



DFG-Research Center MATHEON
Mathematics for Key Technologies

Towards doping optimization of semiconductor lasers

Dirk Peschka

Nella Rotundo

Marita Thomas

Preprint

MATHEON **preprint**

<http://opus4.kobv.de/opus4-mattheon>

Preprint

February 2016

Towards doping optimization of semiconductor lasers

D. Peschka¹, N. Rotundo², and M. Thomas³

¹*peschka@wias-berlin.de*

²*rotundo@wias-berlin.de*

³*thomas@wias-berlin.de*

^{1,2,3}*Weierstrass Institute for Applied Analysis and Stochastics Mohrenstr. 39, 10117 Berlin Germany*

Abstract

We discuss analytical and numerical methods for the optimization of optoelectronic devices by performing optimal control of the PDE governing the carrier transport with respect to the doping profile. First, we provide a cost functional that is a sum of a regularization and a contribution, which is motivated by the *modal net gain* that appears in optoelectronic models of bulk or quantum-well lasers. Then, we state a numerical discretization, for which we study optimized solutions for different regularizations and for vanishing weights.

1 Introduction

Studies of semiconductor device optimization via optimal control methods have been the subject of a number of previous studies, cf. e.g. [6–8]. In recent years there has been an increase in research on optoelectronic devices, e.g. aimed at on-chip integration of lasers in order to increase communication bandwidths for computing or telecommunication applications. For example, it has been demonstrated recently that germanium can be used as an optically active medium, however, advanced engineering techniques such as high doping or application of large strain, are necessary to improve optical properties of laser cavities [1].

In the following, we present our first steps towards the optimal control of an optoelectronic device. In Sec. 2 we introduce a model that has been used to describe optical modes, charge transport and spontaneous emission in the cross-section of an edge-emitting laser. We first simplify this model to the extent we believe is reasonable for a laser below lasing threshold. For this model we discuss well-posedness in Sec. 3 relying on well-known existence results cf. e.g. [4,13,17,18]. In Sec. 4 we set up an optimal control problem with the goal to improve the optical properties of the laser. Finally, in Sec. 5 we discuss the feasibility of the optimization concept at the hand of 1D-examples.

2 Mathematical model

The state of a semiconductor in a bounded domain $\Omega \subset \mathbb{R}^d$, with $d = 1, 2, 3$ is described by the electrostatic potential ψ and by the carrier densities n, p for the electrons in the conduction band and holes in the valence band, respectively. Steady-state solutions (ψ, n, p) solve the van Roosbroeck system [4], which after non-dimensionalization using $[x] = l$, $[\psi] = U_T$, $[n, p, C] = N_0$ reads

$$-\Delta\psi = q(C + p - n), \quad \nabla \cdot j_n = R, \quad \nabla \cdot j_p = -R. \quad (1a)$$

The current densities for electrons $j_n = -\mu_n n \nabla \varphi_n$ and holes $j_p = -\mu_p p \nabla \varphi_p$ are proportional to the gradients of the quasi-Fermi potentials φ_n and φ_p , which are related to the carrier densities through the equation of state

$$n = n_i \exp(\psi - \varphi_n), \quad \text{and} \quad p = n_i \exp(\varphi_p - \psi). \quad (1b)$$

With C we denote the given doping concentration. In this model we rescaled the dimensional mobilities by μ_0 , recombination rates by $R_0 = U_T \mu_0 N_0 / l^2$, intrinsic densities by N_0 , and defined

$q = l^2 q_e N_0 / (U_T \epsilon)$ and $U_T = k_B T / q_e$. At Ohmic contacts Γ_D these equations are supplemented with the following special type of Dirichlet boundary conditions

$$\psi = \psi_{\text{bi}} + \psi_{\text{ext}}, \quad \varphi_n = \psi_{\text{ext}}, \quad \varphi_p = \psi_{\text{ext}}, \quad \text{on } \Gamma_D, \quad (1c)$$

using the built-in voltage ψ_{bi} and the given external voltage ψ_{ext} , and on the insulating part with

$$\nabla \psi \cdot \nu = j_n \cdot \nu = j_p \cdot \nu = 0, \quad \text{on } \Gamma_N. \quad (1d)$$

Rewriting the fluxes in (1) in terms of the densities n, p , one arrives at the well-known form

$$j_n = -\mu_n (n \nabla \psi - \nabla n), \quad j_p = -\mu_p (p \nabla \psi + \nabla p), \quad (2)$$

and at Ohmic contacts the boundary conditions become

$$\psi = \psi_{\text{bi}} + \psi_{\text{ext}} =: \psi_D, \quad n = n_i e^{\psi_{\text{bi}}} =: n_D, \quad p = n_i e^{-\psi_{\text{bi}}} =: p_D, \quad \text{on } \Gamma_D, \quad (3)$$

so that we have $np = n_i^2$ at Γ_D .

In addition to the electronic transport model (1), an optoelectronic model needs to take into account radiative recombination processes and should, in particular, capture the number of coherent photons S generated by stimulated emission. Analogously to [9] we introduce the optical modes Ψ as solutions of a Helmholtz eigenvalue problem, which prototypically for TE modes reads

$$\left[\nabla^2 + (n_r + \frac{i}{2}(g - \ell - \ell_1))^2 \right] \Psi = \beta^2 \Psi,$$

with n_r the refractive index of the material. The imaginary part of the eigenvalue enters the stationary balance of stimulated emission and spontaneous emission r_{sp} as $0 = (2\text{Im}\beta - \ell_2)S + r_{\text{sp}}$, where ℓ_1, ℓ_2 constitute additional loss mechanisms not considered here. With g and ℓ we denote the *gain* and, as a representative loss mechanism, the *free carrier absorption*, which encodes the rate at which photons S are created and annihilated and which both depend on carrier densities, i.e. $g = g(n, p)$ and $\ell = \ell(n, p)$. Motivated by one of our previous studies [2] we use

$$g(n, p) := g_0 (np)^\delta - g_1 (np)^{\delta-1}, \quad (4a)$$

$$\ell(n, p) := f_n n + f_p p, \quad (4b)$$

where $g_0, g_1, f_n, f_p > 0, \delta \in (0, 1/2)$ are material dependent parameters (possibly depending on wavelength, material quality, temperature, mechanical strain etc.) and which in the following are taken to be constant in space. By a perturbation argument one can show that the imaginary part of the eigenvalue can be approximated

$$\text{Im}\beta \simeq \int_{\Omega} (g(n, p) - \ell(n, p)) |\Psi|^2 dx, \quad (5)$$

where, to leading order, the form of Ψ does not depend on the charge distribution.

The recombination-generation term $R = R_{\text{norad}} + R_{\text{rad}}$ in (1) takes into account different non-radiative and radiative recombination processes. The non-radiative Auger and Shockley-Read-Hall recombination terms are of the form $R_{\text{norad}} = \tilde{R}(n, p)(np - n_i^2)$. In addition we have radiative recombinations $R_{\text{rad}} = \tilde{R}_{\text{spont}}(np - n_i^2) + R_{\text{stim}}(n, p)$. In addition to spontaneous recombination it contains the main ingredient for a laser, the stimulated radiative recombination

$$R_{\text{stim}}(n, p) \simeq g(n, p) |\Psi|^2 S, \quad (6)$$

which has the stimulated emission term $\text{Im}\beta$ in the algebraic equation for S as a counter-part. For the purpose of this paper we consider the optimization of a laser below threshold $0 < S \ll 1$. From a practical point of view this implies that we neglect the influence of the optics on the electronic transport and set $S = 0$, whereas it is taken into account in the optimization by using $\text{Im}\beta$ in (5) as one part of the cost functional. Such a strategy is aimed at lowering the threshold current of a semiconductor laser.

