# A REGULARIZATION METHOD FOR THE NUMERICAL SOLUTION OF ELLIPTIC BOUNDARY CONTROL PROBLEMS WITH POINTWISE STATE CONSTRAINTS 

F. TRÖLTZSCH* AND I. YOUSEPT*


#### Abstract

A Lavrentiev type regularization technique for solving elliptic boundary control problems with pointwise state constraints is considered. The main concept behind this regularization is to look for controls in the range of the adjoint control-to-state mapping. After investigating the analysis of the method, a semismooth Newton method based on the optimality conditions is presented. The theoretical results are confirmed by numerical tests. Moreover, they are validated by comparing the regularization technique with standard numerical codes based on the discretize-then-optimize concept.


Key words. Boundary control, state constraints, Lavrentiev type regularization, semismooth Newton method, optimize-then-discretize, nested iteration.

AMS subject classifications. 49J20, 49M05, 65 K 10 .

1. Introduction. In this paper, we suggest a regularization method for solving a class of state-constrained elliptic boundary control problems. We consider the following fairly standard problem:

$$
\begin{equation*}
\text { minimize } J(u, y):=\frac{1}{2} \int_{\Omega}\left(y(x)-y_{d}(x)\right)^{2} d x+\frac{\alpha}{2} \int_{\Gamma} u^{2}(x) d s \tag{P}
\end{equation*}
$$

subject to the elliptic boundary value problem

$$
\begin{align*}
A y & =0 & \text { in } \quad \Omega \\
\partial_{n} y & =u & \text { on } \Gamma \tag{1.1}
\end{align*}
$$

and to the pointwise state constraints

$$
\begin{equation*}
y_{a}(x) \leq y(x) \leq y_{b}(x) \quad \text { for almost all } x \text { in } \Omega \tag{1.2}
\end{equation*}
$$

Here, $A$ is a linear uniformly elliptic differential operator and $\Omega \subset \mathbb{R}^{N}, N \geq 2$, is a bounded domain with sufficiently regular boundary $\Gamma$. The functions $y_{d}, y_{a}, y_{b}$ and the cost parameter $\alpha>0$ are assumed to be given data. We do not require that the state constraints are satisfied for all $x \in \bar{\Omega}$, since this will be meaningful only for two-dimensional domains. If the admissible set of $(P)$ is not empty, then, for any dimension $N$, the problem $(P)$ admits a unique optimal solution.

If $N=2$, necessary optimality conditions for this problem follow from Casas [4] or from the Pontryagin principle presented in Casas [5], or Alibert and Raymond [1] for more general problems with semilinear equations (cf. also our remarks on the dimension $N$ below).

When solving this problem numerically, one of the main difficulties is the lack of regularity of the Lagrange multipliers associated with the state constraints. Moreover, there is another, somewhat hidden obstacle in the analysis: To show the existence of these multipliers - i.e. to gain Karush-Kuhn-Tucker type optimality conditions - the constraints (1.2) should be considered in $\mathcal{C}(\bar{\Omega})$. This,

[^0]however, restricts the theory to two-dimensional domains $\Omega$. In this case, the mapping $u \mapsto y$ defined by (1.1) is continuous from $L^{2}(\Gamma)$ to $\mathcal{C}(\bar{\Omega})$. Another way would consist of restricting the constraints (1.2) to a compact subset $\Omega^{\prime} \subset \Omega$.

This obstacle is overcome after regularizing problem (P). Then necessary optimality conditions can be stated for the regularized problem for any dimension $N \geq 2$. To show, however, convergence for vanishing regularization parameter, the restriction to $N=2$ is more or less needed, again.

It is well known that the numerical treatment of state-constrained problems is a quite difficult issue. On the one hand, the measure type form of Lagrange multipliers complicates the numerical treatment of the problems. On the other hand, in the analysis one is faced with some ill-conditioned equations when dealing with state-constrained problems. This is mostly due to the compactness of the mapping $u \mapsto y$. This is known for distributed optimal control problems and it turns out to be even harder in the case of boundary control.

In the last years, two different regularization concepts were proposed to overcome the difficulties mentioned previously. First, Ito and Kunisch [10] suggested a Moreau-Yosida type regularization approach, which removes the pointwise state inequality constraints by adding a penalty term to the objective functional. Hereafter, the penalized problems are solved in an efficient way. We also refer to [2], [3], and [11].

Later, Meyer et al. [14] came up with a Lavrentiev type regularization to the pointwise state inequality constraints, see also the case of pure state constraints in [15]. In contrast to the first method, it preserves, in some sense, the structure of the state-constrained problem. Let us briefly introduce this regularization technique and compare it with the main issue of our paper. Suppose that the state equation is given by

$$
\begin{array}{lll}
A y & =u & \text { in } \\
\partial_{n} y=0 & \text { on } & \Gamma
\end{array}
$$

with distributed control $u$. Then the Lavrentiev type regularization converts the state constraints (1.2) into the mixed control-state constraints

$$
\begin{equation*}
y_{a} \leq y+\lambda u \leq y_{b} \quad \text { a.e. in } \Omega \tag{1.3}
\end{equation*}
$$

with a small parameter $\lambda>0$. This technique is not applicable in the case of boundary control, since the control is defined only on $\Gamma$, while (1.3) must be considered on $\Omega$. The domains of $u$ and $y$ do not fit together.

Our main idea to overcome this difficulty is as follows: Let $S: L^{2}(\Gamma) \rightarrow L^{2}(\Omega)$ denote the control-to-state mapping defined by (1.1). We look for controls in the range of the adjoint operator $S^{\star}$, i.e., we use the ansatz

$$
u=S^{\star} v
$$

with a new control $v$ defined in $\Omega$. Then, the pointwise state constraints (1.2) admit the form $y_{a} \leq S S^{\star} v \leq y_{b}$ and, substituting a new "control function" $w \in L^{2}(\Omega)$ by $w=S S^{\star} v$, we obtain the "control constraints"

$$
y_{a} \leq w \leq y_{b}
$$

However, this is too formal, since $S S^{\star}$ is compact in $L^{2}(\Omega)$, and hence the equation $S S^{\star} v=w$ is ill-posed, if $w \in L^{2}(\Omega)$ is given. To obtain a well-posed equation, we apply the Lavrentiev type regularization, cf. [12], i.e. we write

$$
\lambda v+S S^{\star} v=w
$$

Finally, we arrive at the constraints

$$
y_{a} \leq \lambda v+S S^{\star} v \leq y_{b} \quad \text { a.e. in } \Omega
$$

There are two reasons for the ansatz $u=S^{\star} v$. We obtain from the necessary optimality condition (2.3) that the optimal control $\bar{u}$ is in the range of the adjoint control-to-state mapping $G^{\star}: \mathcal{C}(\bar{\Omega})^{\star} \rightarrow L^{2}(\Gamma)$ (notice that $p=G^{\star}\left(\bar{y}-y_{d}+\mu_{b}-\mu_{a}\right)$ ). By restricting the domain of $G^{\star}$ to $L^{2}(\Omega)$, we avoid measures and arrive at $S^{\star}$. Moreover, we obtain the representation $y=S S^{\star} v$ with a positive semidefinite and self-adjoint operator $S S^{\star}$, which is useful for computations.

We investigate the analysis of this idea as well as its numerical performance. The method slightly increases the number of unknowns and doubles the number of equations. This is certainly some drawback. However, the numerical results are encouraging. In some cases, the results were even better then those obtained by the discretize-then-optimize concept that was used to compare our method. Our main aim was to find an extension of the Lavrentiev type regularization concept to the case of boundary controls that guarantees the existence of regular Lagrange multipliers. Moreover, it should generate a problem that is equivalent to a control-constrained one. Then, we have access to known results on numerical approximations of control-constrained problems such as error estimates or mesh-independence principles.

To give the reader a better orientation on the possible choices for the dimension $N$, we mention already here that $N=2$ is only needed to show the convergence of optimal solutions for vanishing regularization parameter. All other results on the regularized problems hold true for arbitrary dimension. If $N=2$ is needed, we explicitely state this in the associated theorems.

The paper is organized as follows: After introducing the general assumptions as well as our notation in Section 2, we analyze different aspects of our regularization in Section 3. We discuss the optimality conditions for the regularized version of $(P)$ and based on them, we present a semismooth Newton algorithm in Section 4. Finally, in Section 5, we provide a numerical report including a validation of our theoretical results and a comparison of our technique with the application of the discretize-then-optimize concept.
2. General assumptions and notation. Throughout this paper, we consider a bounded domain $\Omega \subset \mathbb{R}^{N}, N \geq 2$, with a $\mathcal{C}^{0,1}$-boundary $\Gamma$. The lower and upper bounds $y_{a}, y_{b} \in \mathcal{C}(\bar{\Omega})$ satisfy $y_{a}(x)<y_{b}(x)$ for all $x \in \bar{\Omega}$. Furthermore, the desired state $y_{d}$ is given in $L^{2}(\Omega)$. If $V$ is a linear normed function space, then we use the notation $\|\cdot\|_{V}$ for the standard norm used in $V$. By $A$, we denote the second-order elliptic partial differential operator defined by

$$
A y(x)=-\sum_{i, j=1}^{N} D_{i}\left(a_{i j}(x) D_{j} y(x)\right)+c_{0}(x) y(x)
$$

where the coefficient functions $a_{i j} \in C^{0,1}(\bar{\Omega})$ satisfy the ellipticity condition

$$
\sum_{i, j=1}^{n} a_{i j}(x) \xi_{i} \xi_{j} \geq \theta\|\xi\|_{L^{2}(\Omega)}^{2} \quad \forall(\xi, x) \in \mathbb{R}^{n} \times \bar{\Omega}
$$

for some constant $\theta>0$. By $A^{\star}$, the associated formally adjoint operator is denoted. We assume that $c_{0} \in L^{\infty}(\Omega)$ is non-negative with $\left\|c_{0}\right\|_{L^{\infty}(\Omega)} \neq 0$. By $G$, we denote the solution operator $G: L^{2}(\Gamma) \rightarrow H^{1}(\Omega)$ that assigns to each $u \in L^{2}(\Gamma)$ the weak solution $y=y(u) \in H^{1}(\Omega)$ of the
elliptic equation

$$
\begin{aligned}
A y & =0 \text { in } \Omega \\
\partial_{n} y & =u \text { on } \Gamma
\end{aligned}
$$

in which $\partial_{n} y$ denotes the co-normal derivative of $y$ (often denoted by $\partial n_{A}$ ). For later use, we set $S=i_{0} G$, where $i_{0}$ is the compact embedding operator from $H^{1}(\Omega)$ to $L^{2}(\Omega)$. Now, our problem can be expressed as follows:

$$
\left\{\begin{array}{l}
\operatorname{minimize} f(u):=\frac{1}{2}\left\|S u-y_{d}\right\|_{L^{2}(\Omega)}^{2}+\frac{\alpha}{2}\|u\|_{L^{2}(\Gamma)}^{2}  \tag{P}\\
\text { over } u \in L^{2}(\Gamma) \\
\text { subject to } y_{a} \leq G u \leq y_{b} \quad \text { a.e. in } \Omega
\end{array}\right.
$$

