector

$$
w:=S^{-1} u=\left[\begin{array}{c}
\widehat{w} \\
\breve{w} \\
\widetilde{w}
\end{array}\right], \widehat{w}=\left[\begin{array}{c}
\widehat{w}^{(1)} \\
\vdots \\
\widehat{w}^{(m)}
\end{array}\right], \widehat{w}^{(i)}=\left[\begin{array}{c}
\widehat{w}^{(i, 1)} \\
\vdots \\
\widehat{w}^{\left(i, \ell_{i}\right)}
\end{array}\right], \widehat{w}^{(i, j)}=\left[\begin{array}{c}
\widehat{w}_{1}^{(i, j)} \\
\vdots \\
\widehat{w}_{n_{i}}^{(i, j)}
\end{array}\right] \in \mathbb{F}^{n_{i}} .
$$

Then we obtain

$$
\widehat{w}_{n_{i}}^{(i, j)}=\left(\breve{u}_{1}^{(i, j)}\right)^{\star} \widehat{u}_{n_{i}}^{(i, j)} .
$$

By Corollary 2.11, the characteristic polynomial $p_{\hat{\lambda}}$ of $A+B-\widehat{\lambda} I$ has the form

$$
p_{\widehat{\lambda}}(\lambda)=c_{n} \lambda^{n}+\cdots+c_{a-n_{1}+1} \lambda^{a-n_{1}+1}+c_{a-n_{1}} \lambda^{a-n_{1}}
$$

where

$$
c_{a-n_{1}}=M \cdot\left(\sum_{j=1}^{\ell_{1}} \widehat{w}_{n_{1}}^{(1, j)}\right)=M\left(\sum_{j=1}^{\ell_{1}}\left(\breve{u}_{1}^{(1, j)}\right)^{\star T} \widehat{u}_{n_{1}}^{(1, j)}\right) ;
$$

$M \neq 0$ is a constant independent of $B$. Clearly, $c_{a-n_{1}}$ is generically (in the sense indicated in the statement of Theorem 3.2) nonzero and thus $\widetilde{a}=a-n_{1}$ is generically the algebraic multiplicity of the eigenvalue $\hat{\lambda}$ of $A+B$. Together with Corollary 2.11, it follows that the only possible Jordan canonical forms for $A+B$ are

$$
\left(\bigoplus_{j=1}^{\ell_{1}-1} \mathcal{J}_{n_{1}}(\widehat{\lambda})\right) \oplus\left(\bigoplus_{j=1}^{\ell_{2}} \mathcal{J}_{n_{2}}(\widehat{\lambda})\right) \oplus \cdots \oplus\left(\bigoplus_{j=1}^{\ell_{m}} \mathcal{J}_{n_{m}}(\widehat{\lambda})\right) \oplus \widetilde{\mathcal{J}}
$$

where $\widetilde{\mathcal{J}}$ contains all the Jordan blocks associated with eigenvalues different from $\widehat{\lambda}$.

Note that the scenario in Theorems 3.1 and 3.2 corresponds exactly to the scenario under arbitrary unstructured rank one perturbations; cf. Theorem 2.3.

The two theorems apply to the cases of symmetric complex matrix $H$ and $H$ symmetric matrices (see Theorem 2.6) discussed in Section 5, and also to the case of $H$-selfadjoint matrices that will be discussed in a forthcoming paper. Finally, they apply to the case where $J$ is skew-symmetric, and $A$ is $J$-Hamiltonian and invertible (case (iii) in Theorem 2.7). Thus, for $J$-Hamiltonian matrices it remains to study the case of the eigenvalue zero. This will be done in the next section.

## 4 Generic structured rank one perturbations for complex Hamiltonian matrices

In this section we state and prove one of the main results of the paper concerning perturbations of complex Hamiltonian matrices. According to Theorem 2.7, if $\lambda \neq 0$ is an eigenvalue of a complex $J$-Hamiltonian matrix $A$, then so is $-\lambda$ (with the same partial multiplicities), and for every odd $k$, the number of Jordan blocks in the Jordan form of $A$ of size $k$ corresponding to the zero eigenvalue is even.

As Theorem 4.2 shows, in the case the largest partial multiplicity of the zero eigenvalue is odd, the generic behavior of the Jordan structure of the perturbed matrix contrasts sharply with the unstructured situation (Theorem 2.3). To motivate the main result, consider an example first:

Example 4.1 Consider the matrix

$$
Z(w)=\left[\begin{array}{cc}
\mathcal{J}_{2 m+1}(0) & 0 \\
0 & \mathcal{J}_{2 m+1}(0)
\end{array}\right]+w w^{T}\left[\begin{array}{cc}
0 & \Sigma_{2 m+1} \\
-\Sigma_{2 m+1} & 0
\end{array}\right] \in \mathbb{C}^{(4 m+2) \times(4 m+2)}
$$

We will show that generically (with respect to the components of $w \in \mathbb{C}^{2 m+1}$ ) $Z(w)$ has the Jordan from of type $\mathcal{J}_{2 m+2}(0) \oplus k_{1} \oplus k_{2} \oplus \cdots \oplus k_{2 m}$, where the $k_{j}$ 's are distinct nonzero complex numbers.

A standard transformation allows us to consider the matrix

$$
M:=M(u, v):=\left[\begin{array}{cc}
\mathcal{J}_{2 m+1}(0) & 0 \\
0 & -\mathcal{J}_{2 m+1}(0)^{T}
\end{array}\right]+\left[\begin{array}{l}
u \\
v
\end{array}\right]\left[\begin{array}{ll}
-v^{T} & u^{T}
\end{array}\right]
$$

instead of $Z(w)$, where we have put $w=\left[\begin{array}{l}u \\ v\end{array}\right], u, v \in \mathbb{C}^{2 m+1}$. Indeed, it is easy to see that the matrix

$$
A:=\left[\begin{array}{cc}
\mathcal{J}_{2 m+1}(0) & 0 \\
0 & -\mathcal{J}_{2 m+1}(0)^{T}
\end{array}\right]
$$

is $J$-Hamiltonian (see (1.1)), and that the canonical form of the pair $(A, J)$ is the same as the canonical form of $\left(Z(0),\left[\begin{array}{cc}0 & \Sigma_{2 m+1} \\ -\Sigma_{2 m+1} & 0\end{array}\right]\right)$. We shall denote the entries of $u$ and $v$ by $u_{1}, \ldots, u_{2 m+1}$ and $v_{1}, \ldots, v_{2 m+1}$, respectively. Clearly, $M$ is singular for all $u$ and $v$. It is easy to see that for some choice of $u$ and $v$ the rank of $M$ is equal to $4 m+1$, and therefore there exists a generic (with respect to the entries of $u$ and $v$ ) set $\Omega$ such that for every $(u, v) \in \Omega$ the rank of $M$ is equal to $4 m+1$.