3 The existence of solutions

In the subsequent sections we assume that

- the domain Ω is bounded in \mathbb{R}^d , $d \in \{1, 2, 3\}$, with Lipschitz boundary such that $\Gamma_D \neq \emptyset$ and $\Gamma_N = \partial\Omega \setminus \Gamma_D$,
- $R(n, p) = \tilde{R}(n, p)(np - n_i^2)$ with $\tilde{R} : [0, \infty)^2 \rightarrow [0, \infty)$ continuously.

In order to account for general Dirichlet data (ψ_D, n_D, p_D) , we shall understand from now on $\xi_D = (\psi_D, n_D, p_D)$ as the extension of the data prescribed on Γ_D into the domain Ω . The special case of Ohmic contacts (1c) is included in the existence Theorem 1. Moreover, we introduce $H_D^1(\Omega) := \{u \in H^1(\Omega), u = 0 \text{ on } \Gamma_D\}$. Then we set

$$\mathbf{H} := [H^1(\Omega)]^3, \quad \mathbf{X} := [H_D^1(\Omega) \cap L^\infty(\Omega)]^3. \quad (8)$$

Since we want to allow for some freedom in the choice of the cost functional later on in Sections 4 & 5, we keep the space \mathbf{C} as general as possible, this means that \mathbf{C} can be chosen accordingly

$$\mathbf{C} := H^1(\Omega) \quad \text{or} \quad \mathbf{C} := L^p(\Omega) \text{ s.t. } p > d/2. \quad (9)$$

In order to specify our notion of solution we shall interpret the system (1a,2,3) as an operator

$$\rho_{\text{vR}}(\xi_D; \cdot, \cdot) : \mathbf{X} \times \mathbf{C} \rightarrow [H^1(\Omega)^*]^3, \quad \rho_{\text{vR}}(\xi_D; \xi, C) := \begin{pmatrix} -\Delta(\psi + \psi_D) - q(C + p + p_D - n - n_D) \\ \nabla \cdot j_n - R \\ \nabla \cdot j_p + R \end{pmatrix},$$

$$\text{where } j_n := \nabla(n + n_D) + (n + n_D)\nabla(\psi + \psi_D), \quad j_p := -\nabla(p + p_D) + (p + p_D)\nabla(\psi + \psi_D), \quad (10)$$

with $\xi = (\psi, n, p) \in \mathbf{X}$ and $H^1(\Omega)^*$ the dual of $H^1(\Omega)$. In this way, a suitable notion for solutions of the boundary value problem (1) can be stated as follows.

Definition 1 (Solution for the boundary value problem (1)) For the given data $C \in \mathbf{C}$ and $\xi_D = (\psi_D, n_D, p_D) \in [H^1(\Omega) \cap L^\infty(\Omega)]^3$ a triple $\xi = (\psi, n, p) \in \mathbf{X}$ is a solution to the boundary value problem of the van Roosbroeck system (1) if

$$\rho_{\text{vR}}(\xi_D; \psi, n, p, C) = \mathbf{0} \quad \text{in } [H^1(\Omega)^*]^3. \quad (11)$$

Arguing along the lines of [4,18,13,7] the following existence result can be obtained.

Theorem 1 (Existence of solutions) Let (7) be satisfied and $\rho_{\text{vR}}(\xi_D, \cdot, \cdot) : \mathbf{X} \times \mathbf{C} \rightarrow [H^1(\Omega)^*]^3$ as in (8) & (10). Then for all given data $C \in \mathbf{C}$ and $(\psi_D, n_D, p_D) \in [H^1(\Omega) \cap L^\infty(\Omega)]^3$ such that

$$\|\psi_D\|_{L^\infty(\Omega)} \leq K, \quad \text{and} \quad \frac{1}{K} \leq n_D, p_D \leq K \quad \text{a.e. in } \Omega \quad (12)$$

for some $K \geq 1$, there exists a triple $(\psi, n, p) \in \mathbf{X}$ such that (11) is satisfied and

$$\|(\psi, n, p)\|_{\mathbf{X}} \leq L \quad \text{as well as} \quad \frac{1}{L} \leq n, p \leq L \quad \text{a.e. in } \Omega \quad (13)$$

with a constant $L = L(\Omega, K, \|C\|_{\mathbf{C}}) \geq 1$.

Remark 1 (Comments on the existence proof) *For the shortness of the presentation we shall not give a full proof of the existence Theorem 1, for all the details we refer to [4, Sec. 3]. The main idea is to apply a fixed point argument based on Schauder's fixed point theorem using a decoupled iteration scheme. It is important to note that the bounds on the given data (12) and the choice of \mathbf{C} according to (9) allow us to find the bounds (13) on the solutions (ψ, n, p) . In particular, it is important to mention Stampaccia's method [14] as one important step in the existence proof. Thanks to this, it is possible to prove that a solution for the electrostatic potential is bounded in $L^\infty(\Omega)$ by a constant which continuously depends on the $L^p(\Omega)$ -norm of the doping function, on the bound corresponding to K for the Dirichlet datum, and on the space dimension, cf. also [11,12]. Given that the space \mathbf{C} is chosen as in (9), this implies, that ψ is also bounded in $L^\infty(\Omega)$ by a bound featuring a continuous dependence on the \mathbf{C} -norm of the doping function. This information will play an important role in the optimization lateron.*

Introducing, in accordance with the bounds (13), the topology τ_{vR}

$$(\psi_k, n_k, p_k, C_k) \xrightarrow{\tau_{\text{vR}}} (\psi, n, p, C) \Leftrightarrow \begin{cases} (\psi_k, n_k, p_k) \rightharpoonup (\psi, n, p) \text{ in } [H^1(\Omega)]^3, \\ (\psi_k, n_k, p_k) \overset{*}{\rightharpoonup} (\psi, n, p) \text{ in } [L^\infty(\Omega)]^3, \\ (\psi_k, n_k, p_k) \rightarrow (\psi, n, p) \text{ in } [L^q(\Omega)]^3 \text{ for all } q \in (1, \infty), \\ C_k \rightarrow C \text{ in } \mathbf{C} \end{cases} \quad (14)$$

and the set

$$\mathbf{S}_{\text{vR}} := \{(\xi, C) \in \mathbf{X} \times \mathbf{C}, (\xi, C) \text{ satisfy (11) \& (13) with } \xi_{\text{D}} \text{ as in (12)}\}, \quad (15)$$

it can be shown, thanks to (13), that \mathbf{S}_{vR} is compact in $\mathbf{H} \times \mathbf{C}$ with respect to the topology τ_{vR} . This will be the crucial ingredient to verify the existence of a solution for the optimization problem treated in Sec. 4.