We still use $G$ instead of $S$ in the constraints, since we need higher regularity of $y$ in Theorem 2.1.
REmark 2.1. Consider the following more general problem

$$
\left\{\begin{array}{l}
\text { minimize } \frac{1}{2}\left\|y-y_{d}\right\|_{L^{2}(\Omega)}^{2}+\frac{\alpha}{2}\left\|u-u_{d}\right\|_{L^{2}(\Gamma)}^{2}  \tag{2.1}\\
\text { over }(u, y) \in L^{2}(\Omega) \times H^{1}(\Omega) \\
\text { subject to } A y=e \quad \text { in } \Omega, \partial_{n} y=u \quad \text { on } \Gamma \\
y_{a} \leq y \leq y_{b} \quad \text { a.e. in } \Omega
\end{array}\right.
$$

with a fixed function $e \in L^{2}(\Omega)$ and a fixed shift control $u_{d} \in L^{2}(\Gamma)$. By substituting $\mathbf{u}=u-u_{d}$ and splitting the state into two components $y(e, u)=y_{e}+G u$, where $y_{e}$ is solution of

$$
\begin{gathered}
A y=e \quad \text { in } \Omega \\
\partial_{n} y=0 \quad \text { on } \Gamma,
\end{gathered}
$$

the problem (2.1) is equivalent to:

$$
\left\{\begin{array}{l}
\text { minimize } \frac{1}{2}\left\|S \mathbf{u}-y_{\Omega}\right\|_{L^{2}(\Omega)}^{2}+\frac{\alpha}{2}\|\mathbf{u}\|_{L^{2}(\Gamma)}^{2}  \tag{2.2}\\
\text { over } \mathbf{u} \in L^{2}(\Gamma) \\
\text { subject to } y_{a}^{\prime} \leq G \mathbf{u} \leq y_{b}^{\prime} \quad \text { a.e. in } \Omega
\end{array}\right.
$$

with $y_{\Omega}=y_{d}-y_{e}-S u_{d}, y_{a}^{\prime}=y_{a}-G u_{d}-y_{e}$ and $y_{b}^{\prime}=y_{b}-G u_{d}-y_{e}$. Therefore, we are justified to concentrate on the simpler problem (P).
2.1. Standard results. It is well known that the operator $G: L^{2}(\Gamma) \rightarrow H^{1}(\Omega)$ is continuous. In the case of a two-dimensional domain, $N=2$, the mapping $G: u \rightarrow y$ is even continuous from $L^{2}(\Gamma)$ into $H^{1}(\Omega) \cap \mathcal{C}(\bar{\Omega})$, see [5]. One can show that, independently of the dimension $N \geq 2$, the problem $(P)$ admits a unique solution $\bar{u} \in L^{2}(\Gamma)$ provided that the admissible set $\left\{u \in L^{2}(\Gamma) \mid y_{a} \leq\right.$ $G(u) \leq y_{b}$ a.e. in $\left.\Omega\right\}$ is not empty. In the rest of the paper, we assume hence that the admissible set for $(P)$ is not empty and denote the optimal solution to $(P)$ by $\bar{u}$ with the associated state $\bar{y}=G \bar{u}$. In the following, we present the optimality system for $(P)$ in a appropriate sense defined in [5].

ThEOREM 2.1 (First-order optimality conditions for $(P)$ ). Let $N=2$ and assume that the following Slater condition is satisfied: There exists a function $u_{0} \in L^{2}(\Gamma)$ (a so-called Slater point) such that

$$
y_{a}(x)<G\left(u_{0}\right)(x)<y_{b}(x) \forall x \in \bar{\Omega}
$$

Then, $\bar{u}$ is optimal for $(P)$ if and only if there exists an adjoint state $p \in W^{1, s}(\Omega)$ for all $1 \leq s<$ $\frac{N}{N-1}$ and Lagrange multipliers $\mu_{a}, \mu_{b} \in \mathcal{C}^{*}(\bar{\Omega})$ such that

$$
\begin{align*}
& A \bar{y}=0 \quad \text { in } \Omega \quad A^{\star} p=\bar{y}-y_{d}+\left(\mu_{b}-\mu_{a}\right)_{\mid \Omega} \quad \text { in } \Omega \\
& \partial_{n} \bar{y}=\bar{u} \quad \text { on } \Gamma, \quad \partial_{n} p=\left(\mu_{b}-\mu_{a}\right)_{\mid \Gamma} \quad \text { on } \Gamma \text {, } \\
& p+\alpha \bar{u}=0 \quad \text { on } \Gamma,  \tag{2.3}\\
& <\mu_{a}, \bar{y}-y_{a}>_{\mathcal{C}^{\star}, \mathcal{C}}=<\mu_{b}, \bar{y}-y_{b}>_{\mathcal{C}^{\star}, \mathcal{C}}=0, \\
& <\mu_{b}, w>_{\mathcal{C}^{\star}, \mathcal{C}} \geq 0, \quad<\mu_{a}, w>_{\mathcal{C}^{\star}, \mathcal{C}} \geq 0 \quad \forall w \in \mathcal{C}(\bar{\Omega}) \text {. }
\end{align*}
$$

Notice that $<\cdot, \cdot>_{\mathcal{C}^{\star}, \mathcal{C}}$ stands for the duality pairing between $\mathcal{C}^{\star}(\bar{\Omega})$ and $\mathcal{C}(\bar{\Omega})$.
3. Regularization. As pointed out in the introduction, we look for controls $u$ in the range of the adjoint operator $S^{\star}$, i.e., we substitute $u=S^{\star} v$ in $(P)$. This idea leads us to an associated problem with a new control $v$ defined on the domain $\Omega$. In this section, we study this problem with different regularization terms, to motivate the final form of our regularized problem.

We start by investigating the operator $S^{\star}$. By definition, $S^{\star}$ is defined from $L^{2}(\Omega)$ to $L^{2}(\Gamma)$. It is represented by $S^{\star} v=w_{\left.\right|_{\Gamma}}$, where $w=w(v) \in H^{1}(\Omega)$ is defined as the solution of

$$
\left\{\begin{array}{l}
A^{\star} w=v \quad \text { in } \quad \Omega  \tag{3.1}\\
\partial_{n} w=0 \quad \text { on } \Gamma .
\end{array}\right.
$$

Hence, for each $v \in L^{2}(\Omega), w_{\left.\right|_{\Gamma}}$ possesses at least the regularity $w_{\left.\right|_{\Gamma}} \in H^{\frac{1}{2}}(\Gamma)$. In convex domains $\Omega$, we obtain even $w_{\left.\right|_{\Gamma}} \in H^{\frac{3}{2}}(\Gamma)$. This higher regularity raises some difficulties, since it implies that $S^{\star}: L^{2}(\Omega) \rightarrow L^{2}(\Gamma)$ is not surjective. However, we have the following result:

Lemma 3.1. $S^{\star}\left(L^{2}(\Omega)\right)$ is dense in $L^{2}(\Gamma)$.
Proof. Assume the contrary: Then there exists a function $d \in L^{2}(\Gamma)$ with $d \neq 0$ such that

$$
\left(d, S^{\star} v\right)_{\Gamma}=0 \quad \forall v \in L^{2}(\Omega) .
$$

This implies that $(S d, v)_{\Omega}=0$ for all $v \in L^{2}(\Omega)$ and hence $y(d)=S d=0$. Since $y=y(d)$ satisfies

$$
\left\{\begin{array}{ccc}
A y= & \text { in } \quad \Omega \\
\partial_{n} y= & d & \text { on } \Gamma,
\end{array}\right.
$$

one finds then $d=0$, which contradicts our assumption.
This result shows that $\bar{u}$ can be approximated by functions $S^{\star} v$. Compared with the regularity of $\bar{u}$, the images of $S^{\star}$ are too smooth so that $\bar{u}$ cannot in general be represented by $\bar{u}=S^{\star} v$ with $v \in L^{2}(\Omega)$. We would need measures $v$ for this representation.
3.1. Auxiliary problem. Substituting $u=S^{\star} v$ in $(P)$, we obtain the following problem
$\left(P_{a u x}\right) \quad\left\{\begin{array}{l}\text { minimize } \tilde{f}(v):=f\left(S^{\star} v\right)=\frac{1}{2}\left\|S S^{\star} v-y_{d}\right\|_{L^{2}(\Omega)}^{2}+\frac{\alpha}{2}\left\|S^{\star} v\right\|_{L^{2}(\Gamma)}^{2} \\ \text { over } v \in L^{2}(\Omega) \\ \text { subject to } y_{a} \leq G S^{\star} v \leq y_{b} \quad \text { a.e. in } \Omega .\end{array}\right.$

As noticed previously, the operator $S^{\star}$ is not surjective and consequently the auxiliary problem $\left(P_{\text {aux }}\right)$ is not necessarily solvable. However, under some assumptions of approximability, we are able to show that $(P)$ and $\left(P_{a u x}\right)$ have the same infimal value.