Next, we introduce the $(2 m+1) \times(2 m+1)$ matrix $\Upsilon$ :

$$
\Upsilon=\Sigma_{m+1} R_{m+1}=1 \oplus(-1) \oplus 1 \oplus(-1) \oplus \cdots \oplus(-1) \oplus 1
$$

It is useful to note that $\mathcal{J}_{2 m+1}:=\mathcal{J}_{2 m+1}(0)$ and $\Upsilon$ anti-commute:

$$
\Upsilon \mathcal{J}_{2 m+1}=-\mathcal{J}_{2 m+1} \Upsilon
$$

Our first observation is that the vector

$$
x_{1}=\left[\begin{array}{c}
\mathcal{J}_{2 m+1}^{2 m} u \\
\left(-\mathcal{J}_{2 m+1}^{2 m}\right)^{T} v
\end{array}\right] \in \operatorname{Ker} M
$$

Indeed, $A x_{1}=0$, and $\left[\begin{array}{ll}-v^{T} & u^{T}\end{array}\right] x_{1}=0$ as well. Now define for $j=2, \ldots, 2 m+1$ the vectors

$$
x_{j}=(-1)^{j+1}\left[\begin{array}{c}
\Upsilon \mathcal{J}_{2 m+1-j}^{2 m+1-j} u \\
\left(\mathcal{J}_{2 m+1}^{2 m+1-j}\right)^{T} \Upsilon v
\end{array}\right] .
$$

Note that for all $j$ we have $\left[\begin{array}{ll}-v^{T} & u^{T}\end{array}\right] x_{j}=0$, and so

$$
M x_{j}=A x_{j}=(-1)^{j+1}\left[\begin{array}{c}
\mathcal{J}_{2 m+1} \Upsilon \mathcal{J}_{2 m+1}^{2 m+1-j} u \\
-\left(\mathcal{J}_{2 m+1}^{22 m+j}\right)^{T} \Upsilon v
\end{array}\right]=(-1)^{j}\left[\begin{array}{c}
\Upsilon \mathcal{J}_{2 m+1}^{2 m+2-j} u \\
\left(\mathcal{J}_{2 m+1}^{2 m+2-j}\right)^{T} \Upsilon v
\end{array}\right]=x_{j-1}
$$

Thus we see that $x_{1}, \ldots, x_{2 m+1}$ is a Jordan chain of $M$ corresponding to zero.
Next, note that

$$
x_{2 m+1}=\left[\begin{array}{l}
\Upsilon u \\
\Upsilon v
\end{array}\right]
$$

We now define for some complex numbers $a$ and $b$, still to be determined, the vector

$$
x_{2 m+2}=a e_{1}+b e_{4 m+2}+\left[\begin{array}{c}
-(I+\Upsilon) \mathcal{J}_{2 m+1}^{T} u \\
(I+\Upsilon) \mathcal{J}_{2 m+1} v
\end{array}\right]
$$

Then

$$
\begin{aligned}
A x_{2 m+2} & =A\left[\begin{array}{c}
-(I+\Upsilon) \mathcal{J}_{2 m+1}^{T} u \\
(I+\Upsilon) \mathcal{J}_{2 m+1} v
\end{array}\right]=\left[\begin{array}{l}
-\mathcal{J}_{2 m+1}(I+\Upsilon) \mathcal{J}_{2 m+1}^{T} u \\
-\mathcal{J}_{2 m+1}^{T}(I+\Upsilon) \mathcal{J}_{2 m+1} v
\end{array}\right] \\
& =\left[\begin{array}{c}
-(I-\Upsilon) \mathcal{J}_{2 m+1} \mathcal{J}_{2 m+1}^{T} u \\
-(I-\Upsilon) \mathcal{J}_{2 m+1}^{T} \mathcal{J}_{2 m+1} v
\end{array}\right]=\left[\begin{array}{c}
-(I-\Upsilon) u \\
-(I-\Upsilon) v
\end{array}\right]=x_{2 m+1}-\left[\begin{array}{c}
u \\
v
\end{array}\right] .
\end{aligned}
$$

So

$$
\begin{aligned}
M x_{2 m+2} & =\left(x_{2 m+1}-\left[\begin{array}{l}
u \\
v
\end{array}\right]\right)+\left[\begin{array}{l}
u \\
v
\end{array}\right]\left[\begin{array}{ll}
-v^{T} & u^{T}
\end{array}\right] x_{2 m+2} \\
& =x_{2 m+1}+\left[\begin{array}{l}
u \\
v
\end{array}\right]\left(\left[\begin{array}{ll}
-v^{T} & u^{T}
\end{array}\right] x_{2 m+2}-1\right) .
\end{aligned}
$$

If we can choose $a$ and $b$ so that

$$
\left[\begin{array}{ll}
-v^{T} & u^{T}
\end{array}\right] x_{2 m+2}=-a v_{1}+b u_{2 m+1}+2 u^{T}(I+\Upsilon) \mathcal{J}_{2 m+1} v=1
$$

then we have constructed a Jordan chain of length $2 m+2$ for $M$ corresponding to the eigenvalue zero. But it is easily seen that $a$ and $b$ can be chosen as desired, whenever not both $u_{2 m+1}=0$ and $v_{1}=0$. So, generically this can be done.

The next step is to see that generically the Jordan block with eigenvalue zero of $M$ has size $2 m+2$. Here we make essential use of the fact that we already know that the rank of $M$ generically is $4 m+1$, and hence there can be at most one Jordan block with eigenvalue zero in the Jordan normal form of $M$. Then for any Jordan chain it must be possible to extend it to a Jordan chain of length equal to the algebraic multiplicity (this follows, for example, from general results on marked invariant subspaces in [8]). So, it suffices to show that the Jordan chain we have constructed cannot be extended further. For this, observe that vectors $\left[\begin{array}{l}z \\ y\end{array}\right]$ in the range of $M$ are such that $\left[\begin{array}{c}z_{2 m+1} \\ y_{1}\end{array}\right]$ is a multiple of $\left[\begin{array}{c}u_{2 m+1} \\ v_{1}\end{array}\right]$. So, in order for $x_{2 m+2}$ to be in the range of $M$ it is necessary and sufficient that $\left[\begin{array}{c}-u_{2 m} \\ v_{2}\end{array}\right]$ is a multiple of $\left[\begin{array}{c}u_{2 m+1} \\ v_{1}\end{array}\right]$. Obviously, generically this will not be the case.

Next, we show that generically all nonzero eigenvalues are simple eigenvalues. The characteristic polynomial of $M$ is, by what we have shown, generically of the form

$$
x^{2 m+2}\left(x^{2 m}+x^{2 m-2} a_{2 m-2}+\ldots+x^{2} a_{2}+a_{0}\right)
$$

(we also use that $M$ is $J$-Hamiltonian matrix, and so its characteristic polynomial is a polynomial in $x^{2}$ ), and generically, $a_{0} \neq 0$.