Proposition 1 (Compactness of \mathbf{S}_{vR}) *The set \mathbf{S}_{vR} defined in (15) is compact in $\mathbf{H} \times \mathbf{C}$ with respect to the topology τ_{vR} .*

Proof: Consider a sequence $(\hat{\psi}_k, \hat{n}_k, \hat{p}_k, \hat{C}_k)_k \subset \mathbf{S}_{\text{vR}}$ such that $\|(\hat{\psi}_k, \hat{n}_k, \hat{p}_k, \hat{C}_k)\|_{\mathbf{H} \times \mathbf{C}} \leq M$ for a constant $M < \infty$. Thus, $\|\hat{C}_k\|_{\mathbf{C}} \leq M$ and hence, by (13), also $\|(\hat{\psi}_k, \hat{n}_k, \hat{p}_k)\|_{[L^\infty(\Omega)]^3} \leq L(M)$ and $1/L(M) \leq \hat{n}_k, \hat{p}_k \leq L(M)$ a.e. in Ω . By the reflexivity of $[H^1(\Omega)]^3$ we find a (not relabelled) subsequence $(\hat{\psi}_k, \hat{n}_k, \hat{p}_k)_k$ and a triple $(\hat{\psi}, \hat{n}, \hat{p}) \in [H^1(\Omega)]^3$ such that $(\hat{\psi}_k, \hat{n}_k, \hat{p}_k) \rightharpoonup (\hat{\psi}, \hat{n}, \hat{p})$ in $[H^1(\Omega)]^3$. By the compact embedding $[H^1(\Omega)]^3 \Subset [L^2(\Omega)]^3$ we have that $(\hat{\psi}_k, \hat{n}_k, \hat{p}_k) \rightarrow (\hat{\psi}, \hat{n}, \hat{p})$ in $[L^2(\Omega)]^3$, which can be upgraded to strong convergence in $[L^q(\Omega)]^3$ for any $q \in (1, \infty)$ thanks to the uniform L^∞ -bound. Indeed, this bound also implies the weak-star convergence in $L^\infty(\Omega)$ of a further subsequence. Upon extrating a further subsequence that converges pointwise a.e. in Ω we deduce that $1/L(M) \leq \hat{n}, \hat{p} \leq L(M)$ a.e. in Ω . Finally, using Banach-Alaoglu's theorem, the uniform bound of $(\hat{C}_k)_k$ in \mathbf{C} , as a dual of a separable Banach space, allows us to conclude the convergence of a subsequence to a limit $\hat{C} \in \mathbf{C}$. In conclusion, we have obtained that also the limit $(\hat{\psi}, \hat{n}, \hat{p}, \hat{C})$ satisfies (13).

It remains to show that the limit $(\hat{\psi}, \hat{n}, \hat{p}, \hat{C})$ complies with (11). Thanks to the first and the fourth convergence property of (14) we find that $\hat{\psi}$ is a weak solution of the Poisson equation in (1), i.e. the first line of system (11), with right-hand side $q(\hat{C} - \hat{n} - n_{\text{D}} + \hat{p} + p_{\text{D}})$. In order to show that $(\hat{\psi}, \hat{n}, \hat{p})$ also solve the corresponding continuity equations note that with $(\hat{n}_k, \hat{p}_k)_k$ satisfying (13) we have $|R| \leq \max_{(n,p) \in [1/L, L]} \tilde{R}(n + n_{\text{D}}, p + p_{\text{D}}) ((L + K)^2 - n_i^2)$ pointwise a.e. in Ω , uniformly for all

$k \in \mathbb{N}$. Owing to the strong L^q -convergence ensured in the third line of (14), we may extract a not re-labelled subsequence such that $(\hat{n}_k, \hat{p}_k) \rightarrow (\hat{n}, \hat{p})$ pointwise a.e. and since all the reaction terms are continuous according to (7), we find that $R(\hat{n}_k + n_D, \hat{p}_k + p_D) \rightarrow R(\hat{n} + n_D, \hat{p} + p_D)$ along this subsequence. Thus, thanks to the pointwise uniform bound, we find with the aid of the dominated convergence theorem that $\int_{\Omega} R(\hat{n}_k + n_D, \hat{p}_k + p_D) v \, dx \rightarrow \int_{\Omega} R(\hat{n} + n_D, \hat{p} + p_D) v \, dx$ for any test function $v \in H^1(\Omega)$, hence the convergence of the right-hand sides in the second and third line of (11). Finally, we obtain the convergence of the corresponding left-hand sides by weak-strong convergence arguments using the first of (14) and that $\|(\hat{n}_k - \hat{n})v\|_{L^2(\Omega)}^2 \leq \|v\|_{L^{2^*}(\Omega)}^2 \|\hat{n}_k - \hat{n}\|_{L^{(2^*)'}(\Omega)}^2 \rightarrow 0$ thanks to the third of (14), and the corresponding argument for the hole density. Thus, $(\hat{\psi}, \hat{n}, \hat{p}, \hat{C})$ satisfies (11) and altogether we have verified that $(\hat{\psi}, \hat{n}, \hat{p}, \hat{C}) \in \mathbf{S}_{\text{VR}}$. ■

Remark 2 (Comments on the uniqueness of solutions to (11) and its effect on optimization)

Uniqueness of the solution has been proved in [13] under additional smallness assumptions on the Dirichlet data and restrictions on K in (12) which ensure that the applied voltage is sufficiently small and hence keep the carrier densities solving (11) sufficiently close to the thermal equilibrium state. Further away from equilibrium, uniqueness is in general not to be expected. This has the effect that also an optimization problem based on (11) may admit multiple solutions. Also observe that, due to the quasilinear character of the current continuity equations, it is unclear whether \mathbf{S}_{VR} from (15) forms a convex set. This spoils the uniqueness of a minimizer, or at least its verification, for an optimization problem involving functionals with equality constraint (11), even if the functional itself is strictly convex. We refer to [8, Sec. 2], where the non-uniqueness of a minimizer has been demonstrated for a particular choice of the cost functional.

4 Optimization problem

Our optimization goal is to maximize the optical output, i.e. the number of photons available in the device, with respect to the doping profile $C \in \mathbf{C}$ injected into a device, whose electrical properties are governed by the van Roosbroeck system (11). This will amount to a constrained minimization problem for a suitable cost functional $\mathcal{Q} : \mathbf{X} \times \mathbf{C} \rightarrow \mathbb{R} \cup \{\infty\}$ which shall be introduced now. Following [2], the optical output is related to the modal net-gain

$$-Q_1(n, p) = (g - \ell)|\Psi|^2, \quad (16)$$

with the representation of the optical mode $\Psi \in C_0^\infty(\Omega)$ as in (6) and with the material gain g and the optical losses ℓ from (4). In view of this, we introduce the functional

$$\tilde{\mathcal{Q}}_1 : \mathbf{X} \rightarrow \mathbb{R}, \quad \tilde{\mathcal{Q}}_1(\psi, n, p) := \int_{\Omega} Q_1(\psi + \psi_D, n + n_D, p + p_D) \, dx, \quad (17)$$

so that the maximization of the optical output will be realized by minimizing a functional that involves $\tilde{\mathcal{Q}}_1$. In particular, since the doping profile $C \in \mathbf{C}$ shall be the control parameter of the optimization problem, we shall combine $\tilde{\mathcal{Q}}_1$ with a second functional

$$\tilde{\mathcal{Q}}_2 : \mathbf{C} \rightarrow \mathbb{R} \cup \{\infty\} \text{ with domain } \text{dom } \tilde{\mathcal{Q}}_2 = \mathbf{C}_0 \subset \mathbf{C} \text{ a closed, convex subset} \quad (18a)$$

and we impose that $\tilde{\mathcal{Q}}_2 : \mathbf{C} \rightarrow \mathbb{R} \cup \{\infty\}$ is:

- weakly sequentially lower semicontinuous on \mathbf{C} , (18b)

- bounded from below and coercive: $\exists c_{\mathcal{Q}_2}, \bar{c}_{\mathcal{Q}_2} > 0, r \in (1, \infty)$, (18c)

$$\forall C \in \mathbf{C} : \tilde{\mathcal{Q}}_2(C) \geq c_{\mathcal{Q}_2} \|C\|_{\mathbf{C}}^r - \bar{c}_{\mathcal{Q}_2}.$$

Making use of (15), (17), and (18a) we finally introduce the cost functional

$$\mathcal{Q} : \mathbf{X} \times \mathbf{C} \rightarrow \mathbb{R} \cup \{\infty\}, \quad \mathcal{Q}(\psi, n, p, C) := \begin{cases} \tilde{\mathcal{Q}}_1(\psi, n, p) + \tilde{\mathcal{Q}}_2(C) & \text{if } (\psi, n, p, C) \in \mathbf{S}_{\text{vR}}, \\ \infty & \text{otherwise} \end{cases} \quad (19)$$

and in what follows we shall treat the constrained minimization problem

$$\min_{\mathbf{X} \times \mathbf{C}} \mathcal{Q}(\psi, n, p, C). \quad (20)$$

Observe that the functional $\tilde{\mathcal{Q}} := \tilde{\mathcal{Q}}_1 + \tilde{\mathcal{Q}}_2 : \mathbf{X} \times \mathbf{C} \rightarrow \mathbb{R} \cup \{\infty\}$ only carries the compactness property (18c) encoded in $\tilde{\mathcal{Q}}_2$ for the doping profiles, whereas $\tilde{\mathcal{Q}}_1$ does not give any compactness for the densities. Therefore it is important to note that the functional \mathcal{Q} gives this missing compactness by constraining the minimization to the set \mathbf{S}_{vR} , which is compact by Prop. 1. This idea will be used in order to prove the existence of a minimizer in (20). In particular, we now collect the following properties of the functional \mathcal{Q} :

Proposition 2 (Properties of \mathcal{Q}) *The functional $\mathcal{Q} : \mathbf{X} \times \mathbf{C} \rightarrow \mathbb{R} \cup \{\infty\}$, defined by the relations (15), (4a)–(19), enjoys the following properties:*

- \mathcal{Q} is bounded from below and coercive: $\exists C_{\mathcal{Q}_1} \in (-\infty, 0)$, $\exists c_{\mathcal{Q}_2}, \bar{c}_{\mathcal{Q}_2} > 0$, $r \in (1, \infty)$,
 $\forall (\psi, n, p, C) \in \mathbf{X} \times \mathbf{C} : \mathcal{Q}(\psi, n, p, C) \geq C_{\mathcal{Q}_1} + c_{\mathcal{Q}_2} \|C\|_{\mathbf{C}}^r - \bar{c}_{\mathcal{Q}_2}. \quad (21a)$
- \mathcal{Q} is lower semicontinuous wrt. convergence in τ_{vR} . (21b)

Proof: Ad boundedness from below and coercivity (21a): For this, we first check that the density $Q_1 = -(g - \ell)$ is bounded from below. Indeed, for any $n, p > 0$ observe by Young's inequality that $-g(n, p) > -g_0(np)^\delta \geq -\frac{g_0}{2}(n^{2\delta} + p^{2\delta}) =: f(n, p)$. On the interval $[0, \infty)$ the function $F(n, p) := f(n, p) + \ell(n, p)$ has a global minimum at $(n_*, p_*) := \left(\left(\frac{2\delta g_0}{f_n}\right)^{1/(1-2\delta)}, \left(\frac{2\delta g_0}{f_p}\right)^{1/(1-2\delta)}\right)$ since the Hessian $D^2F(n_*, p_*) = -(2\delta - 1)(2\delta)g_0^2 \text{diag}(n_*^{2\delta-2}, p_*^{2\delta-2})$ is positively definite because of $\delta < 1/2$. This yields that

$$-g(n, p) + \ell(n, p) > F(n, p) \geq F(n_*, p_*) \quad \text{for any } (n, p) \in [0, \infty)^2, \quad (22)$$

hence a bound from below on $\tilde{\mathcal{Q}}_1$. In combination with the bound (18c) on $\tilde{\mathcal{Q}}_2$ we thus obtain

$$\mathcal{Q}(\psi, n, p, C) \geq F(n_*, p_*) \mathcal{L}^d(\Omega) + c_{\mathcal{Q}_2} \|C\|_{\mathbf{C}}^r - \bar{c}_{\mathcal{Q}_2}.$$

Ad lower semicontinuity (21b): Consider a sequence $(\psi_k, n_k, p_k, C_k) \xrightarrow{\tau_{\text{vR}}} (\psi, n, p, C)$. In the case that $(\psi_k, n_k, p_k, C_k) \in (\mathbf{X} \times \mathbf{C}) \setminus \mathbf{S}_{\text{vR}}$ for all but a finite number of indices, then there is nothing to check because $\infty = \liminf_{k \rightarrow \infty} \mathcal{Q}(\psi_k, n_k, p_k, C_k) \geq \mathcal{Q}(\psi, n, p, C)$. Assume that $(\psi_k, n_k, p_k, C_k)_k \subset \mathbf{S}_{\text{vR}}$ for a not relabelled subsequence. By the compactness of \mathbf{S}_{vR} , ensured by Prop. 1, we thus find that the limit $(\psi, n, p, C) \in \mathbf{S}_{\text{vR}}$. Therefore we may argue that the lower semicontinuity estimate for \mathcal{Q} in this case coincides with the lower semicontinuity of $\tilde{\mathcal{Q}}_1 + \tilde{\mathcal{Q}}_2$. Then, $\liminf_{k \rightarrow \infty} \tilde{\mathcal{Q}}_2(C_k) \geq \tilde{\mathcal{Q}}_2(C)$ by (18b). We now check that $\tilde{\mathcal{Q}}_1$ is lower semicontinuous with respect to strong convergence of $(n_k, p_k)_k$ in $L^q(\Omega) \times L^q(\Omega)$ for any $q \in (1, \infty)$, which corresponds to convergence property 3 in (14). For this, we use that $-g(n_k, p_k) + \ell(n_k, p_k) > F(n_*, p_*)$, which serves as an integrable minorant according to (22), since in particular $(n_*, p_*) = \left(\left(\frac{2\delta g_0}{f_n}\right)^{1/(1-2\delta)}, \left(\frac{2\delta g_0}{f_p}\right)^{1/(1-2\delta)}\right) \in L^q(\Omega) \times L^q(\Omega)$ for any $q \in [1, \infty]$. Moreover, since strong L^q -convergence implies the convergence of the sequence in measure, Fatou's lemma ultimately yields the lower semicontinuity. ■

The existence of a minimizer can now immediately be concluded by employing the direct method of the calculus of variations, making use of the boundedness from below, the coercivity of \mathcal{Q} and its lower semicontinuity.