Definition 3.1 (Approximability and Slater condition). We say that $\bar{u}$ satisfies the approximability condition, if there exists a sequence $\left\{a_{n}\right\}_{n=1}^{\infty}$ in $L^{2}(\Omega)$ such that

$$
\begin{equation*}
\lim _{n \rightarrow \infty} S^{\star} a_{n}=\bar{u} \quad \text { and } \quad y_{a} \leq G S^{\star} a_{n} \leq y_{b} \tag{3.2}
\end{equation*}
$$

for all $n$.
If there exists a function $v_{0} \in L^{\infty}(\Omega)$ satisfying the condition

$$
\begin{equation*}
y_{a}+\delta \leq G S^{\star} v_{0} \leq y_{b}-\delta \tag{3.3}
\end{equation*}
$$

with some constant $\delta>0$, then we say that the Slater condition is satisfied.
Lemma 3.2. For $N=2$, the approximability condition is satisfied if the Slater condition (3.3) is fulfilled.

Proof. Since $S^{\star}\left(L^{2}(\Omega)\right)$ is dense in $L^{2}(\Gamma)$, there exists a sequence $\left\{\tilde{v}_{n}\right\}_{n=1}^{\infty}$ in $L^{2}(\Omega)$ such that

$$
\begin{equation*}
\bar{u}=\lim _{n \rightarrow \infty} S^{\star} \tilde{v}_{n}=\lim _{n \rightarrow \infty} u_{n} \tag{3.4}
\end{equation*}
$$

holds in $L^{2}(\Gamma)$, where $u_{n}:=S^{\star} \tilde{v}_{n}$. Moreover, by the assumption $N=2, G$ is continuous from $L^{2}(\Gamma)$ in $H^{1}(\Omega) \cap \mathcal{C}(\bar{\Omega})$. Thus, $G u_{n}$ converges uniformly to $G \bar{u}$. For this reason, we have

$$
\begin{equation*}
G u_{n}=G \bar{u}+d_{n} \quad \forall n, \tag{3.5}
\end{equation*}
$$

where $\left\{d_{n}\right\}_{n=1}^{\infty}$ is a sequence in $\mathcal{C}(\bar{\Omega})$ converging to zero. Define now $a_{n}=\left(1-t_{n}\right) \tilde{v}_{n}+t_{n} v_{0}$, where $\left\{t_{n}\right\}_{n=1}^{\infty}$ is a sequence of positive numbers tending to zero, given by $t_{n}=\frac{\left\|d_{n}\right\|_{\mathcal{C}(\bar{\Omega})}}{\delta}$. Obviously, due to (3.4), the sequence $\left\{S^{\star} a_{n}\right\}_{n=1}^{\infty}$ converges strongly to $\bar{u}$ in $L^{2}(\Omega)$. By (3.5) and the Slater condition (3.3), we find further for all sufficiently large $n$ :

$$
\begin{aligned}
& G S^{\star} a_{n}=\left(1-t_{n}\right) G \bar{u}+\left(1-t_{n}\right) d_{n}+t_{n} G S^{\star} v_{0} \\
& \leq\left(1-t_{n}\right) y_{b}+\left(1-t_{n}\right) d_{n}+t_{n} y_{b}-t_{n} \delta \\
& \leq y_{b}+\left(1-t_{n}\right)\left\|d_{n}\right\|_{\mathcal{C}(\bar{\Omega})}-t_{n} \delta \\
& =y_{b}-t_{n}\left\|d_{n}\right\|_{\mathcal{C}(\bar{\Omega})} \leq y_{b} .
\end{aligned}
$$

In a similar way, we infer $y_{a} \leq G S^{\star} a_{n}$ for all sufficiently large $n$ and hence $\left\{a_{n}\right\}_{n=1}^{\infty}$ satisfies the property (3.2).

THEOREM 3.1. Assume that the optimal solution $\bar{u}$ satisfies the approximability condition. Then $(P)$ and $\left(P_{\text {aux }}\right)$ have the same infimal value. Furthermore, there exists an infimal sequence $\left\{v_{n}\right\}_{n=1}^{\infty}$ for $\left(P_{\text {aux }}\right)$ such that $\left\{S^{\star} v_{n}\right\}_{\sim}^{\infty}{ }_{n=1}^{\infty}$ converges strongly to $\bar{u}$ in $L^{2}(\Gamma)$.

Proof. By the nonnegativity of $\tilde{f}$, the infimal values of $\left(P_{a u x}\right)$ and $(P)$, denoted by $j_{a u x}$ and $j$, respectively, exist in $\mathbb{R}_{0}^{+}$. Furthermore, since the range of $S^{\star}$ is dense in $L^{2}(\Gamma)$, we have obviously $j \leq j_{a u x}$. On the other hand, by the approximability assumption, there exists a sequence $\left\{a_{n}\right\}_{n=1}^{\infty}$ in $L^{2}(\Omega)$ such that $S^{\star} a_{n}$ is feasible for $(P)$ for sufficiently large $n$ and $\lim _{n \rightarrow \infty} S^{\star} a_{n}=\bar{u}$. Thus, by the continuity of $S$, we have

$$
\lim _{n \rightarrow \infty}\left(\frac{1}{2}\left\|S S^{\star} a_{n}-y_{d}\right\|_{L^{2}(\Omega)}^{2}+\frac{\alpha}{2}\left\|S^{\star} a_{n}\right\|_{L^{2}(\Gamma)}^{2}\right)=\frac{1}{2}\left\|S \bar{u}-y_{d}\right\|_{L^{2}(\Omega)}^{2}+\frac{\alpha}{2}\|\bar{u}\|_{L^{2}(\Gamma)}^{2}=j
$$

and consequently it holds that $j_{\text {aux }} \leq j$.
Let now $\left\{v_{n}\right\}_{n=1}^{\infty}$ be an infimal sequence for $\left(P_{a u x}\right)$. Thanks to the Tikhonov regularization parameter $\alpha>0$, the sequence $\left\{S^{\star} v_{n}\right\}_{n=1}^{\infty}$ is uniformly bounded in $L^{2}(\Omega)$. Thus, there exists a subsequence, denoted w.l.o.g. by $\left\{S^{\star} v_{n}\right\}_{n=1}^{\infty}$, converging weakly to $\tilde{u} \in L^{2}(\Gamma)$. Due to the compactness of the embedding of $H^{1}(\Omega)$ in $L^{2}(\Omega), G S^{\star} v_{n}$ converges strongly to $G \tilde{u}$ in $L^{2}(\Omega)$ as $n \rightarrow \infty$ and consequently $\tilde{u}$ is feasible for $(P)$, i.e.

$$
y_{a} \leq G \tilde{u} \leq y_{b} \text { a.e. in } \Omega .
$$

The optimality of $\tilde{u}$ for $(P)$ follows now from the semicontinuity of $f$ :

$$
f(\tilde{u}) \leq \lim _{n \rightarrow \infty} \inf f\left(S^{\star} v_{n}\right)=j_{a u x}=j
$$

Therefore, by the uniqueness of the optimal solution, we find $\tilde{u}=\bar{u}$. Finally, due to the weak convergence, the strong convergence of $S^{\star} v_{n}$ to $\bar{u}$ can be directly derived from the convergence of $S^{\star} v_{n}$ in norm. $\square$

In the preceding proof, we did not show the convergence of $\left\{v_{n}\right\}_{n=1}^{\infty}$. As yet, we are only able to show the strong convergence of $S^{\star} v_{n}$ to $\bar{u}$ for the auxiliary problem ( $P_{\text {aux }}$ ). Still, this is not satisfactory due to the possible unsolvability of $\left(P_{a u x}\right)$. To overcome this difficulty, we next study different kinds of regularization and their specific properties. Finally, we end up with our Lavrentiev-regularized problem. We start with the following problem:

$$
\left\{\begin{array}{l}
\operatorname{minimize} \tilde{g}(v):=\frac{1}{2}\left\|S S^{\star} v-y_{d}\right\|_{L^{2}(\Omega)}^{2}+\frac{\alpha}{2}\left\|S^{\star} v\right\|_{L^{2}(\Gamma)}^{2}+\frac{\varepsilon}{2}\|v\|_{L^{2}(\Omega)}^{2} \\
\text { over } v \in L^{2}(\Omega) \\
\text { subject to } y_{a} \leq G S^{\star} v \leq y_{b} \quad \text { a.e. in } \Omega
\end{array}\right.
$$

Lemma 3.3. Assume that the admissible set $\tilde{V}_{a d}=\left\{v \in L^{2}(\Omega) \mid y_{a} \leq G S^{\star} v \leq y_{b}\right.$ a.e. in $\left.\Omega\right\}$ is not empty. Then, for every $\varepsilon>0,\left(P^{\varepsilon}\right)$ admits a unique solution.

Proof. The infimal value of $\left(P^{\varepsilon}\right)$ exists in $\mathbb{R}_{0}^{+}$, since the objective functional $\tilde{g}$ is nonnegative. Let now $\left\{v_{n}\right\}$ be an infimal sequence for $\left(P^{\varepsilon}\right)$. Due to the regularization parameter $\frac{\varepsilon}{2}$ in $\tilde{g}, v_{n}$ is uniformly bounded in $L^{2}(\Omega)$. Consequently, there exists a subsequence of $v_{n}$ converging weakly to $\tilde{v} \in L^{2}(\Omega)$. Owing to the continuity of $G S^{\star}$, the weak limit $\tilde{v}$ belongs obviously to $\tilde{V}_{a d}$. Thus, by the lower semicontinuity of $\tilde{g}, \tilde{v}$ is optimal for $\left(P^{\varepsilon}\right) \square$
3.2. Lavrentiev type regularization. As mentioned previously, our aim is to propose a Lavrentiev type regularization applied to the boundary control problem. We regularize now $(P)$ in the following way:

$$
\left\{\begin{array}{l}
\operatorname{minimize} \tilde{g}(v):=\frac{1}{2}\left\|S S^{\star} v-y_{d}\right\|_{L^{2}(\Omega)}^{2}+\frac{\alpha}{2}\left\|S^{\star} v\right\|_{L^{2}(\Gamma)}^{2}+\frac{\varepsilon}{2}\|v\|_{L^{2}(\Omega)}^{2} \\
\text { over } v \in L^{2}(\Omega) \\
\text { subject to } y_{a} \leq \lambda v+G S^{\star} v \leq y_{b} \quad \text { a.e. in } \Omega
\end{array}\right.
$$

with some regularization parameters $\lambda>0$ and $\varepsilon(\lambda) \geq 0$. We mention here some reasons for regularizing the problem in that way: Without any restriction on the dimension $N$ and without the assumption on approximability and Slater conditions, we always have the solvability of $\left(P_{\lambda}\right)$ and additionally, we are able to show that the associated Lagrange multipliers exist and belong to $L^{2}(\Omega)$. However, for the convergence of the regularized solution to the solution of $(P)$ in the case
of vanishing regularization parameters, one has to restrict again the theory to the two-dimensional case, $N=2$. We require this since for $N=2$, the operator $G$ is continuous from $L^{2}(\Omega)$ to $\mathcal{C}(\bar{\Omega}) \cap H^{1}(\Omega)$.