Now we find particular vectors $u_{0}$ and $v_{0}$ such that for the characteristic polynomial of $M\left(u_{0}, v_{0}\right)$ we have $a_{2}=a_{4}=\ldots=a_{2 k-2}=0$. Indeed, take $u_{0}$, $v_{0}$ with zero entries,
except for $\left(u_{0}\right)_{2 k},\left(u_{0}\right)_{2 k+1},\left(v_{0}\right)_{1},\left(v_{0}\right)_{2}$. Then

$$
\begin{aligned}
& \operatorname{det}\left(A-x I+\left[\begin{array}{l}
u_{0} \\
v_{0}
\end{array}\right]\left[\begin{array}{ll}
-v_{0}^{T} & u_{0}^{T}
\end{array}\right]\right)= \\
&= \operatorname{det}\left\{(A-x I)\left(I+(A-x I)^{-1}\left[\begin{array}{l}
u_{0} \\
v_{0}
\end{array}\right]\left[\begin{array}{ll}
-v_{0}^{T} & u_{0}^{T}
\end{array}\right]\right)\right\}= \\
&= \operatorname{det}(A-x I) \operatorname{det}\left(I+(A-x I)^{-1}\left[\begin{array}{l}
u_{0} \\
v_{0}
\end{array}\right]\left[\begin{array}{ll}
-v_{0}^{T} & u_{0}^{T}
\end{array}\right]\right)= \\
&= \operatorname{det}(A-x I)\left(1+\left[\begin{array}{ll}
-v_{0}^{T} & u_{0}^{T}
\end{array}\right](A-x I)^{-1}\left[\begin{array}{l}
u_{0} \\
v_{0}
\end{array}\right]\right)= \\
&= x^{4 m+2}\left(1+\left[\begin{array}{ll}
-v_{0}^{T} & \left.\left.u_{0}^{T}\right](A-x I)^{-1}\left[\begin{array}{l}
u_{0} \\
v_{0}
\end{array}\right]\right)= \\
= & x^{4 m+2}\left(1-v^{T}\left(\mathcal{J}_{2 m+1}-x I\right)^{-1} u_{0}-u_{0}^{T}\left(\mathcal{J}_{2 m+1}^{T}+x I\right)^{-1} v_{0}\right) .
\end{array} .=\right.\right. \\
&
\end{aligned}
$$

Now take $u_{0}$ and $v_{0}$ as above, so $v_{0}$ having only the first two entries nonzero and $u_{0}$ having only the last two entries nonzero. Then it is clear that we are interested in the $2 \times 2$ block in the right upper corner of $\left(\mathcal{J}_{2 m+1}-x\right)^{-1}$, and the $2 \times 2$ block in the left lower corner of $\left(\mathcal{J}_{2 m+1}^{T}+x\right)^{-1}$. It is easily computed that

$$
\begin{aligned}
& -v_{0}^{T}\left(\mathcal{J}_{2 m+1}-x I\right)^{-1} u_{0}=\frac{\left(v_{0}\right)_{1}\left(u_{0}\right)_{2 m}+\left(v_{0}\right)_{2}\left(u_{0}\right)_{2 m+1}}{x^{2 m}}+\frac{\left(v_{0}\right)_{2}\left(u_{0}\right)_{2 m}}{x^{2 m-1}}+\frac{\left(v_{0}\right)_{1}\left(u_{0}\right)_{2 m+1}}{x^{2 m+1}} \\
& -u_{0}^{T}\left(\mathcal{J}_{2 m+1}^{T}+x I\right)^{-1} v_{0}=\frac{\left(v_{0}\right)_{1}\left(u_{0}\right)_{2 m}+\left(v_{0}\right)_{2}\left(u_{0}\right)_{2 m+1}}{x^{2 m}}-\frac{\left(v_{0}\right)_{2}\left(u_{0}\right)_{2 m}}{x^{2 m-1}}-\frac{\left(v_{0}\right)_{1}\left(u_{0}\right)_{2 m+1}}{x^{2 m+1}}
\end{aligned}
$$

Because the terms with odd powers cancel, the characteristic polynomial of $M\left(u_{0}, v_{0}\right)$ is given by

$$
\operatorname{det}\left(M\left(u_{0}, v_{0}\right)-x I\right)=x^{2 m+2}\left(x^{2 m}+2\left(\left(v_{0}\right)_{1}\left(u_{0}\right)_{2 m}+\left(v_{0}\right)_{2}\left(u_{0}\right)_{2 m+1}\right)\right)
$$

and so for such a perturbation the nonzero eigenvalues are all simple.
Now, there is an open neighborhood $U$ of the pair $\left(u_{0}, v_{0}\right)$ such that for all matrices $M(u, v)$ with $(u, v) \in U$ all nonzero eigenvalues are simple. Choosing $(u, v) \in U$ so that also the multiplicity of zero of $M(u, v)$ is equal to $2 m+2$, we have found an open set of vectors $w$ with the property that $Z(w)$ has the Jordan form of the required type. But then the set of all vectors $w$ for which $Z(w)$ has the Jordan form of the required type is generic; to see that use the Sylvester resultant matrix of the characteristic polynomial of $Z(w)$ and of its derivative, as it was done in the proof of Lemma 2.5.

The next theorem shows that the situation of Example 4.1 is typical for the case of odd largest partial multiplicity corresponding to the zero eigenvalue. We assume in the next theorem that $A$ has zero as an eigenvalue; if $A$ is invertible, then all statements concerning the zero eigenvalue should be considered as void.

Theorem 4.2 Let $J \in \mathbb{C}^{n \times n}$ be skew-symmetric and invertible, let $A \in \mathbb{C}^{n \times n}$ be $J$ Hamiltonian, with pairwise distinct eigenvalues $\lambda_{1}, \lambda_{2}, \cdots, \lambda_{p}, \lambda_{p+1}=0$ and let $B$ be a rank one perturbation of the form $B=u u^{T} J \in \mathbb{C}^{n \times n}$.

For every $\lambda_{j}, j=1,2, \ldots, p+1$, let $n_{1, j}>n_{2, j}>\ldots>n_{m_{j}, j}$ be the sizes of Jordan blocks in the Jordan form of $A$ associated with the eigenvalue $\lambda_{j}$, and let there be exactly $\ell_{k, j}$ Jordan blocks of size $n_{k, j}$ associated with $\lambda_{j}$ in the Jordan form of A, for $k=1,2, \ldots, m_{j}$.
(1) If $n_{1, p+1}$ is even (in particular, if $A$ is invertible), then generically with respect to the components of $u$, the matrix $A+B$ has the Jordan canonical form

$$
\bigoplus_{j=1}^{p+1}\left(\left(\bigoplus_{s=1}^{\ell_{1, j}-1} \mathcal{J}_{n_{1, j}}\left(\lambda_{j}\right)\right) \oplus\left(\bigoplus_{s=1}^{\ell_{2, j}} \mathcal{J}_{n_{2, j}}\left(\lambda_{j}\right)\right) \oplus \cdots \oplus\left(\bigoplus_{s=1}^{\ell_{m_{j}, j}} \mathcal{J}_{n_{m_{j}, j}}\left(\lambda_{j}\right)\right)\right) \oplus \widetilde{\mathcal{J}}
$$