Theorem 2 (Existence of minimizers for the constrained minimization problem) *Let the functional $\mathcal{Q} : \mathbf{X} \times \mathbf{C} \rightarrow \mathbb{R} \cup \{\infty\}$ be defined by (4a), (15), (16)–(19). Then the constrained minimization problem (20) admits at least one solution $(\psi, n, p, C) \in (\mathbf{X} \times \mathbf{C}) \cap \mathbf{S}_{\text{vR}}$.*

Remark 3 (Different choices of \mathbf{C} and \mathcal{Q}_2) *If $\mathbf{C} = H^1(\Omega)$, possible choices for \mathcal{Q}_2 are e.g. \mathcal{Q}_2^{D} or \mathcal{Q}_2^{N} defined by*

$$\mathcal{Q}_2^{\text{D}}(C) := \begin{cases} \frac{\gamma}{2} \|\nabla(C + C_{\text{D}} - \bar{C})\|_{L^2(\Omega)}^2 & \text{if } C \in H_{\text{D}}^1(\Omega), \\ \infty & \text{otherwise,} \end{cases} \quad (23a)$$

$$\mathcal{Q}_2^{\text{N}}(C) := \begin{cases} \frac{\gamma}{2} \|C - \bar{C}\|_{H^1(\Omega)}^2 & \text{if } C \in H^1(\Omega), \\ \infty & \text{otherwise,} \end{cases} \quad (23b)$$

where $\gamma > 0$ is a weight, $\bar{C} \in H^1(\Omega)$ is a given “desired” doping profile and $C_{\text{D}} \in H^1(\Omega) \cap L^\infty(\Omega)$ the extension of a given Dirichlet datum into the domain. Note that in the first case a Poincaré inequality is available to ensure the boundedness from below (18c) in terms of the full $H^1(\Omega)$ -norm.

If $\mathbf{C} = H^1(\Omega)$, the traces to the boundary are well-defined for the doping profile C . Then also the Dirichlet boundary conditions for (ψ, n, p) can be formulated in terms of the corresponding boundary datum of C , thus allowing for Ohmic contacts (1c), cf. [4,7], where the built-in potential ψ_{bi} is determined by

$$C_{\text{D}} + n_i(e^{-\psi_{\text{bi}}} - e^{\psi_{\text{bi}}}) = 0, \quad \text{on } \Gamma_{\text{D}}. \quad (24)$$

which results in $n_{\text{D}} = (C_{\text{D}} + (C_{\text{D}}^2 + 4n_i^2)^{1/2})/2$ and $p_{\text{D}} = C_{\text{D}} - n_{\text{D}}$. Hence, in case of (23a) one has to ensure that $\max\{(\frac{1}{K} - Kn_i), (\frac{n_i}{K} - K)\} \leq C_{\text{D}} \leq \min\{K - \frac{n_i}{K}, (n_i K - \frac{1}{K})\}$, in order to match (3) with (12). In case of (23b) condition (12) is harder to meet as one has to additionally impose that $C \in L^\infty(\Gamma_{\text{D}})$. This can be achieved either by adding a further suitable penalization term to \mathcal{Q}_2^{N} or, as we will do in Sec. 4, by treating the problem in space dimension $d = 1$, where $H^1(\Omega)$ compactly embeds into $C(\bar{\Omega})$, i.e. $\|C\|_{C(\bar{\Omega})} \leq c\|C\|_{\mathbf{C}} =: Z$. For this latter case we find with $K(\|C\|_{\mathbf{C}}) := \max\{m(C), M(C)\}$, where $m(C) := (\min_{z \in [0, Z]} \{n_{\text{D}}(z), p_{\text{D}}(z)\})^{-1}$, and $M(C) = \max_{z \in [0, Z]} \{n_{\text{D}}(z), p_{\text{D}}(z), \psi_{\text{D}}(z)\}$ that $K(\|C\|_{\mathbf{C}})^{-1} \leq n_{\text{D}}(C), p_{\text{D}}(C) \leq K(\|C\|_{\mathbf{C}})$ as well as $\|\psi_{\text{D}}(C)\|_{L^\infty(\Omega)} \leq K(\|C\|_{\mathbf{C}})$. Thus, for the data $(n_{\text{D}}(C), p_{\text{D}}(C), \psi_{\text{D}}(C))$ from (3) in (11) one has $K = K(\|C\|_{\mathbf{C}})$ in (12) and hence, $L = L(\Omega, K(\|C\|_{\mathbf{C}}), \|C\|_{\mathbf{C}})$. According to Remark 1, thanks to the continuous dependence of L on its parameters, both constants are bounded on bounded sets in \mathbf{C} . Hence, it remains true that \mathbf{S}_{vR} is compact in $\mathbf{X} \times \mathbf{C}$ with respect to τ_{vR} .

In the case $\mathbf{C} = L^p(\Omega)$ a possible choice of \mathcal{Q}_2 would be

$$\mathcal{Q}_p^{\text{L}}(C) = \begin{cases} \frac{\gamma}{2} \|C - \bar{C}\|_{L^p(\Omega)}^p & \text{if } C \in L^p(\Omega), \\ \infty & \text{otherwise,} \end{cases} \quad (25)$$

where traces are generally not well-defined. In this case we adopt artificial Ohmic contacts by defining the built-in potential as

$$\bar{C} + n_i(e^{-\psi_{\text{bi}}} - e^{\psi_{\text{bi}}}) = 0, \quad \text{on } \Gamma_{\text{D}}. \quad (26)$$

This choice is not driven by any physical meaning but in order to study its impact on optimized doping profiles. In particular for $\gamma \rightarrow 0$ the artificial Ohmic contact will in general strongly violate the desired charge-neutrality.

Remark 4 (The first order optimality system) *First we notice that the mapping ρ_{vR} defined in (10), is infinitely often Fréchet differentiable with derivatives vanishing for order greater than 2, see [8, Sec. 1.2] and [7]. Moreover the following we assume that*

$$\text{dom } \tilde{Q}_2 = \mathbf{C} \text{ and } \tilde{Q}_2 : \mathbf{C} \rightarrow \mathbb{R} \text{ is twice continuously Fréchet differentiable.} \quad (27)$$

Observe that $\tilde{Q}_1 : \mathbf{X} \rightarrow \mathbb{R}$ is indeed twice continuously Fréchet differentiable. In this way, the cost functional $\tilde{Q} = \tilde{Q}_1 + \tilde{Q}_2 : \mathbf{X} \times \mathbf{C} \rightarrow \mathbb{R}$ satisfies the assumptions posed in [7, Sec.s 3,4] and hence the first order optimality system and the existence of Lagrange multipliers can be obtained following the lines of [7, Sec. 4]. Abbreviating $\xi := (\psi, n, p)$, we write the Lagrangian associated to the minimization problem (20)

$$\mathcal{L}_{\text{vR}} : \mathbf{X} \times \mathbf{C} \times \mathbf{H} \rightarrow \mathbb{R}, \quad \mathcal{L}_{\text{vR}}(\xi, C, \lambda) := \tilde{Q}(\xi, C) + \langle \rho_{\text{vR}}(\xi_{\text{D}}; \xi, C), \lambda \rangle_{\mathbf{H}}. \quad (28)$$

The associated first order optimality condition reads $D_{(\xi, C, \lambda)} \mathcal{L}_{\text{vR}}(\xi, C, \lambda) = \mathbf{0}$ in $\mathbf{X}^* \times \mathbf{C}^* \times \mathbf{H}^*$. The existence and uniqueness of a Lagrange multiplier λ is verified in [7, Thm. 4.2] under an additional smallness condition on the quotients j_n^2/n and j_p^2/p . For Ohmic contacts as in (3) and (24) the derivative $D_C \mathcal{L}_{\text{vR}}$ also acts on the dependence of ξ_{D} on C at Γ_{D} .