THEOREM 3.2. Let $\lambda>0$ and $\varepsilon(\lambda) \geq 0$ be arbitrarily fixed. Then, the regularized problem $\left(P_{\lambda}\right)$ admits a solution and the solution is unique if $\varepsilon(\lambda) \neq 0$.

Proof. In the proof, we consider $G$ again as mapping with range in $L^{2}(\Omega)$, i.e. we subtitute $S$ for $G$. This does not change the admissible set. First of all, we have to show that the admissible set $V_{a d}:=\left\{v \in L^{2}(\Omega) \mid y_{a} \leq \lambda v+S S^{\star} v \leq y_{b}\right.$ a.e. in $\left.\Omega\right\}$ is not empty. To this purpose, we consider the operator $\lambda I+S S^{\star}$, where $I: L^{2}(\Omega) \rightarrow L^{2}(\Omega)$ denotes the identity operator in $L^{2}(\Omega)$. Since $S S^{\star}: L^{2}(\Omega) \rightarrow L^{2}(\Omega)$ is positive semidefinite and $\lambda>0$ holds by assumption, the equation $\lambda v+S S^{\star} v=0$ admits only the trivial solution $v=0$. Hence, by the Fredholm alternative, the compactness of $S S^{\star}$ implies the existence of the inverse operator $\left(\lambda I+S S^{\star}\right)^{-1}$. From this, we infer that $V_{a d}$ is not empty: For instance, we have $\left(\lambda I+S S^{\star}\right)^{-1} y_{b} \in V_{a d}$.

Next, since the objective functional in $\left(P_{\lambda}\right)$ is nonnegative, the infimum in $\left(P_{\lambda}\right)$ exists in $\mathbb{R}_{0}^{+}$. Let now $\left\{v_{n}\right\}_{n=1}^{\infty}$ be an infimal sequence for $\left(P_{\lambda}\right)$. By the presence of the cost parameter $\alpha>0, S^{\star} v_{n}$ is uniformly bounded in $L^{2}(\Omega)$ and hence we can find a subsequence $\left\{S^{\star} v_{n_{j}}\right\}_{j=1}^{\infty}$ of $\left\{S^{\star} v_{n}\right\}_{n=1}^{\infty}$ such that $S^{\star} v_{n_{j}} \rightharpoonup u^{\prime}$. Subsequently, invoking again the compactness of $S$, we infer the strong convergence of $S S^{\star} v_{n_{j}}$ to $S u^{\prime}$. Moreover, since $v_{n_{j}} \in V_{a d}$ for all $j$, the strong convergence of $S S^{\star} v_{n_{j}}$ ensures then the uniform boundedness property of $v_{n_{j}}$. Thus, we find again a subsequence of $v_{n_{j}}$ converging weakly to $\bar{v}_{\lambda} \in L^{2}(\Omega)$. This weak limit $\bar{v}_{\lambda}$ is clearly feasible and finally, by the lower semicontinuity of the objective functional in $\left(P_{\lambda}\right), \bar{v}_{\lambda}$ is optimal. For $\varepsilon(\lambda)>0, \tilde{g}$ is strictly convex and consequently we obtain the uniqueness of $\bar{v}_{\lambda}$.

Next, setting $z=\lambda v+S S^{\star} v$ and hence, $v=\left(\lambda I+S S^{\star}\right)^{-1} z=R z,\left(P_{\lambda}\right)$ is equivalent to the following control problem:

$$
\left\{\begin{array}{l}
\operatorname{minimize} \mathbf{g}(z):=\tilde{g}(R z)  \tag{z}\\
\text { over } z \in L^{2}(\Omega) \\
\text { subject to } y_{a} \leq z \leq y_{b} \quad \text { a.e. in } \Omega
\end{array}\right.
$$

In this way, we have just transformed $\left(P_{\lambda}\right)$ into a control problem with a simple box constraint. Subsequently, by standard arguments, cf. [15], the following optimality system can easily be shown.

Theorem 3.3 (First-order optimality conditions). Let $\lambda>0$ be arbitrarily fixed and let $\bar{v}_{\lambda}$ be an optimal solution to ( $P_{\lambda}$ ). Moreover, we set $\bar{y}_{\lambda}=G S^{\star} \bar{v}_{\lambda}$ and $S^{\star} \bar{v}_{\lambda}=w_{\left.\right|_{\Gamma}}$, with $w=w\left(\bar{v}_{\lambda}\right) \in H^{1}(\Omega)$ solution of (3.1). Then, there exist Lagrange multipliers $\mu_{\lambda}^{a}, \mu_{\lambda}^{b} \in L^{2}(\Omega)$ and adjoint states $p, q \in H^{1}(\Omega)$ such that the following optimality system is satisfied:

$$
\begin{align*}
& \begin{array}{lllllllll}
A \bar{y}_{\lambda} & = & 0 & \text { in } & \Omega & A^{\star} w & = & \bar{v}_{\lambda} & \text { in } \\
\Omega \\
\partial_{n} \bar{y}_{\lambda} & = & & \Omega \\
\text { on } & \Gamma, & { }_{\partial n} w & = & 0 & \text { on } & \Gamma,
\end{array}  \tag{3.6}\\
& \begin{array}{lllllll}
A^{\star} p & = & \bar{y}_{\lambda}-y_{d}+\mu_{\lambda}^{b}-\mu_{\lambda}^{a} & \text { in } & \Omega & A^{\star} q & =0 \\
\text { on } & \Gamma, & A^{n} & \Omega \\
\partial_{n} q & = & \alpha w+p & \text { on } & \Gamma, \\
\partial_{n} p & =0 & & & & &
\end{array}
\end{align*}
$$

$$
\begin{align*}
& \varepsilon(\lambda) \bar{v}_{\lambda}+q+\lambda\left(\mu_{\lambda}^{b}-\mu_{\lambda}^{a}\right)=0  \tag{3.8}\\
& y_{a} \leq \lambda \bar{v}_{\lambda}+\bar{y}_{\lambda} \leq y_{b} \quad \text { a.e. in } \Omega \tag{3.9}
\end{align*}
$$

$$
\begin{gather*}
\mu_{\lambda}^{a} \geq 0, \mu_{\lambda}^{b} \geq 0  \tag{3.10}\\
\left(\mu_{\lambda}^{a}, y_{a}-\lambda \bar{v}_{\lambda}-\bar{y}_{\lambda}\right)_{L^{2}(\Omega)}=\left(\mu_{\lambda}^{b}, \lambda \bar{v}_{\lambda}+\bar{y}_{\lambda}-y_{b}\right)_{L^{2}(\Omega)}=0 .
\end{gather*}
$$

3.3. Pass to the limit $\lambda \rightarrow 0$. We study now the convergence of the solution to the regularized problem in the case of vanishing Lavrentiev parameter $\lambda$.

Assumption 3.1. The regularization parameter $\varepsilon=\varepsilon(\lambda)$ satisfies

$$
\begin{equation*}
\varepsilon=\sigma_{0} \lambda^{1+\sigma_{1}} \tag{3.11}
\end{equation*}
$$

with some constants $\sigma_{0}>0$ and $0 \leq \sigma_{1}<1$. Moreover, there exists a function $v_{0} \in L^{\infty}(\Omega)$ such that

$$
\begin{equation*}
y_{a}+\delta \leq G\left(\bar{u}+S^{\star} v_{0}\right) \leq y_{b}-\delta \tag{3.12}
\end{equation*}
$$

is satisfied with some constant $\delta>0$.
Let $\left\{\lambda_{n}\right\}_{n=1}^{\infty}$ be a sequence of positive real numbers converging to zero and by $\left\{v_{n}\right\}_{n=1}^{\infty}$, we denote the sequence of optimal solutions to $\left(P_{\lambda_{n}}\right)$. The presence of the Tikhonov parameter $\alpha>0$ in $\left(P_{\lambda_{n}}\right)$ ensures the boundedness of the sequence $\left\{S^{\star} v_{n}\right\}_{n=1}^{\infty}$ in $L^{2}(\Gamma)$. For this reason, we can find a subsequence of $\left\{S^{\star} v_{n}\right\}_{n=1}^{\infty}$, denoted w.l.o.g. again by $\left\{S^{\star} v_{n}\right\}_{n=1}^{\infty}$, converging weakly to some $\tilde{u} \in L^{2}(\Gamma)$. Our goal now is to show that $\tilde{u}$ minimizes the original unregularized problem. For this purpose, we should show first the feasibility of $\tilde{u}$ for $(P)$, i.e., $y_{a} \leq G \tilde{u} \leq y_{b}$ a.e. in $\Omega$.

Lemma 3.4. Let Assumption 3.1, (3.11), be satisfied. Then the weak limit $\tilde{u}$ of the sequence $\left\{S^{\star} v_{n}\right\}_{n=1}^{\infty}$ defined above is feasible for $(P)$.