where $\widetilde{\mathcal{J}}$ contains all the Jordan blocks of $A+B$ associated with eigenvalues different from any of $\lambda_{1}, \ldots, \lambda_{p+1}$.
(2) If $n_{1, p+1}$ is odd (in this case $\ell_{1, p+1}$ is even), then generically with respect to the components of $u$, the matrix $A+B$ has the Jordan canonical form

$$
\begin{align*}
& \bigoplus_{j=1}^{p}\left(\left(\bigoplus_{s=1}^{\ell_{1, j}-1} \mathcal{J}_{n_{1, j}}\left(\lambda_{j}\right)\right) \oplus\left(\bigoplus_{s=1}^{\ell_{2, j}} \mathcal{J}_{n_{2, j}}\left(\lambda_{j}\right)\right) \oplus \cdots \oplus\left(\bigoplus_{s=1}^{\ell_{m_{j}, j}} \mathcal{J}_{n_{m_{j}, j}}\left(\lambda_{j}\right)\right)\right) \\
& \oplus\left(\bigoplus_{s=1}^{\ell_{1, p+1}-2} \mathcal{J}_{n_{1, p+1}}(0)\right) \oplus\left(\bigoplus_{s=1}^{\ell_{2, p+1}} \mathcal{J}_{n_{2, p+1}}(0)\right) \oplus \cdots \oplus\left(\bigoplus_{s=1}^{\ell_{m_{p+1}, p+1}} \mathcal{J}_{n_{m_{p+1}, p+1}}(0)\right) \\
& \oplus \mathcal{J}_{n_{1, p+1}+1}(0) \oplus \widetilde{\mathcal{J}} \tag{4.1}
\end{align*}
$$

where $\tilde{\mathcal{J}}$ contains all the Jordan blocks of $A+B$ associated with eigenvalues different from any of $\lambda_{1}, \ldots, \lambda_{p+1}$.
(3) In either case (1) or (2), generically the part $\widetilde{\mathcal{J}}$ has simple eigenvalues.

Proof. If (1) holds, then it follows from Theorem 2.7 that we can apply Theorem 3.1 or Theorem 3.2, and we immediately obtain the desired result; here we also use the easily verifiable fact that the intersection of finitely many generic sets is again generic.

Consider the case (2). In this case, generically the part of the Jordan form of $A+B$ that involves nonzero eigenvalues has again the form as given in (2), in view of Theorems 2.7, 3.1, and 3.2. It remains to prove that generically the part of the Jordan form of $A+B$ corresponding to the zero eigenvalue has the form

$$
\begin{equation*}
\mathcal{J}_{n_{1}+1}(0) \oplus\left(\bigoplus_{s=1}^{\ell_{1}-2} \mathcal{J}_{n_{1}}(0)\right) \oplus\left(\bigoplus_{s=1}^{\ell_{2}} \mathcal{J}_{n_{2}}(0)\right) \oplus \cdots \oplus\left(\bigoplus_{s=1}^{\ell_{m}} \mathcal{J}_{n_{m}}(0)\right) \oplus \widehat{\mathcal{J}} \tag{4.2}
\end{equation*}
$$

Here, we let $m=m_{p+1} ; n_{k}=n_{k, p+1}$ for $k=1,2, \ldots, m$; and $\ell_{k}=\ell_{k, p+1}$ for $k=$ $1,2, \ldots, m$, and $\widehat{\mathcal{J}}$ contains all the Jordan blocks of $A+B$ associated with nonzero eigenvalues.

To this end, we may assume without loss of generality that $A$ and $J$ are in the form (2.10), where we assume in addition that the diagonal blocks of $A$ and $J$ have been permuted in such a way that the blocks associated with the eigenvalue 0 appear first and that they are ordered with decreasing sizes. Thus, we assume that $A$ and $J$ have the forms

$$
\begin{equation*}
A=\left(\bigoplus_{s=1}^{\ell_{1}} \mathcal{J}_{n_{1}}(0)\right) \oplus \cdots \oplus\left(\bigoplus_{s=1}^{\ell_{m}} \mathcal{J}_{n_{m}}(0)\right) \oplus \widetilde{A} \tag{4.3}
\end{equation*}
$$

where $\sigma(\widetilde{A}) \subseteq \mathbb{C} \backslash\{0\}$ and

$$
J=\bigoplus_{s=1}^{\ell_{1} / 2}\left[\begin{array}{cc}
0 & \Sigma_{n_{1}}  \tag{4.4}\\
-\Sigma_{n_{1}} & 0
\end{array}\right] \oplus J_{2} \oplus \cdots \oplus J_{m} \oplus \widetilde{J}
$$

Then the algebraic and geometric multiplicity $a$ and $g$ of the eigenvalue 0 of $A$ are given by

$$
a=\sum_{s=1}^{m} \ell_{s} n_{s}, \quad g=\sum_{s=1}^{m} \ell_{s}
$$

respectively. The corresponding $J$-Hamiltonian rank one perturbation $B$ has the form $B=u v^{T}=u u^{T} J$, where we partition

$$
u=\left[\begin{array}{c}
u^{(1)} \\
\vdots \\
u^{(m)} \\
\widetilde{u}
\end{array}\right], \quad u^{(i)}=\left[\begin{array}{c}
u^{(i, 1)} \\
\vdots \\
u^{\left(i, \ell_{i}\right)}
\end{array}\right], \quad u^{(i, s)}=\left[\begin{array}{c}
u_{1}^{(i, s)} \\
\vdots \\
u_{n_{i}}^{(i, s)}
\end{array}\right] \in \mathbb{C}^{n_{i}},
$$

for $s=1, \ldots, \ell_{i} ; i=1, \ldots, m$. Thus, $\widetilde{u} \in \mathbb{C}^{n-a}$. We will now show in two steps that generically $A+B$ has the Jordan canonical form (4.2). By Corollary 2.11 we know that $A+B$ has $\ell_{1}-1$ Jordan chains of length $n_{1}$ and $\ell_{j}$ Jordan chains of length $n_{j}, j=2, \ldots, m$ associated with the eigenvalue 0 . In the first step, we will show that generically there exists a Jordan chain of length $n_{1}+1$. In the second step, we will show that the algebraic multiplicity of the eigenvalue zero of $A+B$ generically is $\widetilde{a}=\left(\sum_{s=1}^{m} \ell_{s} n_{s}\right)-n_{1}+1=a-n_{1}+1$. Both steps together obviously imply that (4.2) represents the only possible Jordan canonical forms for $A+B$.