5 Numerical results

The numerics is based on a stationary formulation of (1) in terms of the original quasi-Fermi potentials φ_n, φ_p introduced in (1b) and the electrostatic potential ψ in one spatial dimension. Unlike the standard approach by Scharfetter-Gummel schemes [15], we use finite differences on a mesh $0 = x_1 < \dots < x_k < \dots < x_N = 1$ where the discrete electron current is expressed using $\varphi_{n;k} = \varphi_n(x_k)$, $\varphi_{p;k} = \varphi_p(x_k)$, $\psi_k = \psi(x_k)$, $\eta_{p;k} = (\varphi_{p;k} - \psi_k)$ and $\eta_{n;k} = (\psi_k - \varphi_{n;k})$ as

$$\begin{aligned} j_{n;k+1/2} &= -\mu_n n_i \exp\left(\frac{1}{2}(\eta_{n;k+1} + \eta_{n;k})\right) \left(\frac{\varphi_{n;k+1} - \varphi_{n;k}}{x_{k+1} - x_k} \right), \\ j_{p;k+1/2} &= -\mu_p n_i \exp\left(\frac{1}{2}(\eta_{p;k+1} + \eta_{p;k})\right) \left(\frac{\varphi_{p;k+1} - \varphi_{p;k}}{x_{k+1} - x_k} \right), \end{aligned}$$

and approximate the divergence by $\partial_x j_n \approx (j_{n;k+1/2} - j_{n;k-1/2})/(x_{k+1/2} - x_{k-1/2})$. The hole current $j_p = -\mu_p p \partial_x \varphi_p$ is discretized analogously. For sake of notation we use $\xi = (\psi, \varphi_n, \varphi_p)^\top$. Introducing the flux $j_{\psi;k+1/2} = -(\psi_{k+1} - \psi_k)/(x_{k+1} - x_k)$ the previous expressions can be combined into the following discrete and nonlinear residual $\tilde{\rho}_{\text{vR}} : \mathbb{R}^{3N} \times \mathbb{R}^N \rightarrow \mathbb{R}^{3N}$ as

$$\tilde{\rho}_{\text{vR}}(\xi, C) = \begin{pmatrix} \frac{j_{\psi;k+1/2} - j_{\psi;k-1/2}}{x_{k+1/2} - x_{k-1/2}} - \widehat{Q}_k \\ \frac{j_{n;k+1/2} - j_{n;k-1/2}}{x_{k+1/2} - x_{k-1/2}} - R_k \\ \frac{j_{p;k+1/2} - j_{p;k-1/2}}{x_{k+1/2} - x_{k-1/2}} + R_k \end{pmatrix}$$

for $k = 2, \dots, N-1$ and with the total charge $\widehat{Q}_k(\xi, C) = q(C_k + n_i \exp(\eta_{p;k}) - n_i \exp(\eta_{n;k}))$. This residual is the equivalent to the operator introduced in (10). It depends on the doping through the total charge and through the boundary conditions. The boundary conditions (1c) in $\tilde{\rho}_{\text{vR}}$ for $k = 1$ and $k = N$ are $\psi_k = \psi_{\text{bi}}(C_k) + \psi_{\text{ext};k}$ and $\varphi_{n/p;k} = \psi_{\text{ext};k}$. The discretized van Roosbroeck equation, which now reads $\tilde{\rho}_{\text{vR}}(\xi, C) = \mathbf{0}$, is then solved using a Newton method $\xi^{n+1} = \xi^n - (\partial_\xi \tilde{\rho}_{\text{vR}})^{-1} \tilde{\rho}_{\text{vR}}$. In thermal equilibrium $\psi_{\text{ext}} = 0$ we have $\varphi_n = \varphi_p \equiv 0$, so that since $q \gg 1$ it makes sense to choose $\xi^0 = (\psi^0, 0, 0)^\top$ with ψ^0 so that $\widehat{Q}_k = q(C_k + n_i(e^{-\psi_k^0} - e^{\psi_k^0})) = 0$. As

the magnitude of the bias ψ_{ext} is increased, the previous solution with smaller bias is used as the starting-point for the Newton method.

For the present paper the recombination in (7) is $\tilde{R} = (C_n n + C_p p) + (\tau_p(n + n_i) + \tau_n(p + n_i))^{-1}$, with the parameters $C_n = C_p = 10^{-1}$, $\tau_n = \tau_p = 10$, $n_i = 10^{-2}$. Additionally we have $q = 10^2$, $\mu_n = \mu_p = 1$. The external potential is $\varphi_{\text{ext}} = 0.4 \text{ V}/U_T \approx 15.47$ for $T = 300 \text{ K}$. These values are somewhat realistic in the sense that usually one has $0 < n_i \ll 1$ and $q \gg 1$. The optical mode is set to a Gaussian $|\Psi|^2 = \exp(-(10(x - 1/2))^2)$ and the parameters in the net-gain (4) are $f_n = 1/5$, $f_p = 1$, $\delta = 0.3$, $g_0 = e^{-1}$, $g_1 = n_i^2$.

The optimization strategy for (20) is similar to [8], where we discretize the cost functional \mathcal{Q}_1 by

$$\mathcal{Q}_1(\xi) = \sum_{k=1}^N h Q_1(n_k, p_k), \quad (29a)$$

with $Q_1 = -(g(n_k, p_k) - \ell_k)|\Psi_k|^2$ and free-carrier absorption $\ell_k = f_n n_k + f_p p_k$.

The different regularization terms $\mathcal{Q}_2^{N/D}$ from (23) for $\mathbf{C} = H^1(\Omega)$ or $\mathbf{C} = H_0^1(\Omega)$ are treated equally within our discretization by a functional \mathcal{Q}_2 . Together with the functional \mathcal{Q}_2^L for $\mathbf{C} = L^2(\Omega)$ we discretize using

$$\mathcal{Q}_2(C) = \frac{\gamma}{2} \sum_{k=1}^{N-1} h \left(\frac{C_{k+1} - C_k + \bar{C}_{k+1} - \bar{C}_k}{h} \right)^2, \quad \mathcal{Q}_2^L(C) = \frac{\gamma}{2} \sum_{k=1}^N h (C_k - \bar{C}_k)^2 \quad (29b)$$