Proof. We know that it holds

$$
y_{a} \leq \lambda_{n} v_{n}+G S^{\star} v_{n} \leq y_{b} \quad \forall n
$$

Therefore, it suffices to show that $\lambda_{n} v_{n}$ converges to zero. Clearly, by the presence of the regularization parameter $\varepsilon(\lambda)>0$ in the objective functional $\tilde{g}$, one has:

$$
\frac{\varepsilon\left(\lambda_{n}\right)}{2}\left\|v_{n}\right\|_{L^{2}(\Omega)}^{2} \leq c \quad \forall n
$$

with some constant $c>0$ and hence

$$
\frac{\varepsilon\left(\lambda_{n}\right)}{2 \lambda_{n}^{2}}\left\|\lambda_{n} v_{n}\right\|_{L^{2}(\Omega)}^{2} \leq c \quad \forall n
$$

¿From Assumption 3.1, we infer then

$$
\begin{equation*}
\left\|\lambda_{n} v_{n}\right\|_{L^{2}(\Omega)}^{2} \leq \frac{2 c \lambda_{n}^{2}}{\varepsilon\left(\lambda_{n}\right)} \leq \lambda_{n}^{1-\sigma_{1}} \frac{2 c}{\sigma_{0}} \tag{3.13}
\end{equation*}
$$

This implies $\lambda_{n} v_{n} \rightarrow 0$ in $L^{2}(\Omega)$ as $n \rightarrow \infty$ and hence the Lemma is shown.
ThEOREM 3.4. Let $N=2$. Then, under Assumption 3.1, the sequence $\left\{S^{\star} v_{n}\right\}$ converges strongly in $L^{2}(\Gamma)$ to the optimal solution of the unregularized problem $(P)$.

Proof. Since $\mathcal{C}(\bar{\Omega})$ is dense in $L^{2}(\Omega)$ and due to Lemma 3.1, we can find a sequence $\left\{z_{k}\right\}_{k=1}^{\infty}$ in $\mathcal{C}(\bar{\Omega})$ such that

$$
\begin{equation*}
\left\|\bar{u}-S^{\star} z_{k}\right\|_{L^{2}(\Gamma)} \leq \frac{1}{k} \quad \forall k \in \mathbb{N} . \tag{3.14}
\end{equation*}
$$

Moreover, continuity and linearity of $G$ from $L^{2}(\Gamma)$ to $H^{1}(\Omega) \cap \mathcal{C}(\bar{\Omega})$ ensure the existence of a real positive number $c_{0}$ such that

$$
\begin{equation*}
\left\|G\left(\bar{u}-S^{\star} z_{k}\right)\right\|_{\mathcal{C}(\bar{\Omega})} \leq c_{0}\left\|\bar{u}-S^{\star} z_{k}\right\|_{L^{2}(\Gamma)} \leq c_{0} \frac{1}{k} \quad \forall k \in \mathbb{N} \tag{3.15}
\end{equation*}
$$

We define now an auxiliary sequence $\left\{v_{k}^{0}\right\}_{k=1}^{\infty}$ in $L^{\infty}(\Omega)$ by

$$
\begin{equation*}
v_{k}^{0}=z_{k}+\frac{2 c_{0}}{k \delta} v_{0}=z_{k}+\frac{c_{1}}{k} v_{0} \tag{3.16}
\end{equation*}
$$

where $v_{0}$ satisfies (3.12) and $c_{1}=\frac{2 c_{0}}{\delta}$. In a view of (3.14), the definition above implies the strong convergence of $S^{\star} v_{k}^{0}$ in $L^{2}(\Gamma)$ to $\bar{u}$, i.e.

$$
\begin{equation*}
\lim _{k \rightarrow \infty}\left\|S^{\star} v_{k}^{0}-\bar{u}\right\|_{L^{2}(\Gamma)}=0 \tag{3.17}
\end{equation*}
$$

First, we show that, for every $k \in \mathbb{N}$ with $k \geq c_{1}$, one can find an index number $n_{k} \in \mathbb{N}$ such that $v_{k}^{0}$ is feasible for $\left(P_{\lambda_{n}}\right)$ for all $n \geq n_{k}$. Let now $k \in \mathbb{N}$ with $k \geq c_{1}$ be arbitrarily fixed. Then, by our assumptions, it holds that

$$
\begin{align*}
& \lambda_{n} v_{k}^{0}+G S^{\star} v_{k}^{0}=\lambda_{n} v_{k}^{0}+G\left(S^{\star} z_{k}-\bar{u}\right)+\left(1-\frac{c_{1}}{k}\right) G \bar{u}+\frac{c_{1}}{k}\left(G S^{\star} v_{0}+G \bar{u}\right) \\
& \leq \lambda_{n}\left\|v_{k}^{0}\right\|_{L^{\infty}(\Omega)}+c_{0}\left\|S^{\star} z_{k}-\bar{u}\right\|_{L^{2}(\Gamma)}+\left(1-\frac{c_{1}}{k}\right) y_{b}+\frac{c_{1}}{k}\left(y_{b}-\delta\right) \\
& \leq y_{b}+\left(\lambda_{n}\left\|v_{k}^{0}\right\|_{L^{\infty}(\Omega)}+\frac{c_{0}}{k}-2 \frac{c_{0}}{k}\right)  \tag{3.18}\\
& =y_{b}+\left(\lambda_{n}\left\|v_{k}^{0}\right\|_{L^{\infty}(\Omega)}-\frac{c_{0}}{k}\right) .
\end{align*}
$$

Since $\lambda_{n}$ converges to zero, we can find then an index $n_{k} \in \mathbb{N}$ such that $\lambda_{n}\left\|v_{k}^{0}\right\|_{L^{\infty}(\Omega)} \leq \frac{c_{0}}{k}$ for all $n \geq n_{k}$. Inserting this in (3.18), we obtain

$$
\lambda_{n} v_{k}^{0}+G S^{\star} v_{k}^{0} \leq y_{b} \quad \forall n \geq n_{k}
$$

In the same way, one derives $y_{a} \leq \lambda_{n} v_{k}^{0}+G S^{\star} v_{k}^{0}$ for all $n \geq n_{k}$. This implies the feasibility of $v_{k}^{0}$ for all $\left(P_{\lambda_{n}}\right)$ with $n \geq n_{k}$.

Since $v_{n}$ is optimal to $\left(P_{\lambda_{n}}\right)$ for each $n$, we infer further that

$$
f\left(S^{\star} v_{n}\right) \leq f\left(S^{\star} v_{n}\right)+\frac{\varepsilon_{n}}{2}\left\|v_{n}\right\|_{L^{2}(\Omega)}^{2} \leq f\left(S^{\star} v_{k}^{0}\right)+\frac{\varepsilon_{n}}{2}\left\|v_{k}^{0}\right\|_{L^{2}(\Omega)}^{2} \forall n \geq n_{k}
$$

Hence, due to the lower semicontinuity of $f$, one finds by passing to the limit $n \rightarrow \infty$

$$
f(\tilde{u}) \leq \lim _{n \rightarrow \infty} \inf f\left(S^{\star} v_{n}\right) \leq \lim _{n \rightarrow \infty} \sup f\left(S^{\star} v_{n}\right) \leq f\left(S^{\star} v_{k}^{0}\right)
$$

where $\tilde{u}$ is the weak limit introduced in Lemma 3.4. Finally, letting $k$ pass to infinity, we infer from (3.17) that

$$
f(\tilde{u}) \leq \lim _{k \rightarrow \infty} f\left(S^{\star} v_{k}^{0}\right)=f(\bar{u})
$$

Hence, we have shown the optimality of $\tilde{u}$ to $(P)$ and again, due to the uniqueness of $\bar{u}$, we have $\tilde{u}=\bar{u}$. Notice that the latter equality $\lim _{n \rightarrow \infty} f\left(S^{\star} v_{n}\right)=f(\bar{u})$ implies the convergence of $\left\{S^{\star} v_{n}\right\}$ in norm and hence, together with the weak convergence, the strong convergence of $\left\{S^{\star} v_{n}\right\}$ to $\tilde{u}$ is verified.
4. Semismooth Newton algorithm. Based on the experience in [11, 8, 9], the semismooth Newton method is quite efficient when dealing with state-constrained optimal control problems. Mainly due to its superlinear convergence and mesh-independence properties, the method is highly efficient in many applications. This was analyzed and verified numerically, quite recently, also for a similar Lavrentiev regularization technique applied to distributed optimal control problem, [9]. Our
goal in this section is to present a semismooth Newton algorithm based on the first-order optimality conditions (3.6)-(3.10) for the regularized problem $\left(P_{\lambda}\right)$. The analysis for the mesh-independence principle is not included here. This is a subject of our ongoing research. Following [6, 8], we define now the Newton generalized differentiability.

Definition 4.1. Let $X, Y$ be Banach spaces and $U$ be an open set in $X$. A mapping $F: U \rightarrow Y$ is said to be semismooth (or Newton differentiable) in $U$ if there exists a (possibly set-valued) mapping $\partial F: U \rightrightarrows \mathscr{L}(X, Y)$ such that

$$
\begin{equation*}
\sup _{V \in \partial F(x+s)}\|F(x+s)-F(x)-V s\|_{Y}=o\left(\|s\|_{X}\right) \quad a s\|s\|_{X} \rightarrow 0 \tag{4.1}
\end{equation*}
$$

is satisfied for all $x \in U$. We call $\partial F$ the Newton differential, and its elements $V$ are referred to as Newton maps.

In the following, we derive a semismooth Newton algorithm based on the concept above. Let us start by reformulating the complementarity system (3.10) in the optimality conditions by the following max-formulation.