Step 1: Existence of a Jordan chain of length $n_{1}+1$.
Generically, the hypothesis of Corollary 2.11 is satisfied (i.e., specific entries of vectors are nonzero), so generically the matrix $S$ as in Corollary 2.11 exists so that
$S^{-1}(A+B) S$ is in partial Brunovsky form. We first investigate the structure of the vector $v^{T}=u^{T} J$. From (4.4), we obtain that $v$ has the form

$$
v=\left(u^{T} J\right)^{T}=-J u=\left[\begin{array}{c}
v^{(1)}  \tag{4.5}\\
\vdots \\
v^{(m)} \\
\widetilde{v}
\end{array}\right], \quad v^{(i)}=\left[\begin{array}{c}
v^{(i, 1)} \\
\vdots \\
v^{\left(i, \ell_{i}\right)}
\end{array}\right], \quad v^{(i, s)}=\left[\begin{array}{c}
v_{1}^{(i, s)} \\
\vdots \\
v_{n_{i}}^{(i, s)}
\end{array}\right] \in \mathbb{C}^{n_{i}},
$$

for $s=1, \ldots, \ell_{i}$ and $i=1, \ldots, m$, where

$$
v^{(1,2 s-1)}=-\Sigma_{n_{1}} u^{(1,2 s)}=\left[\begin{array}{c}
-u_{n_{1}}^{(1,2 s)} \\
u_{n_{1}-2 s}^{(1,2)} \\
\mp \\
\vdots \\
-u_{1}^{(1,2 s)}
\end{array}\right], \quad v^{(1,2 s)}=\Sigma_{n_{1}} u^{(1,2 s-1)}=\left[\begin{array}{c}
u_{n_{1}}^{(1,2 s-1)} \\
-u_{n_{1}-1}^{(1,2 s-1)} \\
\pm \\
\pm \\
u_{1}^{(1,2 s-1)}
\end{array}\right]
$$

for $s=1, \ldots, \ell_{1} / 2$. Thus, $S^{-1}$ takes the form

$$
S^{-1}=\left(\bigoplus_{s=1}^{\ell_{1}} \operatorname{Toep}\left(v^{(1, s)}\right)\right) \oplus \cdots \oplus\left(\bigoplus_{s=1}^{\ell_{m}} \operatorname{Toep}\left(v^{(m, s)}\right)\right) \oplus I_{n-a}
$$

and it follows that

$$
\begin{equation*}
S^{-1} B S=w(\underbrace{e_{1, n_{1}}^{T}, \ldots, e_{1, n_{1}}^{T}}_{\ell_{1} \text { times }}, \ldots, \underbrace{e_{1, n_{m}}^{T}, \ldots, e_{1, n_{m}}^{T}}_{\ell_{m} \text { times }}, z^{T}) \tag{4.6}
\end{equation*}
$$

for some $z \in \mathbb{C}^{n-a}$. Thus,

$$
w=S^{-1} u=\left[\begin{array}{c}
w^{(1)}  \tag{4.7}\\
\vdots \\
w^{(m)} \\
\widetilde{w}
\end{array}\right], \quad w^{(i)}=\left[\begin{array}{c}
w^{(i, 1)} \\
\vdots \\
w^{\left(i, \ell_{i}\right)}
\end{array}\right], \quad w^{(i, s)}=\left[\begin{array}{c}
w_{1}^{(i, s)} \\
\vdots \\
w_{n_{i}}^{(i, s)}
\end{array}\right] \in \mathbb{C}^{n_{i}}
$$

for $s=1, \ldots, \ell_{i}$ and $i=1, \ldots, m$, where

$$
\begin{equation*}
w_{n_{1}}^{(1,2 s-1)}=-u_{n_{1}}^{(1,2 s)} u_{n_{1}}^{(1,2 s-1)}, \quad w_{n_{1}}^{(1,2 s)}=u_{n_{1}}^{(1,2 s-1)} u_{n_{1}}^{(1,2 s)}=-w_{n_{1}}^{(1,2 s-1)} \tag{4.8}
\end{equation*}
$$

and, provided that $n_{1}>1$,

$$
\begin{align*}
w_{n_{1}-1}^{(1,2 s-1)} & =u_{n_{1}-1}^{(1,2 s)} u_{n_{1}}^{(1,2 s-1)}-u_{n_{1}}^{(1,2 s)} u_{n_{1}-1}^{(1,2 s-1)},  \tag{4.9}\\
w_{n_{1}-1}^{(1,2 s)} & =-u_{n_{1}}^{(1,2 s)} u_{n_{1}-1}^{(1,2 s-1)}+u_{n_{1}-1}^{(1,2 s)} u_{n_{1}}^{(1,2 s-1)}=w_{n_{1}-1}^{(1,2 s-1)}, \tag{4.10}
\end{align*}
$$

for $s=1, \ldots, \ell_{1} / 2$. Consider the following Jordan chains associated with the eigenvalue 0 of $S^{-1}(A+B) S$ denoted by $C_{i, s}$ :
length $n_{1}: \quad C_{1, s}: \quad e_{2(s-1) n_{1}+1}-e_{(2 s-1) n_{1}+1}, \ldots, e_{(2 s-1) n_{1}}-e_{2 s n_{1}}, \quad s=1, \ldots, \frac{\ell_{1}}{2}$ length $n_{i}: \quad C_{i, s}: \quad-e_{1}+e_{\Sigma_{k=1}^{i-1} \ell_{k} n_{k}+(s-1) n_{i}+1}, \ldots,-e_{n_{i}}+e_{\Sigma_{k=1}^{i-1} \ell_{k} n_{k}+s n_{i}}, \quad s=1, \ldots, \ell_{i}$,
where $i=2, \ldots, m$. Observe that $C_{i, s}, i \neq 1$, are just the Jordan chains from Corollary 2.11 multiplied by -1 while the chains $C_{1, s}$ are linear combinations of the Jordan chains from Corollary 2.11. Namely, in the notation of (2.13), and numbering the chains in (2.13) first, second, etc., from the top to the bottom, we see that the chains $C_{1,1}, \ldots, C_{1, \ell_{1} / 2}$ are the first chain, the negative of the second chain plus the third chain, $\ldots$, the negative of the $\left(\ell_{1}-2\right)$ th chain plus the $\left(\ell_{1}-1\right)$ th chain, respectively. Now consider the Jordan chain

$$
C:=\left(\sum_{s=1}^{\ell_{1} / 2} \alpha_{1, s} C_{1, s}\right)+\sum_{i=2}^{m} \sum_{s=1}^{\ell_{i}} \alpha_{i, s} C_{i, s}
$$

of length $n_{1}$ (see Definition 2.13), and let $y$ denote the $n_{1}$-th (and thus last) vector of this chain. We next show that the Jordan chain $C$ can be extended by a vector $y$ to a Jordan chain of length $n_{1}+1$ associated with the eigenvalue 0 , for some particular choice of the parameters $\alpha_{i, s}$ (depending on $u$ ) such that generically at least one of $\alpha_{1,1}, \ldots \alpha_{1, \ell_{1} / 2}$ is nonzero. To see this, we have to show that $y$ is in the range of $S^{-1}(A+B) S$. First, partition