where $h = x_{k+1} - x_k$, which for this paper we assume to be constant. Note that for \mathcal{Q}_2^N from (23b) the boundary values for C are not kept fixed. Hence, (23b) suggests to use the full $H^1(\Omega)$ -norm for regularization. However, here we study the effect of different regularizations and thereby use the $L^2(\Omega)$ -norm and $H^1(\Omega)$ -seminorm separately. As before we use $\hat{\mathcal{Q}}(\xi, C) = \mathcal{Q}_1(\xi) + \mathcal{Q}_2(C)$ and introduce a Lagrange multiplier $\lambda \in \mathbb{R}^n$, so that the first-order optimality conditions using the discrete counterpart $\mathcal{L}_{\text{vR}} : \mathbb{R}^{3N} \times \mathbb{R}^{\hat{N}} \times \mathbb{R}^{3N} \rightarrow \mathbb{R}$ of the Lagrangian in (28) is $\mathcal{L}'_{\text{vR}} = (\mathcal{L}_{\text{vR}, \xi}, \mathcal{L}_{\text{vR}, C}, \mathcal{L}_{\text{vR}, \lambda}) = 0$. For $\langle u, v \rangle$ we use the standard inner product in \mathbb{R}^{3N} . Note that, unlike [8], $\tilde{\rho}_{\text{vR}}$ is not linear in C , if we do not fix the value of C_1, C_N . In this case $\hat{N} = N$, otherwise $\hat{N} = N - 2$. Due to the nonlinearity of the boundary conditions (24) in C , this produces the slightly different reduced Hessian $\hat{\mathcal{Q}}''$ of $\hat{\mathcal{Q}}(C) = \mathcal{Q}(\xi(C), C)$ being

$$\hat{\mathcal{Q}}'' = \mathcal{Q}_{,CC} + \langle \tilde{\rho}_{\text{vR}, CC}, \lambda \rangle + \tilde{\rho}_{\text{vR}, C}^* \tilde{\rho}_{\text{vR}, \lambda}^{-*} [\mathcal{Q}_{, \xi \xi} + \langle \tilde{\rho}_{\text{vR}, \xi \xi}, \lambda \rangle] \tilde{\rho}_{\text{vR}, \xi}^{-1} \tilde{\rho}_{\text{vR}, C} \in \mathbb{R}^{\hat{N} \times \hat{N}}.$$

As in [8] we have an outer Newton iteration $C^{r+1} = C^r + \delta C^r$, which utilizes an inner CG iteration in order to solve $\hat{\mathcal{Q}}''(C^r) \delta C^r = -\hat{\mathcal{Q}}'(C^r)$. In all the considered cases, the Newton scheme converges quadratically within 7 or less iterations to our precision goal set to $< 10^{-8}$ in the max-norm of the discrete first order optimality condition. The number of iterations depends on the value of γ , which can be explained by a loss of convexity as $\gamma \rightarrow 0$. We deal with this fact using a path-following approach constructing a sequence of optimizers for $\gamma_{k+1} = \gamma_k \alpha_{\text{path}}$, here with $\alpha_{\text{path}} = 10^{-1}$, and using the previous optimizer as the initial guess for the next Newton method $k + 1$ until we reach the desired $\gamma = 10^{-6}$. The inner CG method, as expected, requires \hat{N} iterations for the residual to drop below 10^{-12} . In one spatial dimension this is not a serious restriction, whereas for higher dimensions there are some practical options to keep the number of iterations sufficiently small. One option is to modify the problem and reduce the dimension \hat{N} of the space of admissible doping profiles to be much smaller than the number of unknowns in the PDE problem, which in fact is reasonable when restricting to producible doping profiles. Another option is to define suitable alternative stopping criteria for the CG iteration, as discussed in [8].

The primary question in designing a realistic optoelectronic optimization problem is to find a cost functional, which combines the desired optimality with mathematical simplicity. Furthermore the cost functional should be feasible from an optimization point of view, for which it might be necessary to add regularizing terms to the cost. Here we study the influence of such regularizing terms Q_2 in (29b) of tracking type on the optimal solution. We study the influence of tracking the values in $L^2(\Omega)$ or the gradients in $H^1(\Omega)$ with respect to a reference doping \bar{C} . Furthermore we study the influence of the corresponding regularization parameter γ and the boundary condition on the control space \mathbf{C} .

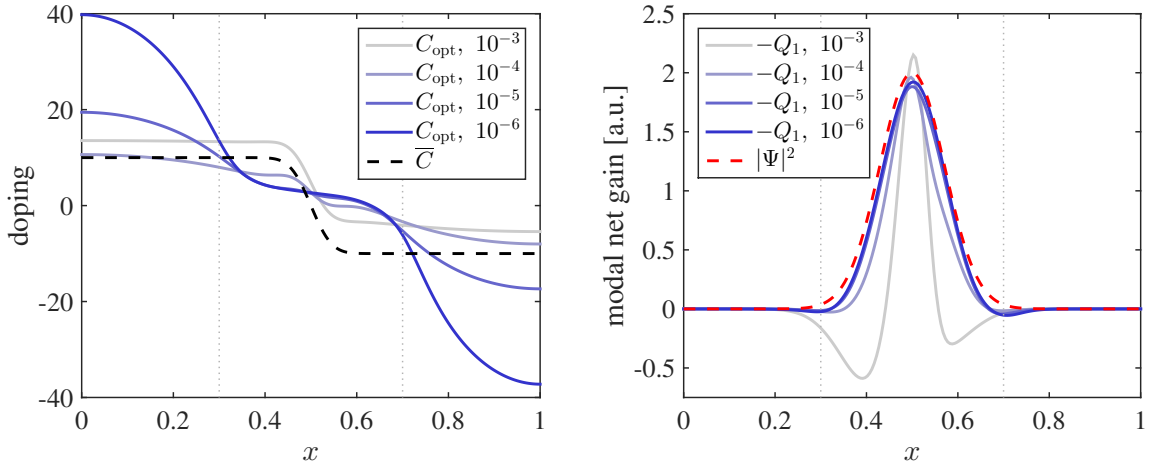


Figure 1: (Left) optimal doping densities C_{opt} for different regularization parameters γ and reference doping \bar{C} and (Right) corresponding modal net gain $-Q_1 = (g - \ell)|\Psi|^2$ compared to mode intensity $|\Psi|^2$ (arbitrary scaling). The values of the cost functional are $Q_1 = \{0.05, 0.25, 0.30, 0.31\}$ for $\gamma = \{10^{-3}, 10^{-4}, 10^{-5}, 10^{-6}\}$ and $Q_1 = -0.24$ for $C = \bar{C}$

In the left panel of Fig. 1 we show the optimal doping C_{opt} for various regularization parameters γ for gradient regularization without a fixed boundary condition, i.e. the case of $H_0^1(\Omega)$. As $\gamma \rightarrow 0$ we observe that the optimal solution stabilizes/converges where the major part of the support of $|\Psi|^2$ is located (as indicated by the dotted light gray vertical lines), whereas the doping still increases where $|\Psi|^2 \simeq 0$. For larger regularization parameters $\gamma > 10^{-5}$ the shape is still affected by the single-step reference doping, whereas for smaller regularization parameters there is very little dependence on the reference. At the same time one can see that $-Q_1 = (g - \ell)|\Psi|^2$ converges to a mostly positive solution as $\gamma \rightarrow 0$, whereas for $\gamma > 10^{-3}$ large parts of $-Q_1$ are still negative.