LEmma 4.1. The complementarity conditions (3.10) are equivalent to:

$$
\begin{align*}
\mu_{\lambda}^{a} & =\max \left(0, \mu_{\lambda}^{a}-\mu_{\lambda}^{b}+\gamma\left(y_{a}-\lambda \bar{v}_{\lambda}-\bar{y}_{\lambda}\right)\right),  \tag{4.2}\\
\mu_{\lambda}^{b} & =\max \left(0, \mu_{\lambda}^{b}-\mu_{\lambda}^{a}+\gamma\left(\lambda \bar{v}_{\lambda}+\bar{y}_{\lambda}-y_{b}\right)\right), \tag{4.3}
\end{align*}
$$

for arbitrarily fixed $\gamma>0$.
We sketch the proof for the convenience of the reader. First, the complementarity conditions (3.10) imply obviously (4.2)-(4.3). Furthermore, (4.2)-(4.3) imply the nonnegativity of $\mu_{\lambda}^{a}$ and $\mu_{\lambda}^{b}$. Let now $\mathcal{A}$ be a set in $\Omega$ such that $\mu_{\lambda}^{a}=\mu_{\lambda}^{a}-\mu_{\lambda}^{b}+\gamma\left(y_{a}-\lambda \bar{v}_{\lambda}-\bar{y}_{\lambda}\right)$ on $\mathcal{A}$. Thus $\mu_{\lambda}^{b}=\gamma\left(y_{a}-\lambda \bar{v}_{\lambda}-\bar{y}_{\lambda}\right)$ on $\mathcal{A}$. Furthermore, (4.3) implies

$$
\mu_{\lambda}^{b}=\max \left(0,-\mu_{\lambda}^{a}+\gamma\left(y_{a}-y_{b}\right)\right)=0 \quad \text { on } \mathcal{A} .
$$

Therefore, $y_{a}-\lambda \bar{v}_{\lambda}-\bar{y}_{\lambda}=0$ on $\mathcal{A}$. For this reason, we have

$$
\left(\mu_{\lambda}^{a}, y_{a}-\lambda \bar{v}_{\lambda}-\bar{y}_{\lambda}\right)_{L^{2}(\Omega)}=0 .
$$

In the same way, we infer

$$
\left(\mu_{\lambda}^{b}, \lambda \bar{v}_{\lambda}+\bar{y}_{\lambda}-y_{b}\right)_{L^{2}(\Omega)}=0
$$

and hence (4.2)-(4.3) imply the complementarity system (3.10).
In [8], the investigation of the Newton differentiability of the maximum-operation has been analyzed. As demonstrated, a class of corresponding Newton maps for the maximum operator $\mathbb{M}(z)=\max (0, z)$ are given by

$$
\partial \mathbb{M}_{\xi}(z)(x)=\left\{\begin{array}{l}
1 \text { if } z(x)>0  \tag{4.4}\\
0 \text { if } z(x)<0 \\
\xi \text { if } z(x)=0
\end{array}\right.
$$

with arbitrarily fixed $\xi \in \mathbb{R}$, provided that $\mathbb{M}$ is defined from $L^{q_{2}}(\Omega)$ to $L^{q_{1}}(\Omega)$ with $1 \leq q_{1}<q_{2} \leq$ $\infty$.

Setting now a special choice $\gamma:=\varepsilon(\lambda) / \lambda^{2}$ in (4.2)-(4.3), short computations show that due to the equation (3.8) in the optimality system, i.e.,

$$
\varepsilon(\lambda) \bar{v}_{\lambda}+q+\lambda\left(\mu_{\lambda}^{b}-\mu_{\lambda}^{a}\right)=0
$$

the complementarity system is equivalent to

$$
\begin{align*}
& \mu_{\lambda}^{a}=\max \left(0, \frac{1}{\lambda} q+\frac{\varepsilon(\lambda)}{\lambda^{2}}\left(y_{a}-\bar{y}_{\lambda}\right)\right),  \tag{4.5}\\
& \mu_{\lambda}^{b}=\max \left(0,-\frac{1}{\lambda} q+\frac{\varepsilon(\lambda)}{\lambda^{2}}\left(\bar{y}_{\lambda}-y_{b}\right)\right) . \tag{4.6}
\end{align*}
$$

The maximum operators above are defined from $L^{q_{2}}(\Omega)$ to $L^{q_{1}}(\Omega)$ with $1 \leq q_{1}<q_{2}$ and thus the application of semismooth Newton method is justified. Our algorithm is based on the particular choice $\xi=0$ for the Newton map. Then, the semismooth Newton algorithm is equivalent to an active-set-strategy. The complete algorithm is defined by the following steps, cf. also $[8,9]$ for the details.

Algorithm 4.1.
(i) Initialization: Choose initial data $q^{0}, y^{0} \in L^{2}(\Omega)$ and set $l=0$.
(ii) Set the active and inactive sets:

$$
\begin{aligned}
\mathcal{A}_{a}^{l} & =\left\{x \in \Omega: \frac{1}{\lambda} q^{l}(x)+\frac{\varepsilon(\lambda)}{\lambda^{2}}\left(y_{a}(x)-y^{l}(x)\right)>0 \text { a.e. in } \Omega\right\}, \\
\mathcal{A}_{b}^{l} & =\left\{x \in \Omega:-\frac{1}{\lambda} q^{l}(x)+\frac{\varepsilon(\lambda)}{\lambda^{2}}\left(y^{l}(x)-y_{b}(x)\right)>0 \text { a.e. in } \Omega\right\}, \\
\mathcal{I}^{l} & =\Omega \backslash\left(\mathcal{A}_{a}^{l} \cup \mathcal{A}_{b}^{l}\right) .
\end{aligned}
$$

(iii) Find the solution $\left(y^{l+1}, q^{l+1}, p^{l+1}, w^{l+1}, v^{l+1}, \mu_{a}^{l+1}, \mu_{b}^{l+1}\right)$ of

$$
\begin{aligned}
& A y^{l+1}=0 \quad \text { in } \Omega, \quad A^{\star} w^{l+1}=v^{l+1} \quad \text { in } \Omega \text {, } \\
& \partial_{n} y^{l+1}=w^{l+1} \text { on } \Gamma, \quad \partial_{n} w^{l+1}=0 \quad \text { on } \Gamma \text {, } \\
& A^{\star} p^{l+1}=y^{l+1}-y_{d}+\mu_{b}^{l+1}-\mu_{a}^{l+1} \quad \text { in } \quad \Omega, \\
& \partial_{n} p^{l+1}=0 \quad \text { on } \Gamma \text {, } \\
& A^{\star} q^{l+1}=0 \quad \text { in } \Omega, \\
& \partial_{n} q^{l+1}=\alpha w^{l+1}+p^{l+1} \quad \text { on } \Gamma, \\
& \varepsilon(\lambda) v^{l+1}+q^{l+1}+\lambda\left(\mu_{b}^{l+1}-\mu_{a}^{l+1}\right)=0, \\
& \mu_{a}^{l+1}=\frac{1}{\lambda} q^{l+1}+\frac{\varepsilon(\lambda)}{\lambda^{2}}\left(y_{a}-y^{l+1}\right) \quad \text { on } \quad \mathcal{A}_{a}^{l}, \quad \mu_{a}^{l+1}=0 \text { on } \mathcal{I}^{l} \cup \mathcal{A}_{b}^{l} \text {, } \\
& \mu_{b}^{l+1}=-\frac{1}{\lambda} q^{l+1}+\frac{\varepsilon(\lambda)}{\lambda^{2}}\left(y^{l+1}-y_{b}\right) \quad \text { on } \quad \mathcal{A}_{b}^{l}, \quad \mu_{b}^{l+1}=0 \text { on } \mathcal{I}^{l} \cup \mathcal{A}_{a}^{l} \text {. }
\end{aligned}
$$

(iv) Stop or set $l=l+1$ and go to (ii).

Unless otherwise specified, we initialize Algorithm 4.1 with $y^{0}=q^{0}=0$ and terminate it using the stopping criterion: $\mathcal{A}_{a}^{n}=\mathcal{A}_{a}^{n-1}$ and $\mathcal{A}_{b}^{n}=\mathcal{A}_{b}^{n-1}$.
5. Numerical Experiments. Our numerical report is splitted into two parts. First, we confirm the numerical reliability of our regularization approach to solve some classes of stateconstrained optimal boundary control problems. In particular, we aim at investigating the influence of the regularization parameter on the algorithm. To this purpose, a test example with analytically known solution to the problem (2.1) will be considered. As pointed out in Remark 2.1, our theory is applicable also to this more general problem (2.1). By means of this example, the numerical approximation of our method as well as its convergence behavior will be analyzed. We also study briefly a nested iteration technique, based on a multigrid concept, to gain a higher efficiency of the algorithm.

Second, we compare our technique based on the "optimize-then-discretize" concept with the standard numerical optimization code QUADPROG of the MATLAB optimization toolbox applied to the discretized problem ("discretize-then-optimize"). We mainly aim at showing that our method exhibits a reasonable performance. We do not intend to compare the concepts "optimize-thendiscretize" and "discretize-then-optimize", since this would require various test runs with common available nonlinear optimization codes. A detailed study of the method "discretize-then-optimize" was carried out by Maurer and Mittelmann [13].

Our experience showed that the semismooth Newton method applied to the Lavrentiev type regularization was as efficient as QUADPROG and partially even more advantageous. We point out that all the numerical computations in this paper were carried out on a PC with a $250-\mathrm{GHz}$ AMD processor and a 16-gigabyte memory.
5.1. Discretization. As noticed early, the regularization technique that we propose here is based on the optimize-then-discretize strategy. In the following, we explain the discretization of Algorithm 4.1 and call later the algorithm under this discretization "OTD".