$$
y=\left[\begin{array}{c}
y^{(1)} \\
\vdots \\
y^{(m)} \\
\widetilde{y}
\end{array}\right], \quad y^{(i)}=\left[\begin{array}{c}
y^{(i, 1)} \\
\vdots \\
y^{\left(i, \ell_{i}\right)}
\end{array}\right], \quad y^{(i, s)}=\left[\begin{array}{c}
y_{1}^{(i, s)} \\
\vdots \\
y_{n_{i}}^{(i, s)}
\end{array}\right] \in \mathbb{C}^{n_{i}}
$$

for $s=1, \ldots, \ell_{i} ; i=1, \ldots, m$. Then by the definition of $y$, we have $\widetilde{y}=0 \in \mathbb{C}^{n-a}$,

$$
\begin{aligned}
y_{n_{1}}^{(1,2 s-1)} & =\alpha_{1, s}, \quad y_{n_{1}}^{(1,2 s)}=-\alpha_{1, s}, \quad s=1, \ldots, \ell_{1} / 2 \\
y_{n_{i}}^{(i, s)} & =\alpha_{i, s}, \quad s=1, \ldots, \ell_{i} ; i=2, \ldots, m
\end{aligned}
$$

We have to solve the linear system

$$
\begin{equation*}
S^{-1}(A+B) S x=y \tag{4.11}
\end{equation*}
$$

Partitioning

$$
x=\left[\begin{array}{c}
x^{(1)} \\
\vdots \\
x^{(m)} \\
\widetilde{x}
\end{array}\right], \quad x^{(i)}=\left[\begin{array}{c}
x^{(i, 1)} \\
\vdots \\
x^{\left(i, \ell_{i}\right)}
\end{array}\right], \quad x^{(i, s)}=\left[\begin{array}{c}
x_{1}^{(i, s)} \\
\vdots \\
x_{n_{i}}^{(i, s)}
\end{array}\right] \in \mathbb{C}^{n_{i}},
$$

and making the ansatz $\widetilde{x}=0$, then equation (4.11) becomes (we here use (4.6, (4.7))):

$$
\begin{align*}
& w_{k}^{(i, s)}\left(\sum_{\nu=1}^{m} \sum_{\mu=1}^{\ell_{\nu}} x_{1}^{(\nu, \mu)}\right)+x_{k+1}^{(i, s)}=y_{k}^{(i, s)}, \quad k=1, \ldots, n_{i}-1 ; s=1, \ldots, \ell_{i} ; i=1, \ldots, m,  \tag{4.12}\\
& w_{n_{i}}^{(i, s)}\left(\sum_{\nu=1}^{m} \sum_{\mu=1}^{\ell_{\nu}} x_{1}^{(\nu, \mu)}\right)=\alpha_{i, s},  \tag{4.13}\\
& s=1, \ldots, \ell_{i} ; i=2, \ldots, m,  \tag{4.14}\\
& w_{n_{1}}^{(1,2 s-1)}\left(\sum_{\nu=1}^{m} \sum_{\mu=1}^{\ell_{\nu}} x_{1}^{(\nu, \mu)}\right)=\alpha_{1, s}, \quad s=1, \ldots, \ell_{1} / 2,  \tag{4.15}\\
& w_{n_{1}}^{(1,2 s)}\left(\sum_{\nu=1}^{m} \sum_{\mu=1}^{\ell_{\nu}} x_{1}^{(\nu, \mu)}\right)=-\alpha_{1, s}, \quad s=1, \ldots, \ell_{1} / 2 .
\end{align*}
$$

Set $x_{1}^{(1,1)}=1$ and $x_{1}^{(\nu, \mu)}=0$, for $\mu=1, \ldots, \ell_{\nu} ; \nu=1, \ldots, m ;(\nu, \mu) \neq(1,1)$, as well as $\alpha_{i, s}=w_{n_{i}}^{(i, s)}$ for $s=1, \ldots, \ell_{i} ; i=2, \ldots, m$ and $\alpha_{1, s}=w_{n_{1}}^{(1,2 s-1)}$ for $s=1, \ldots, \ell_{1} / 2$. Then (4.13) and (4.14) are satisfied and so is (4.15), because by (4.8) we have

$$
w_{n_{1}}^{(1,2 s)}=u_{n_{1}}^{(1,2 s)} u_{n_{1}}^{(1,2 s-1)}=-w_{n_{1}}^{(1,2 s-1)}=-\alpha_{1, s}, \quad s=1, \ldots, \ell_{1} / 2 .
$$

Finally, (4.12) can be solved by choosing $x_{k+1}^{(i, s)}=y_{k}^{(i, s)}-w_{k}^{(i, s)}$ for $k=1, \ldots, n_{i}-1$; $s=1, \ldots, \ell_{i} ; i=1, \ldots, m$.

Step 2: We show that the algebraic multiplicity of the eigenvalue 0 of $A+B$ generically is $\widetilde{a}=\left(\sum_{s=1}^{m} \ell_{s} n_{s}\right)-n_{1}+1=a-n_{1}+1$.

Let $\mu_{1}, \ldots, \mu_{q}$ denote the pairwise distinct nonzero eigenvalues of $A$ and let $r_{1}, \ldots, r_{q}$ be their algebraic multiplicities. By Corollary 2.11, the lowest possible power of $\lambda$ associated with a nonzero coefficient in $p_{0}(\lambda)$ is $a-n_{1}$ and the corresponding coefficient $c_{a-n_{1}}$ is

$$
c_{a-n_{1}}=(-1)^{a}\left(\prod_{i=1}^{q} \mu_{i}^{r_{i}}\right)\left(\sum_{s=1}^{\ell_{1}} w_{n_{1}}^{(1, s)}\right)=0
$$

because of (4.8). If $n_{1}=1$ then $\widetilde{a}=a$ and there is nothing to show as the algebraic multiplicity of the eigenvalue zero cannot increase when a generic perturbation is applied. Otherwise, we distinguish the cases $n_{2}<n_{1}-1$ and $n_{2}=n_{1}-1$. If $n_{2}<n_{1}-1$, then by Corollary 2.11 the coefficient $c_{a-n_{1}+1}$ of $\lambda^{a-n_{1}+1}$ in $p_{0}(\lambda)$ is

$$
\begin{aligned}
c_{a-n_{1}+1} & =-\left(\sum_{\nu=1}^{q} r_{\nu} \mu_{\nu}^{r_{\nu}-1} \prod_{\substack{i=1 \\
i \neq \nu}}^{q} \mu_{i}^{r_{i}}\right) \underbrace{\left(\sum_{s=1}^{\ell_{1}} w_{n_{1}}^{(1, s)}\right)}_{=0 \text { using (4.8) }}+\left(\prod_{i=1}^{q} \mu_{i}^{r_{i}}\right)\left(\sum_{s=1}^{\ell_{1}} w_{n_{1}-1}^{(1, s)}\right) \\
\text { using } & =(4.10) \quad\left(\prod_{i=1}^{q} \mu_{i}^{r_{i}}\right)\left(\sum_{s=1}^{\ell_{1} / 2} 2\left(u_{n_{1}-1}^{(1,2 s)} u_{n_{1}}^{(1,2 s-1)}-u_{n_{1}}^{(1,2 s)} u_{n_{1}-1}^{(1,2 s-1)}\right)\right)
\end{aligned}
$$

which generically is nonzero. If, on the other hand, $n_{2}=n_{1}-1$, then again by Corollary 2.11 the coefficient $c_{a-n_{1}+1}$ of $\lambda^{a-n_{1}+1}$ in $p_{0}(\lambda)$ is

$$
c_{a-n_{1}+1}=\left(\prod_{i=1}^{q} \mu_{i}^{r_{i}}\right)\left(\sum_{s=1}^{\ell_{1}} w_{n_{1}-1}^{(1, s)}+\sum_{s=1}^{\ell_{2}} w_{n_{2}}^{(2, s)}\right) .
$$