This can also be observed in Fig. 2. In its left panel the optimal doping is shown for $H^1(\Omega)$ regularization with $\gamma = 10^{-6}$ and two different reference dopings, e.g. a smoothed single-step and a smoothed double-step doping. Where $|\Psi|^2$ is located both optimal dopings agree and where $|\Psi|^2 \simeq 0$ the solutions are close. Even for entirely different regularization, the left panel shows that the optimal dopings are basically the same for $H^1(\Omega)$, $L^2(\Omega)$, and $H_0^1(\Omega)$ regularization. The main effect can be observed near the boundary, where $|\Psi|^2 \simeq 0$. The $L^2(\Omega)$ regularization is only a slight restriction on the doping and hence the solution is largest near the boundaries. In this case also Q_2 is slightly larger than for the $H^1(\Omega)$ regularization, which, in turn, is slightly larger than for the $H_0^1(\Omega)$ regularization. However, since the gains in all three cases are very similar, we can conclude also here that for sufficiently small γ also the influence on Q_1 is very small. We observed that when the regularization parameter is sufficiently small, then the choice of the regularization mechanism has an effect on the optimal solution but only small effect on the optimal gain.

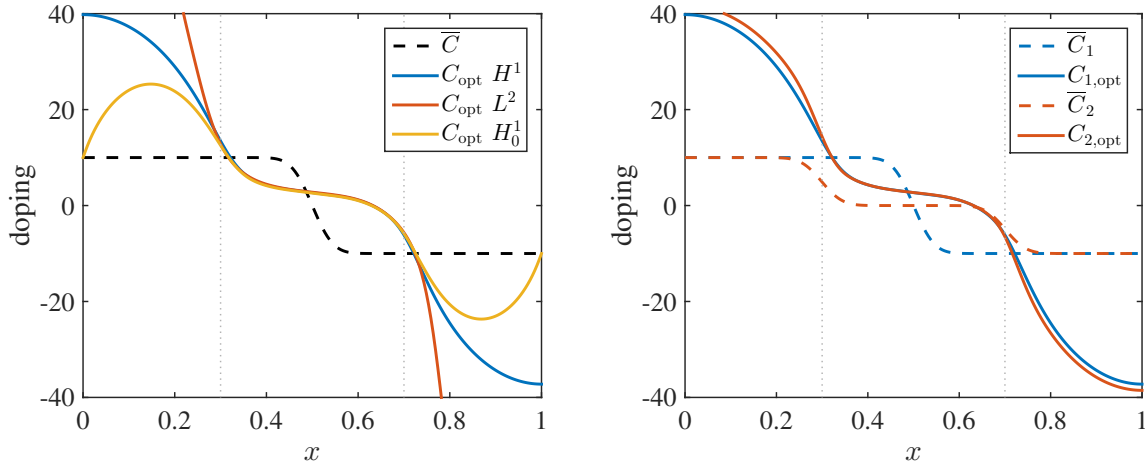


Figure 2: (Left) optimal doping densities C_{opt} for different regularization functionals and spaces $H^1(\Omega)$, $H_0^1(\Omega)$, $L^2(\Omega)$ for $\gamma = 10^{-6}$ and reference doping \bar{C} and (Right) optimal doping densities C_{opt} for $\gamma = 10^{-6}$ and two different reference dopings $\bar{C}_1 = 10 \tanh(25(\frac{1}{2} - x))$ and $\bar{C}_2 = 5 \tanh(25(0.3 - x)) + 5 \tanh(25(0.7 - x))$

We conclude that, in 1D the choice regularization of the optimization problem has little influence on the optimizer in regions, where the main part of the optical mode is located.

Acknowledgement

This work is supported by the Einstein Foundation Berlin in the framework of MATHEON, project OT1 (D.P., M.T.), and by the ERC-2010-AdG no.267802 Analysis of Multiscale Systems Driven by Functionals (N.R.). The authors kindly acknowledge the discussions with Thomas Koprucki & Annegret Glitzky (WIAS, Berlin), and Klaus Gärtner (USI, Lugano), which helped a lot to make the authors more familiar with analytical and modeling aspects of semiconductor devices.

References

- [1] B. Dutt, D.S. Sukhdeo, D. Nam, B.M. Vulovic, Z. Yuan, K.C. Saraswat, “Roadmap to an Efficient Germanium-on-Silicon Laser: Strain vs. n-Type Doping,” *IEEE Photonics*, vol. 4, (2012)
- [2] D. Peschka, M. Thomas, A. Glitzky, R. Nürnberg, K. Gärtner, M. Virgilio, S. Guha, T. Schroeder, G. Capellini, T. Koprucki, “Modeling of Edge-Emitting Lasers Based on Tensile Strained Germanium Microstrips,” *IEEE Photonics*, vol. 7,(2015).
- [3] M. Virgilio, C.L. Manganelli, G. Grosso, G. Pizzi, G. Capellini, “Radiative recombination and optical gain spectra in biaxially strained n-type germanium,” *Physical Review B*, vol. 87, p. 235313, (2013)
- [4] P. Markowich, “The Stationary Semiconductor Device Equations”, Springer-Verlag, Wien, 1986
- [5] S. Selberherr, “Analysis and Simulation of Semiconductor Devices” Springer-Verlag, Wien, 1984
- [6] M. Burger, R. Pinnau, “Fast optimal design of semiconductor devices,” *SIAM Journal on Applied Mathematics* vol. 64 (2003)

- [7] M. Hinze, R. Pinnau, “An optimal control approach to semiconductor design,” *Math. Models Methods Appl. Sci.* vol. 12, pp. 89–107 (2002)
- [8] M. Hinze, R. Pinnau, “Second-Order Approach to Optimal Semiconductor Design,” *J Optim Theory Appl* vol. 133, pp. 179–199 (2007)
- [9] U. Bandelow, H. Gajewski, R. Hünlich, “Thermodynamics-based modeling of edge-emitting quantum well lasers,” *WIAS* (2004)
- [10] O.A. Ladyzhenskaya, N.N. Ural'tseva, “Linear and quasilinear elliptic equations,” *Academic Press* vol. 46 (1968)
- [11] D. Peschka, N. Rotundo, M. Thomas, “Towards doping optimization of semiconductor lasers,” *WIAS Preprint* No. 2180, (2015)
- [12] J. Naumann, “On the existence of weak solutions of the stationary semiconductor equations with velocity saturation” *Le Matematiche*, Italy, 44, may. (1991)
- [13] J. Naumann, M. Wolff, “A uniqueness theorem for weak solutions of the stationary semiconductor equations,” *Appl. Math. Optim.* 24, 223–232 (1991)
- [14] G. Stampacchia, “Contributi alla regolarizzazione delle soluzioni dei problemi al contorno per equazioni del secondo ordine ellittiche” *Ann. Scuola Norm. Sup. Pisa* 3, 12, 223–245 (1958)
- [15] D.L. Scharfetter, H.K. Gummel, “Large signal analysis of a silicon Read diode”. *IEEE Trans. Electron. Dev.* 16, 64–77 (1969)
- [16] I. Fonseca, G. Leoni, “Modern Methods in the Calculus of Variations: L^p -spaces,” Springer, 2007.
- [17] H. Gajewski, “On the existence of steady-state carrier distributions in semiconductors.” *Probleme und Methoden der Mathematischen Physik. Teubner-Texte zur Mathematik* 63 (1984), 76–82.
- [18] K. Gröger, “On steady-state carrier distributions in semiconductor devices”, *Applikace matematiky* 32 (1987), No. 1, 49–56.