Throughout the experiment, we use for simplicity the unit square domain $\Omega=(0,1) \times(0,1)$ and set $A=-\Delta+I$. We discretize $\Omega$ by a regular Friedrichs-Keller triangulation with mesh size $h$ and the mesh on $\Gamma$ is induced by that on $\Omega$. The partial differential equations (PDEs) are approximated by the finite element method. The state space $H^{1}(\Omega) \cap \mathcal{C}(\bar{\Omega})$ is discretized by the span of the standard finite element basis $\left\{\phi_{h}^{1}(x), \ldots, \phi_{h}^{N_{h}}(x)\right\}$ consisting of the piecewise linear and continuous hat functions on $\bar{\Omega}$. Analogously, we define the discrete control space with the standard finite element basis $\left\{\psi_{h}^{1}(x), \ldots, \psi_{h}^{M_{h}}(x)\right\}$, composed of the piecewise linear and continuous hat functions defined on the boundary $\Gamma$. Hence, the state equation is approximated by the system of linear equations

$$
A_{h} y_{h}=B_{h} u_{h}+M_{h} e_{h}
$$

where the matrices $A_{h}, M_{h} \in \mathbb{R}^{N_{h} \times N_{h}}$ and $B_{h} \in \mathbb{R}^{N_{h} \times M_{h}}$ are given by

$$
\begin{aligned}
\left(A_{h}\right)_{i j} & =\left(\nabla \phi_{h}^{i}, \nabla \phi_{h}^{j}\right)_{L^{2}(\Omega)}+\left(\phi_{h}^{i}, \phi_{h}^{j}\right)_{L^{2}(\Omega)} \\
\left(B_{h}\right)_{i, j} & =\left(\phi_{h}^{i}, \psi_{h}^{j}\right)_{L^{2}(\Gamma)} \\
\left(M_{h}\right)_{i j} & =\left(\phi_{h}^{i}, \phi_{h}^{j}\right)_{L^{2}(\Omega)}
\end{aligned}
$$

Here, the vectors $y_{h} \in \mathbb{R}^{N_{h}}, u_{h} \in \mathbb{R}^{M_{h}}$ and $e_{h} \in \mathbb{R}^{N_{h}}$ serve for the discrete approximations of the state, the control and the fixed function $e$, respectively, with mesh size $h$. For instance, $y_{h_{i}}$ is the numerical approximation of the value $y\left(x_{i}\right)$ in the node point $x_{i}$ that is associated with the ansatz function $\phi_{h}^{i}$. The remaining PDEs in the optimality system associated with $\left(P_{\lambda}\right)$ are analogously
discretized. The active sets are discretized by the approximated values of corresponding functions at the nodes. For instance, the discretization of $\mathcal{A}_{a}$ is given by

$$
\left\{i: \frac{1}{\lambda} q_{h_{i}}+\frac{\varepsilon(\lambda)}{\lambda^{2}}\left(y_{a, h_{i}}-y_{h_{i}}\right)>0\right\}
$$

where the vectors $q_{h}, y_{a, h} \in \mathbb{R}^{N_{h}}$ serve for the discrete approximations of $q$ and $y_{a}$, respectively, with mesh size $h$. Algorithm 4.1 is implemented in this way.
5.2. Test example 1. To construct the example, we first define the optimal state $y_{o p t}$, the adjoint state $p_{\text {opt }}$, the upper bound $\psi$ and the fixed function $e$ by

$$
\begin{aligned}
y_{o p t}(x) & =\frac{2}{\pi} \sin \left(\pi x_{1}\right) \sin \left(\pi x_{2}\right) \\
p_{o p t}(x) & =-0.5 \\
\psi(x) & =\max \left(\frac{1}{\pi}, y_{o p t}(x)\right) \\
e(x) & =-\Delta y_{o p t}(x)+y_{o p t}(x)
\end{aligned}
$$

Notice that, in our test examples, we only consider an upper bound constraint

$$
y_{o p t}(x) \leq \psi(x) \quad \forall x \in \bar{\Omega}
$$

Clearly, our theory applies to this case as well. Short computations show that

$$
\mu_{o p t}(x):=\left\{\begin{array}{lll}
1.7 & \text { if } & y(x)>\frac{1}{\pi} \\
0 & \text { if } & y(x) \leq \frac{1}{\pi}
\end{array}\right.
$$

fulfills the complementarity slackness condition for $(P)$. Setting for the desired state

$$
y_{d}=y_{o p t}-p_{o p t}+\mu_{o p t}
$$

$\mu_{\text {opt }}$ can serve for the Lagrange multiplier associated with $(P)$. Next, by computing the normal derivative of $y_{o p t}$, one obtains the optimal boundary control $u_{o p t}$, which is identical on all edges of $\Omega$. For example, on the lower boundary of $\Omega, u_{o p t}=-2 \sin \left(\pi x_{1}\right)$. Finally, for the cost parameter $\alpha$ and the desired control $u_{d}$, we set:

$$
\begin{array}{ccc}
\alpha & = & 10^{-2} \\
u_{d} & = & u_{o p t}+\frac{1}{\alpha} p_{o p t}{ }_{\Gamma}
\end{array}
$$

For the choice of the parameter $\varepsilon=\varepsilon(\lambda)$, we select throughout the numerical test $\varepsilon(\lambda)=\lambda^{1+\frac{1}{2}}$. This satisfies clearly Assumption 3.1.

Our aim now consists of investigating the numerical approximations based on Algorithm 4.1 in the case of vanishing Lavrentiev parameter $\lambda$. In Table 5.1, we report on the numerical results when solving the problem utilizing Algorithm 4.1. We found out in our test runs $(h=1 / 128)$ that the problem is becoming harder to be solved for decreasing $\lambda$. Observing the second and third columns of Table 5.1, we notice that the distance to the optimal solution is quite satisfactory and getting smaller with respect to decreasing Lavrentiev parameter $\lambda$. This confirms the result of

Table 5.1
Convergence behavior of Algorithm 4.1 with respect to decreasing Lavrentiev parameter.

| $\lambda$ | $\left\\|u_{h}-u_{\text {opt }}\right\\|_{L^{2}}$ | $\left\\|y_{h}-y_{\text {opt }}\right\\|_{L^{2}}$ | $\left\\|\mu_{h}-\mu_{\text {opt }}\right\\|_{L^{2}}$ | $\left\\|p_{h}-p_{\text {opt }}\right\\|_{L^{2}}$ |
| :---: | :---: | :---: | :---: | :---: |
| $10^{-2.0}$ | $3.1880078 \mathrm{e}-01$ | $3.9198985 \mathrm{e}-01$ | $6.2201011 \mathrm{e}-01$ | $3.5549336 \mathrm{e}-03$ |
| $10^{-3.0}$ | $3.2990915 \mathrm{e}-02$ | $2.0899831 \mathrm{e}-02$ | $2.4068944 \mathrm{e}-01$ | $1.1464553 \mathrm{e}-03$ |
| $10^{-4.0}$ | $4.4104922 \mathrm{e}-03$ | $4.5454556 \mathrm{e}-03$ | $1.2453488 \mathrm{e}-01$ | $7.7483476 \mathrm{e}-04$ |
| $10^{-5.0}$ | $2.5020193 \mathrm{e}-03$ | $2.6584714 \mathrm{e}-03$ | $6.3426567 \mathrm{e}-02$ | $3.5686997 \mathrm{e}-04$ |
| $10^{-6.0}$ | $1.8128938 \mathrm{e}-03$ | $1.0053254 \mathrm{e}-03$ | $4.2474621 \mathrm{e}-02$ | $2.6608656 \mathrm{e}-04$ |
| $10^{-7.0}$ | $1.1850217 \mathrm{e}-03$ | $3.5760613 \mathrm{e}-04$ | $4.0462811 \mathrm{e}-02$ | $4.1388376 \mathrm{e}-04$ |
| $10^{-8.0}$ | $8.1408314 \mathrm{e}-04$ | $1.3624889 \mathrm{e}-04$ | $9.2715530 \mathrm{e}-02$ | $1.0450310 \mathrm{e}-03$ |
| $10^{-9.0}$ | $7.3457286 \mathrm{e}-04$ | $6.5919685 \mathrm{e}-05$ | $2.6964435 \mathrm{e}-01$ | $3.0716652 \mathrm{e}-03$ |
| $10^{-10.0}$ | $7.2968216 \mathrm{e}-04$ | $4.5600222 \mathrm{e}-05$ | $7.6555488 \mathrm{e}-01$ | $8.5845442 \mathrm{e}-03$ |

Theorem 3.4. At the same time, it indicates the applicability of our technique when dealing with state-constrained optimal boundary control problems.

If the selection of the parameter $\lambda$ is too small, the approximation to the Lagrange multiplier turns out to be rather poor. In the fourth column of Table 5.1, we observe that the distance to the Lagrange multiplier is increasing with respect to decreasing $\lambda \leq 10^{-7}$, see Figure 5.2. This effect occurred most likely due to the noticeably-increasing ill-conditioning, which is monitored in the experiment for decreasing $\lambda<10^{-8}$.

Therefore, we suggest to choose a moderate Lavrentiev parameter, $\lambda \approx 10^{-7}$. For this selection, we obtained the best approximation of the desired Lagrange multiplier, see Figure 5.1.
5.3. Comparison with a commercial optimization code. As noted earlier, it is one of our goals to compare our regularization technique based on "optimize-then-discretize" (OTD) with a commercial code applied to the discretized problem ( "discretize-then-optimize", DTO). To this aim, we selected the code $Q U A D P R O G$ from the MATLAB optimization toolbox, since this is often applied by users of MATLAB.

We start by defining briefly the discretized version of $(P)$ :
$\left(P_{h}\right)$

$$
\left\{\begin{array}{l}
\text { minimize } \frac{1}{2}\left(y_{h}-y_{d_{h}}\right)^{T} M_{h}\left(y_{h}-y_{d_{h}}\right)+\frac{\alpha}{2}\left(u_{h}-u_{d_{h}}\right)^{T} \tilde{M}_{h}\left(u_{h}-u_{d_{h}}\right) \\
\text { subject to } A_{h} y_{h}=M_{h} e_{h}+B_{h} u_{h} \\
y_{h} \leq \psi_{h} \\
\left(u_{h}, y_{h}\right) \in \mathbb{R}^{M_{h}} \times \mathbb{R}^{N_{h}}
\end{array}\right.
$$

where the mass matrix $\tilde{M}_{h}$ is given by $\tilde{M}_{h_{i j}}=\left(\psi_{h}^{i}, \psi_{h}^{j}\right)_{L^{2}(\Gamma)}$ and the vectors $\psi_{h}, y_{d_{h}} \in \mathbb{R}^{N_{h}}$ and $u_{d_{h}} \in \mathbb{R}^{M_{h}}$ stand for the discretization of the upper bound function $\psi$, the desired state $y_{d}$ and the desired control $u_{d}$, respectively. $Q U A D P R O G$ was used to solve this linear quadratic problem $\left(P_{h}\right)$.

Table 5.3 shows the results when solving the previous test example. Considering the first and second columns of Table 5.3, the numerical approximations of DTO to the optimal state and optimal control with respect to decreasing mesh size are quite satisfactory just as those of OTD. We point out here that, in contrast to the previous numerical test for OTD, the computation based on DTO for the test example with mesh size $h=1 / 128$ failed due to exceeding the memory of the PC.

Fig. 5.1. Computed solution based on Algorithm 4.1 for $h=1 / 128$ and $\lambda=10^{-7}$ : Optimal state (upper left), optimal control (upper right), adjoint state (lower left) and Lagrange multiplier (lower right).