Since $n_{1}>1$ is odd, $n_{2} \geq 2$ is even and the block $J_{2}$ in (4.4) takes the form

$$
J_{2}=\Sigma_{n_{2}} \oplus \cdots \oplus \Sigma_{n_{2}}
$$

Hence, for the component $v_{2}$ in (4.5) we obtain that

$$
v^{(2, s)}=-\Sigma_{n_{2}} u^{(2, s)}=\left[\begin{array}{c}
-u_{n 2}^{(2, s)} \\
u_{n_{2}-1}^{(2, s)} \\
\mp \\
\vdots \\
u_{1}^{(2, s)}
\end{array}\right], \quad s=1, \ldots, \ell_{2}
$$

and thus

$$
w_{n_{2}}^{(2, s)}=-u_{n_{2}}^{(2, s)} v_{1}^{(2, s)}=-\left(u_{n_{2}}^{(2, s)}\right)^{2}
$$

which gives

$$
c_{a-n_{1}+1}=\left(\prod_{i=1}^{q} \mu_{i}^{r_{i}}\right)\left(\sum_{s=1}^{\ell_{1} / 2} 2\left(u_{n_{1}-1}^{(1,2 s)} u_{n_{1}}^{(1,2 s-1)}-u_{n_{1}}^{(1,2 s)} u_{n_{1}-1}^{(1,2 s-1)}\right)-\sum_{s=1}^{\ell_{2}}\left(u_{n_{2}}^{(2, s)}\right)^{2}\right)
$$

Again, this is nonzero generically. In all cases, we have shown that zero is a root of $p_{0}(\lambda)$ with multiplicity $a-n_{1}+1$. Thus, the algebraic multiplicity of the eigenvalue zero of $A+B$ is $a-n_{1}+1$. Together with Step 1 , we obtain that (4.2) generically are the only possible Jordan canonical forms of $A+B$.

Finally, we prove part (3) by following the arguments of the proof of part (b) of Theorem 2.3, and using Examples 4.3-4.5 and Lemma 4.6 (instead of Lemma 2.5 that was used the proof of Theorem 2.3) presented in the remainder of the section.

Example 4.3 Let

$$
Z^{(2)}(\alpha)=\mathcal{J}_{2 m}(0)+\left(\alpha e_{2 m}\right)\left(\alpha e_{2 m}^{T}\right) \Sigma_{2 m} \in \mathbb{C}^{2 m \times 2 m}, \quad \alpha \in \mathbb{C} \backslash\{0\}
$$

Analogously to Example 2.4, we have $\chi\left(Z^{(2)}(\alpha)\right)=x^{2 m}+\alpha^{2}$, in particular, $Z^{(2)}(\alpha)$ has $2 m$ distinct nonzero eigenvalues.

Example 4.4 Comsider the $(4 m+2) \times(4 m+2)$ matrix

$$
Z^{(3)}(\alpha, w)=\left[\begin{array}{cc}
\mathcal{J}_{2 m+1}(0) & 0 \\
0 & \mathcal{J}_{2 m+1}(0)
\end{array}\right]+(\alpha w)\left(\alpha w^{T}\right)\left[\begin{array}{cc}
0 & \Sigma_{2 m+1} \\
-\Sigma_{2 m+1} & 0
\end{array}\right]
$$

where $\alpha \in \mathbb{C} \backslash\{0\}$. It follows from Example 4.1 that there exist a nonzero vector $w$ and $\epsilon>0$ with the property that the matrix $Z^{(3)}(\alpha, w)$ has the Jordan form $\mathcal{J}_{2 m+2}(0) \oplus \mathcal{K}$, where $\mathcal{K}$ is a diagonal invertible matrix with distinct diagonal entries, for every $\alpha$ in the punctured disc $0<|\alpha|<\epsilon$.

Example 4.5 Let

$$
\begin{aligned}
Z^{(4)}(\lambda, \alpha)= & {\left[\begin{array}{cc}
\mathcal{J}_{m}(\lambda) & 0 \\
0 & -\mathcal{J}_{m}(\lambda)^{T}
\end{array}\right]+(\alpha u)(\alpha u)^{T}\left[\begin{array}{cc}
0 & I_{m} \\
-I_{m} & 0
\end{array}\right] \in \mathbb{C}^{2 m \times 2 m} } \\
& \lambda \in \mathbb{C} \backslash\{0\}, \quad \alpha \in \mathbb{C} \backslash\{0\}
\end{aligned}
$$

Let

$$
u=\left[\begin{array}{c}
e_{m} \\
e_{1}
\end{array}\right]
$$

We shall prove that there exists $\epsilon>0$ which depends only on $\lambda$ and on $m$, such that for all $\alpha$ with $0<|\alpha|<\epsilon$, the matrix $Z^{(4)}(\lambda, \alpha)$ has $2 m$ distinct eigenvalues and none of them is equal to $\pm \lambda$.

Using the Laplace theorem for determinants with respect to the first $m$ rows of $\operatorname{det}\left(x I-Z^{(4)}(\lambda, \alpha)\right)$, and omitting terms that are obviously zeros, we easily compute

$$
\begin{gathered}
\chi\left(Z^{(4)}(\lambda, \alpha)\right)=\left((x-\lambda)^{m}+\alpha^{2}\right)\left((x+\lambda)^{m}+(-1)^{m} \alpha^{2}\right)+(-1)^{m+1} \alpha^{4}= \\
(x-\lambda)^{m}(x+\lambda)^{m}+\alpha^{2}(x+\lambda)^{m}+(-1)^{m} \alpha^{2}(x-\lambda)^{m}
\end{gathered}
$$

Clearly, $\pm \lambda$ are not zeros of $\chi\left(Z^{(4)}(\lambda, \alpha)\right.$ because $\lambda \neq 0, \alpha \neq 0$. Assuming that $\chi\left(Z^{(4)}(\lambda, \alpha)\right)$ and $\frac{\partial}{\partial x} \chi\left(Z^{(4)}(\lambda, \alpha)\right)$ have a common root $x_{0}$, we have the equalities

$$
\begin{gather*}
\left(x_{0}-\lambda\right)^{m}\left(x_{0}+\lambda\right)^{m}+\beta\left(x_{0}+\lambda\right)^{m}+(-1)^{m} \beta\left(x_{0}-\lambda\right)^{m}=0  \tag{4.16}\\
\left(x_{0}-\lambda\right)^{m-1}\left(x_{0}+\lambda\right)^{m}+\left(x_{0}-\lambda\right)^{m}\left(x_{0}+\lambda\right)^{m-1}+\beta\left(x_{0}+\lambda\right)^{m-1}+(-1)^{m} \beta\left(x_{0}-\lambda\right)^{m-1} \tag{4.17}
\end{gather*}
$$

where $\beta=\alpha^{2}$. Multiplying (4.17) by $x_{0}-\lambda$ and using (4.16) yields after simple algebra

$$
\left(x_{0}-\lambda\right)^{m+1}=2 \beta \lambda
$$

Analogously $\left(x_{0}+\lambda\right)^{m+1}=(-1)^{m+1} 2 \beta \lambda$ is obtained. These equalities are contradictory if $|\alpha|$ is sufficiently small.

Using Examples 4.3, 4.4, and 4.5, and the already proved parts (1) and (2) of Theorem 4.2, the following lemma is proved in the same way as Lemma 2.5. We omit the details of proof.

Lemma 4.6 Let $\Omega^{\prime}$ be the (open) generic set of vectors $u \in \mathbb{C}^{n}$ for which (1) or (2) of Theorem 4.2 holds. Then there is $\epsilon>0$ and an open dense (in the ball $\left\{u \in \mathbb{C}^{n}:\|u\|<\right.$ $\epsilon\}$ ) set $\Omega^{\prime \prime} \subseteq \Omega^{\prime}$ such that for every $u \in \Omega^{\prime \prime},\|u\|<\epsilon$, the Jordan form of $A+u u^{T} J$ is of the type described in items (1) - (3) of Theorem 4.2.

We conclude that in case (1) of Theorem 4.2 generically all Jordan blocks associated with eigenvalues $\lambda_{1}, \ldots, \lambda_{p+1}$ remain unchanged except for one block of the largest size for every eigenvalue $\lambda_{j}$ which disappears (leading to eigenvalues different from $\lambda_{j}$ ). In the case (2), the generic behavior of Jordan blocks of nonzero eigenvalues is the same as in the case (1), whereas all Jordan blocks associated with the zero eigenvalue remain unchanged except for two of the largest size ones of which one of them disappears (leading to nonzero eigenvalues), while the other one increases its size by one.

## 5 Generic structured rank one perturbations for complex $H$-symmetric matrices

Our next result concerns perturbations of $H$-symmetric matrices. According to Theorem 2.6, every (complex) matrix is similar to an $H$-symmetric matrix, for any fixed symmetric invertible matrix $H$.
Theorem 5.1 Let $H \in \mathbb{C}^{n \times n}$ be symmetric and invertible, $A \in \mathbb{C}^{n \times n}$ be $H$-symmetric, with distinct eigenvalues $\lambda_{1}, \lambda_{2}, \cdots, \lambda_{p}$ and having the Jordan canonical form

$$
\bigoplus_{j=1}^{p}\left(\left(\bigoplus_{s=1}^{\ell_{1, j}} \mathcal{J}_{n_{1, j}}\left(\lambda_{j}\right)\right) \oplus\left(\bigoplus_{s=1}^{\ell_{2, j}} \mathcal{J}_{n_{2, j}}\left(\lambda_{j}\right)\right) \oplus \cdots \oplus\left(\bigoplus_{s=1}^{\ell_{m_{j}, j}} \mathcal{J}_{n_{m_{j}, j}}\left(\lambda_{j}\right)\right)\right)
$$

where $n_{1, j}>n_{2, j}>\ldots>n_{m_{j}, j}, j=1, \ldots p$. Let $B \in \mathbb{C}^{n \times n}$ be a rank one perturbation of the form $B=u u^{T} H, u \in \mathbb{C}^{n}$. Then:
(1) generically (with respect to the components of $u$ ), the matrix $A+B$ has the Jordan canonical form

$$
\bigoplus_{j=1}^{p}\left(\left(\bigoplus_{s=1}^{\ell_{1, j}-1} \mathcal{J}_{n_{1, j}}\left(\lambda_{j}\right)\right) \oplus\left(\bigoplus_{s=1}^{\ell_{2, j}} \mathcal{J}_{n_{2, j}}\left(\lambda_{j}\right)\right) \oplus \cdots \oplus\left(\bigoplus_{s=1}^{\ell_{m_{j}, j}} \mathcal{J}_{n_{m_{j}, j}}\left(\lambda_{j}\right)\right)\right) \oplus \widetilde{\mathcal{J}}
$$

where $\widetilde{\mathcal{J}}$ contains all the Jordan blocks of $A+B$ associated with eigenvalues different from any of $\lambda_{1}, \ldots, \lambda_{p}$;
(2) generically, all eigenvalues of $A+B$ different from any of $\lambda_{1}, \ldots, \lambda_{p}$, are simple.

Proof. Part (1) follows immediately from Theorem 2.6 and Theorem 3.1.
Part (2) is proved completely analogously to the proofs of part (b) of Theorem 2.3 and part (3) of Theorem 4.2 by using the following lemma which is based on Example 2.4. We omit details.

Lemma 5.2 Let $\Omega^{\prime}$ be the open generic set of vectors $u \in \mathbb{C}^{n}$ for which (1) of Theorem 5.1 holds. Then there is $\epsilon>0$ and an open dense (in the ball $\left\{u \in \mathbb{C}^{n}:\|u\|<\epsilon\right\}$ ) set $\Omega^{\prime \prime} \subseteq \Omega^{\prime}$ such that for every $u \in \Omega^{\prime \prime},\|u\|<\epsilon$, the Jordan form of $A+u u^{T} H$ is of the type described in items (1) and (2) of Theorem 5.1.

## 6 Conclusion

We have presented several results on Jordan structures of matrices under structured and unstructured rank one perturbations in a general context, and studied the perturbation analysis for the Jordan structures of complex $J$-Hamiltonian and complex $H$-symmetric under structured rank one perturbations. We have shown that as in the case of unstructured perturbations, generically only (one of) the largest Jordan blocks is destroyed. However in the structured case, there is a particular situation, where the effect of generic structured perturbation differs from the effect of generic unstructured perturbations. If the largest Jordan blocks associated with the eigenvalue zero of a complex Hamiltonian matrix have odd size, then this Jordan block occurs multiple times. As a result of the perturbation, one of the largest Jordan blocks is destroyed, but one obtains an increase of the size of one of the other blocks of largest size.

In subsequent papers, this perturbation analysis will be extended to the case of $H$-selfadjoint and real $H$-symmetric matrices under real perturbations.

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