TABLE 5.2
The $L^{2}$-error to the optimal values

| $N$ | $\left\\|u_{h}-u_{o p t}\right\\|_{L^{2}}$ | $\left\\|y_{h}-y_{o p t}\right\\|_{L^{2}}$ |
| :---: | :---: | :---: |
| 8 | $2.453490834802477 \mathrm{e}-01$ | $4.004580211912924 \mathrm{e}-03$ |
| 16 | $7.119225785142243 \mathrm{e}-02$ | $1.533151564674966 \mathrm{e}-03$ |
| 32 | $1.901332987811985 \mathrm{e}-02$ | $4.525371992082597 \mathrm{e}-04$ |
| 64 | $4.927627144442628 \mathrm{e}-03$ | $1.156125655672012 \mathrm{e}-04$ |

Next, we aim at comparing the efficiency of our methods OTD with that of $Q U A D P R O G$ based on the discretize-then-optimize concept. For this purpose, we test again OTD at $\lambda=10^{-7}$ as well as DTO with various mesh sizes and report on their required CPU-time for each grid. Observing Table 5.3 , we detect that OTD was more efficient. On the finest grid of the numerical tests $(h=1 / 64)$, DTO required about $6.46 \mathrm{e}+03$ seconds to converge, whereas OTD was at least 50 times faster. This demonstrates the computational efficiency of our regularization strategy.
5.4. Test example 2. We consider now an example without given analytical solution.

Fig. 5.2. Spurious oscillation of the Lagrange multipliers with respect to decreasing $\lambda$.

Table 5.3
CPU times for various mesh sizes.

| $h$ | $\frac{1}{8}$ | $\frac{1}{16}$ | $\frac{1}{32}$ | $\frac{1}{64}$ |
| :---: | :---: | :---: | :---: | :---: |
| DTO | $4.00 \mathrm{e}-02$ | $1.18 \mathrm{e}+00$ | $1.23 \mathrm{e}+02$ | $6.46 \mathrm{e}+03$ |
| OTD $\lambda=10^{-7}$ | $2.00 \mathrm{e}-01$ | $1.85 \mathrm{e}+00$ | $1.33 \mathrm{e}+01$ | $9.89 \mathrm{e}+01$ |

Example 5.1. We choose:

$$
\begin{aligned}
& y_{d}=\left(x_{1}-0.5\right)^{2}+\left(x_{2}-0.5\right)^{2}-3 \\
& u_{d}=0 \\
& e=0 \\
& \alpha=10^{-2}
\end{aligned}
$$

and the state constraint is given by

$$
y(x) \leq-2.723 \quad \forall x \in \bar{\Omega}
$$

Figure 5.3 displays the computed solution for $\lambda=10^{-9}$. Again, the optimal control is identical on all edges of $\Omega$ and hence is only plotted on the lower boundary.

Fig. 5.3. Computed solution based on Algorithm 4.1 for $h=1 / 128$ and $\lambda=10^{-9}$

TABLE 5.4
Number of iterations for several Lavrentiev-parameter choices with fixed mesh size $h=1 / 128$.

| $\lambda$ | $10^{-3}$ | $10^{-4}$ | $10^{-5}$ | $10^{-6}$ | $10^{-7}$ | $10^{-8}$ | $10^{-9}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| \# It. | 7 | 9 | 10 | 10 | 16 | 29 | 34 |

In Table 5.4, we provide the required iteration numbers when solving the problem using our algorithm with fixed initial data $p_{0}=y_{0}=0$. Just as before, we detect that the problem is harder to be solved if the choice of $\lambda$ is too small. Based on our numerical observation, this effect occurred mostly due to the following two reasons: the linear equations solved by the algorithm are ill-conditioned. This behavior is monitored particularly if the selection of the Lavrentiev-parameter is too small. Furthermore, the measure structure of the Lagrange multipliers associated with the upper bound complicates considerably the numerical computation. As $\lambda \rightarrow 0$, the computed Lagrange multiplier approaches a Dirac measure concentrated in a single point, cf. Figure 5.4.
5.5. Nested iteration. In this last part of the paper, we briefly introduce an interpolation technique with the objective to gain a faster method. Our experience in handling the regularized distributed control problem, cf. [9], indicates that a simple nested iteration scheme, based on the coarse-to-fine grid strategy, may improve the efficiency of the algorithm. In fact, this method is also reliable for boundary control problems and can significantly accelerate the convergence. First, let us shortly explain the nested iteration method: We set a sequence of grids $\Omega_{k}$ with mesh size $h_{k}=2 h_{k-1}$ and start by solving the problem over the coarsest grid $\Omega_{0}$. Subsequently, we interpolate

FIG. 5.4. Computed Lagrange multipliers associated with the upper bound: $\lambda=10^{-4}, \lambda=10^{-6}, \lambda=10^{-7}$ and $\lambda=10^{-9}$ (from left to right).
the result to the next finer grid $\Omega_{1}$ by the nine prolongation method, see [7], and utilize it as initial data for the algorithm with finer mesh size $h_{1}$. We repeat this process until the desired mesh size is reached. In table 5.5 , we present a comparison of the results for the first test example with $\lambda=10^{-5}$ based on Algorithm 4.1 (OTD) and those based on the nested iteration strategy.

In the second row of Table 5.5, we report on the iteration numbers as well as the CPU-time required by OTD with fixed initial data $y_{0}=q_{0}=0$. The row with title Interpolation displays the performance of the algorithm combined with the nested iteration scheme (NIS) including its required CPU-time for each grid. Clearly, we find a significant speed-up of the algorithm under this coarse to fine grid sweep: By comparing the accumulated CPU-time of NIS with that of OTD at $h=1 / 128$ with fixed initial data, we infer that NIS is more than three times faster than OTD with fixed initial data.

Certainly, under this nested iteration scheme, a similar improvement of the performance might have obtained for DTO based on $Q U A D P R O G$. We did not investigate this alternative.

Our second noticeable observation is the mesh-independence behavior of OTD. Regardless of the mesh size of the discretization, the iteration numbers of OTD with fixed initial data remain constant equal to six. This might correspond to the mesh independence principle for a similar regularization technique applied to distributed optimal control problems [9]. Notice that after regularization, our problem is equivalent to a control-constrained one (cf. our remarks before Theorem 3.3).

Table 5.5
Mesh independence behavior of OTD and speed-up under a coarse-to-fine mesh sweep.

| $h$ | $\frac{1}{32}$ | $\frac{1}{64}$ | $\frac{1}{128}$ |
| :---: | :---: | :---: | :---: |
| Fixed mesh (OTD) | 6 | 6 | 6 |
| CPU-time | $6.420 \mathrm{e}+00$ | $4.600 \mathrm{e}+01$ | $4.208 \mathrm{e}+02$ |
| Interpolation (NIS) | 6 | 1 | 1 |
| CPU-time | $6.420 \mathrm{e}+00$ | $1.221 \mathrm{e}+01$ | $1.278 \mathrm{e}+02$ |

## REFERENCES

[1] J.-J. Alibert and J.-P. Raymond. Boundary control of semilinear elliptic equations with discontinuous leading coefficients and unbounded controls. Numer. Funct. Anal. and Optimization, 3\&4:235-250, 1997.
[2] M. Bergounioux, K. Ito, and K. Kunisch. Primal-dual strategy for constrained optimal control problems. SIAM J. Control and Optimization, 37:1176-1194, 1999.
[3] M. Bergounioux and K. Kunisch. Primal-dual active set strategy for state-constrained optimal control problems. Computational Optimization and Applications, 22:193-224, 2002.
[4] E. Casas. Control of an elliptic problem with pointwise state constraints. SIAM J. Control and Optimization, 4:1309-1322, 1986.
[5] E. Casas. Boundary control of semilinear elliptic equations with pointwise state constraints. SIAM J. Control and Optimization, 31:993-1006, 1993.
[6] X. Chen, Z. Nashed, and L. Qi. Smoothing methods and semismooth methods for nondifferentiable operator equations. SIAM J. Numer. Anal., 38(4):1200-1216 (electronic), 2000.
[7] W. Hackbusch. Multigrid methods and applications, volume 4 of Springer Series in Computational Mathematics. Springer-Verlag, Berlin, 1985.
[8] M. Hintermüller, K. Ito, and K. Kunisch. The primal-dual active set strategy as a semismooth Newton method. SIAM J. Optim., 13:865-888, 2003.
[9] M. Hintermüller, F. Tröltzsch, and I. Yousept. Mesh-independence of semismooth Newton methods for Lavrentiev-regularized state constrained nonlinear optimal control problems. 2006.
[10] K. Ito and K. Kunisch. Augmented Lagrangian methods for nonsmooth, convex optimization in Hilbert spaces. Nonlinear Analysis TMA, 41:591-616, 2000.
[11] K. Ito and K. Kunisch. Semi-smooth Newton methods for state-constrained optimal control problems. Systems and Control Letters, 50:221-228, 2003.
[12] Lavrentiev, M. M. Some Improperly Posed Problems of Mathematical Physics. Springer, New York, 1967.
[13] H. Maurer and H. D. Mittelmann. Optimization techniques for solving elliptic control problems with control and state constraints. I: Boundary control. J. of Computational and Applied Mathematics, 16:29-55, 2000.
[14] C. Meyer, A. Rösch, and F. Tröltzsch. Optimal control of PDEs with regularized pointwise state constraints. Computational Optimization and Applications, 33:209-228, 2006.
[15] C. Meyer and F. Tröltzsch. On an elliptic optimal control problem with pointwise mixed control-state constraints. In A. Seeger, editor, Recent Advances in Optimization. Proceedings of the 12 th French-GermanSpanish Conference on Optimization held in Avignon, September 20-24, 2004, Lectures Notes in Economics and Mathematical Systems, Vol. 563, pages 187- 204. Springer-Verlag, 2006.


[^0]:    *Institut für Mathematik, Technische Universität Berlin, D-10623 Berlin, Str. des 17. Juni 136, Germany (troeltz@math.tu-berlin.de, yousept@math.tu-berlin.de). The authors acknowledge support through DFG Research Center "Mathematics for Key Technologies" (FZT 86) in Berlin.

